

## Future sudden singularities in Palatini cosmology

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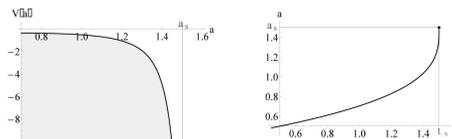
**Abstract.** We show that future singularities which have appeared in the Palatini cosmological models investigated in [1] are of finite size at finite time type [2].

Doomsday scenarios for cosmological models are subject of numerous investigations and speculations. For example the standard LCDM model predicts a thermal death caused by the expansion forever. In this short note we solve similar problem for the Palatini cosmological models investigated already in [1].

First of all we recall that, similarly to LCDM case, Palatini models are described by dynamical systems of Newtonian type. They are fully determined by effective potential functions  $V(a)$  defined by Friedmann equation (see [1]). The evolution of the Universe described by such system is given by:

$$\frac{\dot{a}^2}{2} = -V(a) \Rightarrow \ddot{a} = -\frac{\partial V}{\partial a} \tag{1}$$

where  $a$  denotes the cosmic (FRWL) scale factor. From the potential plot (fig . 1) one can see that:



**Figure 1.** Left panel: typical future singularity of the potential function for cosmological models presented in [1]. Right panel: example of time dependency of the scale factor near singularity.

$$\lim_{a \rightarrow a_s^-} V(a) \rightarrow -\infty, \quad \lim_{a \rightarrow a_s^-} \ddot{a} > +\infty \tag{2}$$

as well as

$$\lim_{a \rightarrow a_s^-} \dot{a} = +\infty \tag{3}$$

for some finite  $a_s < \infty$  singular (final) point.

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The question is whether the singular (finite) point has to be reached at finite time  $t_s < \infty$ . Assuming (a' contrario)  $\lim_{t \rightarrow +\infty} a(t) = a_s^-$  gives rise to

$$\lim_{t \rightarrow +\infty} \dot{a}(t) = 0^+ \quad (4)$$

which contradicts with (3). Therefore  $t_s$  has to be finite (sudden).

As illustrative examples of time (analytic) dependency nearby singular point one can consider the family of functions (see (fig. 1) and [2])

$$a(t) = a_s - b(t_s - t)^\alpha \quad (5)$$

with  $b > 0$  and  $0 < \alpha < 1$ . Thus  $\dot{a} = b\alpha(t_s - t)^{\alpha-1}$  and  $\ddot{a} = b\alpha(1 - \alpha)(t_s - t)^{\alpha-2}$ . It is worth to mention that typical values of  $a_s$  obtained in [1], after constraining model parameters by astrophysical data, are  $1.4 < a_s < 2.5$ .

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