

## Investigation of $0^+$ states in mercury isotopes after two-neutron transfer

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**Abstract.** Using the high-resolution Q3D magnetic spectrograph at the Maier-Leibnitz Laboratory (MLL) Tandem accelerator in Munich, we studied  $0^+$  excitations in the mercury isotopes  $^{198}\text{Hg}$ ,  $^{200}\text{Hg}$ , and  $^{202}\text{Hg}$  after two-neutron transfer. We only observed 4–6 excited  $0^+$  states per nucleus up to about 3-MeV excitation energy, far fewer than in other experiments of this  $(p, t)$  campaign. The results reveal a sharp drop in the number of low-lying  $0^+$  states towards the  $^{208}\text{Pb}$  shell closure. We discuss the low-energy  $0^+$  state density as a function of the valence nucleon number  $N_{\text{val}}$ . The  $0^+$  excitation energies and the measured  $(p, t)$  transfer cross sections indicate a structural change throughout the Hg isotopes, with the most notable result being the peaking in the cross section of the low-lying excited  $0_2^+$  state in  $^{200}\text{Hg}$ .

### 1 Introduction

Because of its capability to measure the characteristic forward-peaking of  $L = 0$  transfers, the Q3D magnetic spectrograph [1] at the Maier-Leibnitz Laboratory (MLL) Tandem accelerator in Munich – in combination with the focal plane detector [2] – has turned out to be a very successful instrument for the identification of  $0^+$  excitations. This feature has been extensively used to identify  $0^+$  states throughout the rare-earth region from gadolinium up to platinum [3–7]. For some nuclei, the number of identified  $0^+$  states increased sharply with the analysis of the Q3D measurement. The large number of low-lying  $0^+$  states in  $^{154}\text{Gd}$  [8] was interpreted as a new signature for the shape-phase transition from spherical to deformed nuclei [9]. The unexpected high number of observed  $0^+$  excitations triggered various calculations reproducing the density and distribution of  $0^+$  states [10–13]. Now, with the experiments on the even mercury isotopes  $^{198-202}\text{Hg}$  [14, 15], we move further to the  $^{208}\text{Pb}$  shell closure. This allows us to investigate  $0^+$  states in the vicinity of the prolate-oblate shape-phase transition in the

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**Table 1.** Assigned  $0^+$  states in  $^{198}, ^{200}, ^{202}\text{Hg}$  and their relative  $R(5/17.5)$  ratio. Newly assigned  $0^+$  states are marked with an asterisks, tentative ( $0^+$ ) assignments are denoted in *italic*. Each Hg isotope is labeled with its corresponding valence neutron number  $N_{\text{val}}$ .

$^{198}\text{Hg}$ ( $N_{\text{val}} = 10$ )		$^{200}\text{Hg}$ ( $N_{\text{val}} = 8$ )		$^{202}\text{Hg}$ ( $N_{\text{val}} = 6$ )	
Energy (keV)	$R(5/17.5)$	Energy (keV)	$R(5/17.5)$	Energy (keV)	$R(5/17.5)$
0.0 (0)	13.3 (4)	0.0 (0)	13.4 (4)	0.1 (1)	15.7 (4)
1401.0 (3)	9.4 (15)	1029.3 (1)	4.8 (2)	1411.0 (3)	12.3 (18)
1646.4 (8) *	>3.35	1515.5 (3)	10.5 (12)	1643.0 (3)	7.3 (5)
1779.6 (2)	4.9 (6)	1856.6 (2)	8.0 (21)	<i>1655.8 (13)</i>	<i>3.1 (13)</i>
		2246.1 (2) *	5.0 (7)	<i>1778.9 (6)</i>	<i>4.5 (15)</i>
		2331.8 (3)	3.2 (9)	<i>2126.7 (7)</i>	<i>3.2 (17)</i>
		2475.2 (1) *	4.4 (5)	<i>2570.7 (10)</i>	<i>2.3 (25)</i>
				2598.5 (2) *	4.3 (4)
				<i>2685.7 (5)</i>	<i>4.7 (32)</i>

Hf-Hg region [16] and to test if the prolate-oblate shape-phase transition affects the low-energy  $0^+$  state density of these nuclei.

## 2 Experiment & Analysis

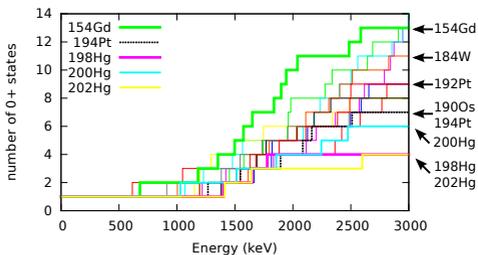
To identify  $L = 0$  transfers from the Hg  $0^+$  ground state (GS) in the target material to an excited  $0^+$  state of the Hg isotope of interest, e.g., a transfer from the  $^{200}\text{Hg}$  GS to an excited  $0^+$  state in  $^{198}\text{Hg}$ , we measured spectra at  $5^\circ$ ,  $17.5^\circ$ , and  $30^\circ$  laboratory angle relative to the incoming proton beam. A more detailed description of the experimental setup at the Q3D, the specifications of the enriched Hg targets, and the analysis is given in Refs. [14, 15]. The characteristic forward-peaking of  $L = 0$  transfers was determined by evaluating the ratio  $R(5/17.5) \equiv \sigma(5^\circ)/\sigma(17.5^\circ)$  for each observed transfer. As a safe lower limit to prevent incorrect  $0^+$  assignments, we used  $R(5/17.5) > 3$ .

## 3 Results

The resulting  $0^+$  assignments for  $^{198}\text{Hg}$ ,  $^{200}\text{Hg}$ , and  $^{202}\text{Hg}$  –based on the  $R(5/17.5)$  ratio of our data– are listed in Table 1. In total, we assigned four new  $0^+$  states up to  $\sim 3$ -MeV excitation energy in these three Hg isotopes investigated. Tentative assignments are mostly due to a poor population of these particular states. Absolute cross sections for each state are given in Refs. [14, 15]. Some  $0^+$  assignments in the data sheets [17–19] were not confirmed. Whenever possible, we tested assignments from the  $(p, t)$  data for consistency with  $\gamma\gamma$  coincidence data of a fusion-evaporation experiment [20] and data of a neutron-capture experiment at ILL Grenoble [21, 22].

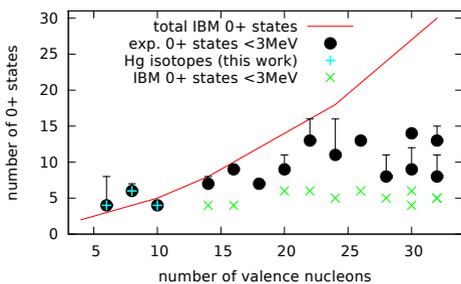
## 4 Discussion

Figure 1 shows the complete data set of  $0^+$  excitations assigned in this  $(p, t)$  campaign at the Q3D spectrograph, ranging from Gd to Hg isotopes. One notes a significantly higher  $0^+$  state density in  $^{154}\text{Gd}$ , which was interpreted as a new signature for the shape-phase transition from spherical to deformed nuclei [8]. The prolate-oblate shape-phase transition observed in Ref. [16] does not show an effect on the low-energy  $0^+$  density. The number of low-lying  $0^+$  states rather declines with larger nucleon mass, down to only four low-energy  $0^+$  assignments in  $^{198}\text{Hg}$  and  $^{202}\text{Hg}$ .



**Figure 1.** Histogram of low-energy  $0^+$  states observed in this  $(p, t)$  Q3D campaign. Data are taken from Refs. [3–7, 14, 15]. The labeled nuclei on the right-hand side show the decline of  $0^+$  state density throughout the prolate-oblate shape-phase transitional region, with  $^{194}\text{Pt}$  concluded to be the closest to the critical point [16]. Figure based on Ref. [14].

The effect of a decline in the number of low-energy  $0^+$  states towards the  $^{208}\text{Pb}$  shell closure is reproduced by Interacting Boson Model (IBM) [23] calculations and illustrated in Fig. 2. Please see Ref. [14] for more details on the calculations. Figure 2 shows that the number of observed  $0^+$  states strongly increases with  $N_{\text{val}}$  as the valence space expands, but saturates near midshell starting at about  $N_{\text{val}} = 22$  valence nucleons.



**Figure 2.** Number of  $0^+$  states up to 3 MeV assigned in this  $(p, t)$  campaign as a function of valence nucleon number  $N_{\text{val}}$ . The error bars include tentative assignments. The red line corresponds to the maximum number of  $sd$  IBM  $0^+$  states, where as the green crosses indicate the calculated number of  $0^+$  states using realistic parameters, if available. Based on Ref. [15].

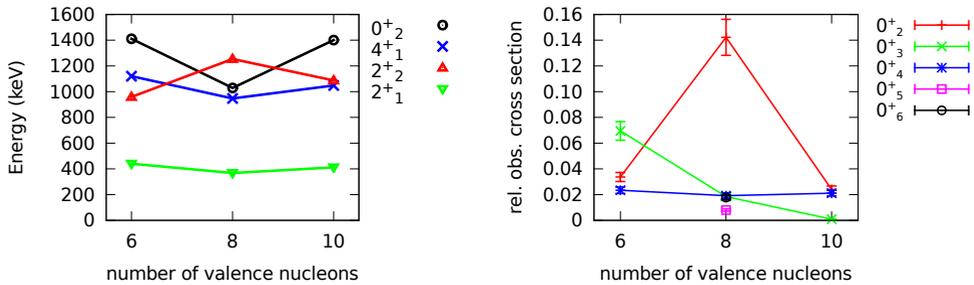
Figure 2 shows a small peak at  $N_{\text{val}} = 8$ , corresponding to  $^{200}\text{Hg}$ . By directly comparing further properties of the investigated Hg isotopes one notes more differences shown in Fig. 3: the sequence and level energies of the low-lying states change rapidly at  $^{200}\text{Hg}$  and the  $0_2^+$  state in  $^{200}\text{Hg}$  is strongly populated with about 12% of the ground-state cross section. Historically,  $0^+$  state two-nucleon transfer cross sections approaching or exceeding 15% have signaled structural effects such as phase transitional regions or shape coexistence [9]. In Ref. [24], the enhanced  $0_2^+$  cross section in  $^{200}\text{Hg}$  has been associated with an oblate single particle energy gap, but other explanations like mixing or coexistence cannot be ruled out. Another indicator for structural changes, the two-neutron separation energy  $S_{2n}$  (or the differential  $\delta S_{2n}$ ), shows at most a weak anomaly [25, 26].

## 5 Conclusion

The experiments on the Hg isotopes complete a high-resolution  $(p, t)$  campaign using the Q3D spectrograph. We observe fewer low-energy  $0^+$  states in  $^{198}, ^{200}, ^{202}\text{Hg}$  than in other nuclei investigated in this  $(p, t)$  campaign. Plotted as a function of valence nucleon number  $N_{\text{val}}$ , we note a sharp drop in the number of  $0^+$  states in the near-magic region compared to the transitional and deformed regions, but a saturation towards midshell. The low-energy  $0^+$  density seems not to be affected by the prolate-oblate shape phase transition [16] in the Hf-Hg region.

The  $0_2^+$  state in  $^{200}\text{Hg}$  has a large two-nucleon transfer cross section, which might be an indicator for a structural change throughout the Hg isotopes at  $N = 120$ . To understand this effect, it would be very helpful to learn more about the  $E0$  ground-state transition of the  $0_2^+$  state.

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**Figure 3.** Low-energy states (left) and relative observed cross sections of  $0^+$  excitations (right) for  $^{198, 200, 202}\text{Hg}$  as a function of valence nucleon number  $N_{\text{val}}$ . The significant changes in excitation energies for the low-lying states and the unusual strong population of the  $0^+_{2}$  state at  $N_{\text{val}} = 8$ , corresponding to  $^{200}\text{Hg}$ , are often interpreted as an indicator for a structural change. Based on Ref. [15].

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