

## $\Lambda nn$ bound state with three-body potential

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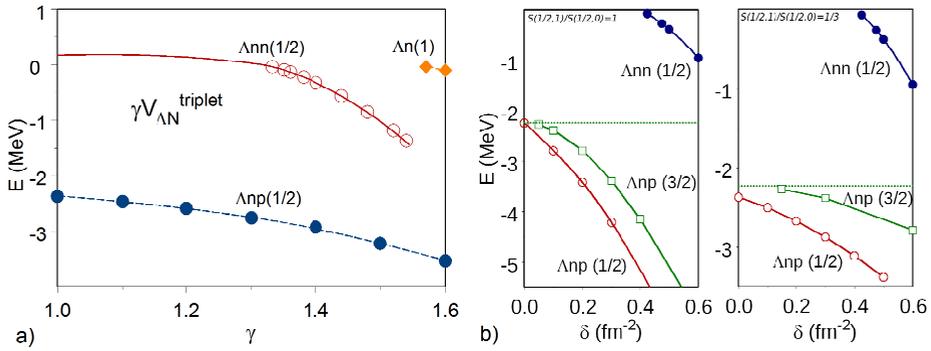
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**Abstract.** Formalism of the configuration-space Faddeev equations is applied to analyze the bound state problem of  $\Lambda NN$  system. The simulated NSC97f potential for  $\Lambda N$  interaction and semi-realistic  $NN$  potentials are used for calculations. We predict a broad  $\Lambda nn$  resonance near threshold. A spin dependent three-body  $\Lambda NN$  potential has been proposed to explain experimental evidence for  ${}^3_{\Lambda}n$  bound state. The potential is defined by a Gaussian function of a hyper-radius. It was found that such potential with appropriate spin dependence does not support a bound state of  ${}^3_{\Lambda}n$ . A short range form of the potential generates non-physical solutions of the Faddeev equations.

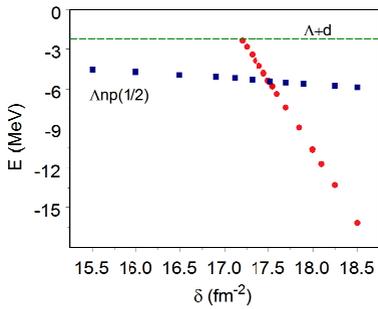
The HypHI Collaboration has recently reported evidence for a bound state of  $\Lambda + n + n$  system [1]. However, the bound  ${}^3_{\Lambda}n$  state has not been found under theoretical analysis (see for instance [2]). In the present work we construct a phenomenological three-body  $\Lambda NN$  potential with spin dependence that forms a bound state in the channel  $\Lambda + n + n$  ( $T=1, S=1/2$ ). Such potential does not affect the ground state energy of  ${}^3_{\Lambda}H$  hypernucleus fixed by the experimental value of 0.13(5) MeV for the hyperon separation (or binding) energy

Formalism of the configuration-space Faddeev equations was applied for  $\Lambda + n + n$  and  $\Lambda + n + p$  systems within an  $s$ -wave approach. The simulated NSC97f potential [3] for  $\Lambda N$  interaction and MT-I-III semi-realistic  $NN$  potential [4] were used for calculations. This model reproduces well the hyperon binding energy for  ${}^3_{\Lambda}H$ , with the value of 0.14 MeV (see also Ref. [5]). We have found a broad resonance in the channel  $\Lambda + n + n$ . The results of the calculations are presented in figure 1a). The method of analytical continuation in the coupling constant is used to calculate the parameters of the resonance [6]. The coupling constant is the coefficient  $\gamma$  which scales triplet component of the  $\Lambda N$  potential. The complex value of the Pade approximant calculated by set of negative energies (circles in figure 1a) for  $\gamma=1$  gives the energy and width of resonance:  $E(\gamma = 1) = E_r - i\frac{\Gamma}{2}$ . Our result for the energy of the resonance is 0.2 MeV, that agrees with [7]. Within the presented model, the three-body  $\Lambda NN$  potential is treated as a perturbation of the Hamiltonian. It is defined as one range Gaussian:  $V_{3bf}(\rho) = -\delta \exp(-\alpha\rho^2)S(s_{\Lambda}, s_{NN})$ , where  $\rho$  is the hyper-radius:  $\rho^2 = x^2 + y^2$ , with  $x, y$  the mass scaled Jacobi coordinates. The function  $S(s_{\Lambda}, s_{NN})$  depends on the spin variables of a hyperon ( $s_{\Lambda}$ ) and  $NN$  pair ( $s_{NN}$ ). Two free parameters  $\delta$  and  $\alpha$  have to be adjusted. We assume the ratio  $S(1/2, 1)/S(1/2, 0)$  to be equal 1 or 1/3 [2]. The calculations for  $\alpha=0.1 \text{ fm}^{-2}$  are shown in figure 1b). The three-body  $\Lambda NN$  potential resulted in a bound state of  ${}^3_{\Lambda}n$  generates  $\Lambda np$  over-bound states for  $S=1/2$  and  $3/2$ . The spin dependence, which provides a bound state of  ${}^3_{\Lambda}n$  and keeps the  ${}^3_{\Lambda}H$  ground state energy, cannot be considered as realistic. It was found that the short range form (when  $\alpha>0.1 \text{ fm}^{-2}$ ) of the potential generates non-physical solutions of the Faddeev equations (see figure 2).

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**Figure 1.** a) Bound state  $\Lambda + n + n$  and  $\Lambda + n + p$  energies (circles) calculated for different values of  $\gamma$ . The line for  $\Lambda nn$  represents the real part of the Padé approximant. The resonance energy corresponds to the value of the approximant for  $\gamma=1$ . The three-body potential is not taken into account. The energy is measured from the  $\Lambda + n + n$  threshold. b)  $\Lambda + n + n$  and  $\Lambda + n + p$  energies (solid and open circles) calculated for different values of  $\delta$  ( $\alpha=0.1$  fm $^{-2}$ ). The  $\delta$  parameter is scaled by the value  $\hbar^2/m$ ,  $m$  is mass of nucleon. (Left)  $S(1/2, 1)/S(1/2, 0)=1$ , (Right)  $S(1/2, 1)/S(1/2, 0)=1/3$ .



**Figure 2.** Bound state  $\Lambda + n + p$  energies calculated for different values of the  $\delta$  parameter of the three body potential ( $\alpha=0.9$  fm $^{-2}$ ). The squares (circles) correspond to physical (non-physical) solution of the Faddeev equations.  $\Lambda + d$  threshold is shown by horizontal dashed line.

We have found a broad  $\Lambda nn$  resonance near threshold with the energy of 0.2 MeV and the width of 2 MeV. It is not possible to construct a three-body potential for both  $\Lambda nn$  and  $\Lambda np$  systems with reasonable spin dependence under the condition of existence of a  ${}^3_{\Lambda}n$  bound state. The fact of the appearance of non-physical solutions reflects the known problem of implementation of three-body force into the Faddeev equation formalism.

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