

Pseudoscalar mesons parton distribution amplitudes

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Abstract. We compute all kaon and pion parton distribution amplitudes (PDAs) to twist-three and find that all of the pion's distributions are symmetric and all of the kaon's distributions are skewed in favour of the heavier s -quark, which are clear signals of dynamical chiral symmetry breaking (DCSB). And we can also notice that only the pseudotensor PDA can reasonably be approximated by its conformal limit. At any realistic energy scale, the pseudovector and pseudoscalar PDAs differ markedly from the corresponding functions in QCD's conformal limit. We can use these PDAs as the inputs to study hard exclusive processes and get trustable results.

1 Introduction

The PDAs of pion and kaon are very crucial to study hard exclusive processes[1, 2]. In the previous studies, people can only use the expressions of PDAs in the QCD's conformal limit[3, 4]. Now people can recalculate mesons' more realistic PDAs by using a new method. With this new method[5, 6], we calculate six PDAs of pion and kaon to twist-three.

2 Distribution amplitudes and Bethe-Salpeter wave functions

A pseudoscalar meson, $P_{\bar{q}f}(q)$, has three two-particle distribution amplitudes on the light front, they can be expressed as[7]

$$f_P \varphi_P(u) = N_c \text{tr} Z_2 \int_{dk}^{\Lambda} \delta(n \cdot k_\eta - u n \cdot q) \gamma_5 \gamma \cdot n \chi_P^q(k_\eta, k_{\bar{\eta}}) \quad (1a)$$

$$i\rho_P^\zeta \omega_P(u) = N_c \text{tr} Z_4 \int_{dk}^{\Lambda} \delta(n \cdot k_\eta - u n \cdot q) \gamma_5 \chi_P^q(k_\eta, k_{\bar{\eta}}) \quad (1b)$$

$$\frac{1}{4} \rho_P^\zeta v'(u) = N_c \text{tr} Z_4 \int_{dk}^{\Lambda} \delta(n \cdot k_\eta - u n \cdot q) \gamma_5 \sigma_{\mu\nu} q_\mu n_\nu \chi_P^q(k_\eta, k_{\bar{\eta}}) \quad (1c)$$

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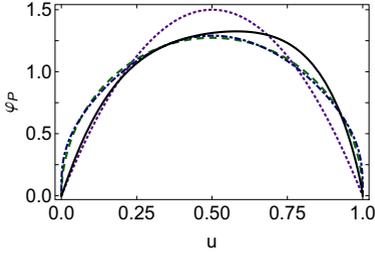


Figure 1. Two-particle twist-two PDAs, computed at ζ_2 : dot-dashed curve (dark blue) – pion, $\varphi_\pi^{\text{DB}}(u; \zeta_2)$; solid curve (black) – kaon, $\varphi_K^{\text{DB}}(u; \zeta_2)$; dashed curve (dark green) – $\varphi_\pi^{\text{model}}(u)$, Eq. (5); and dotted curve (indigo) – asymptotic distribution, Eq. (4).

where $N_c = 3$; the trace is over spinor indices; \int_{dk}^Λ is a Poincaré-invariant regularisation of the four-dimensional integral, with Λ the ultraviolet regularisation mass-scale; $Z_{2,4}(\zeta, \Lambda)$ are renormalisation constants; ζ is the renormalisation scale; $\chi_P^q(k_\eta, k_{\bar{\eta}})$ is the Bethe-Salpeter wave function; and f_P, ρ_P^ζ are, respectively, the pseudovector and pseudoscalar projections of the meson's Bethe-Salpeter wave function onto the origin in configuration space. Every PDA should be normalized, *viz.*

$$\int_0^1 du \{ \varphi_P(u; \zeta), \omega_P(u; \zeta), \nu_P(u; \zeta) \} = 1. \quad (2)$$

We can numerically solve the coupled integral equations[8–11], namely, QCD's gap and Bethe-Salpeter equations, to get the expressions of quark propagators and the mesons' Bethe-Salpeter amplitudes. Here we solve these equations in two symmetry-preserving truncations[12]: the rainbow-ladder (RL) truncation, which is the most widely used in the calculation; and the modern DCSB-improved (DB) kernels, which can give the most realistic predictions by now.

Then we can employ the procedure in previous works to calculate PDAs[5, 6, 13]. First we need to calculate PDAs' moments

$$\langle u_\Delta^m \rangle_\phi = \int_0^1 du (2u - 1)^m \phi_P(u), \quad (3)$$

where $\phi = \varphi_P, \omega_P, \nu_P$, with these moments, we can reconstruct the PDAs by using Gegenbauer polynomials of order α , $\{C_n^\alpha(2u - 1) | n = 0, \dots, \infty\}$. Gegenbauer polynomials form a complete orthonormal set on $u \in [0, 1]$ with respect to the measure $[u(1 - u)]^{\alpha-}$, $\alpha_- = \alpha - 1/2$. They therefore enable reconstruction of any function defined on $u \in [0, 1]$.

3 Numerical results

3.1 Pseudovector

We depict our pseudovector results in Fig. 1, and compare them with the asymptotic two-particle distribution

$$\varphi^{\text{cl}}(u) = 6u(1 - u) \quad (4)$$

which is the asymptotic distribution in QCD's conformal limit, and with

$$\varphi_\pi^{\text{model}}(u) = (8/\pi) \sqrt{u(1 - u)}, \quad (5)$$

which is a model result, it is practically indistinguishable from the DSE prediction.

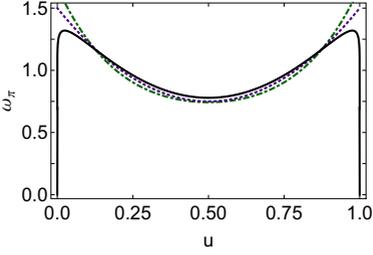


Figure 2. Pion two-particle, twist-three PDAs, computed at ζ_2 . Solid curve (black), $\omega_\pi(u; \zeta_2)$; dot-dashed curve (dark green), QCD sum rules estimate [14]; and dotted curve (indigo), ω_P^{asy} in Eq. (6).

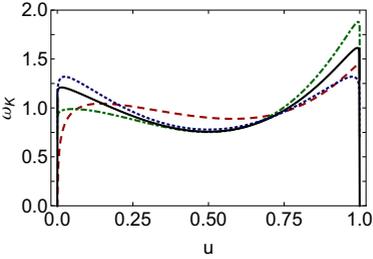


Figure 3. Kaon two-particle, twist-three PDAs, computed at ζ_2 . Solid curve (black), $\omega_K^{\text{DB}}(u; \zeta_2)$; dot-dashed curve (green), $\omega_K^{\text{RL}}(u; \zeta_2)$; dashed curve (red), QCD sum rules estimate [14]; and dotted curve (blue), $\omega_\pi(u; \zeta_2)$.

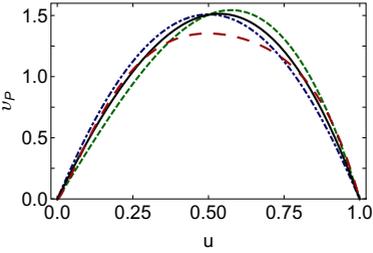


Figure 4. Pseudotensor two-particle, twist-three PDAs, computed at ζ_2 : dot-dashed curve (dark blue) – pion, $\nu_\pi(u)$; solid curve (black) – kaon in DB truncation, $\nu_K^{\text{DB}}(u)$; dashed curve (dark green) – kaon in RL truncation $\nu_K^{\text{RL}}(u)$; and long-dashed curve (red) – QCD sum rules result for the kaon from Ref. [14]. We do not plot the asymptotic form because it is effectively indistinguishable from our prediction for the pion.

3.2 Pseudoscalar

We depict our pseudoscalar results in Fig. 2 and 3, and compare them with the asymptotic distribution

$$\omega_P^{\text{asy}}(u) = 1 + \frac{1}{2}C_2^{(1/2)}(2u - 1). \quad (6)$$

3.3 Pseudotensor

We depict our pseudotensor results in Fig. 4, we do not plot the asymptotic form because it is effectively indistinguishable from our prediction for the pion..

4 Discussion and conclusion

We can get many important information from our results. First, we can see that all of the pion's distributions are symmetric and all of the kaon's distributions are skewed in favour of the heavier s -quark, which are clear signals of DCSB. In kaon's pseudoscalar case, the RL PDA is more skewed than the DB result, *viz.* the RL truncation allocates more fraction of the in-kaon condensate to its valence s -quark. The RL-kernels ignore DCSB in the quark-gluon vertex, meanwhile the DB-kernels

put DCSB into the quark-gluon vertex, so the the DB-PDA is more realistic. In kaon's pseudovector and pseudotensor case, we can get similar observations. It can be shown that in all PDAs, the $SU(3)$ flavour-symmetry breaking is a 13% effect[5, 15]. And we can also notice that only the pseudotensor PDA can reasonably be approximated by its conformal limit. At any realistic energy scale, the pseudovector and pseudoscalar PDAs differ markedly from the corresponding functions in QCD's conformal limit.

Now we have calculated pion's and kaon's all the PDAs to twist-three, which are remarkably different from their expressions in the conformal limit except for the pseudotensor case. We can use these PDAs as the inputs to study hard exclusive processes and get trustable results.

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