

# Superheavy Nuclei in the Quark-Meson-Coupling Model

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**Abstract.** We present a selection of the first results obtained in a comprehensive calculation of ground state properties of even-even superheavy nuclei in the region of  $96 < Z < 136$  and  $118 < N < 320$  from the Quark-Meson-Coupling model (QMC). Ground state binding energies, the neutron and proton number dependence of quadrupole deformations and  $Q_\alpha$  values are reported for even-even nuclei with  $100 < Z < 136$  and compared with available experimental data and predictions of macro-microscopic models. Predictions of properties of nuclei, including  $Q_\alpha$  values, relevant for planning future experiments are presented.

## 1 Introduction

One of the main difficulties in solving the nuclear many-body problem is to incorporate nuclear medium effects into the calculation. In contrast to, for example, elementary quantum electrodynamics, where the force between two electric charges is quantified by the Coulomb law and forces among many electrons can be calculated precisely using the principle of superposition, no equivalent approach exists in nuclear physics. Scattering experiments with free nucleons provide information about bare N-N potentials, but an involved numerical treatment is necessary to transform these potentials to a form which describes the forces between hadrons in the hadronic environment. There have been many theoretical attempts to overcome this problem but they all suffer from the uncertainty associated with many empirical parameters which are not uniquely constrained, thus preventing reliable predictions of nuclear properties in regions not accessible to experiment. One particular area of interest concerns the transuranic superheavy nuclei, where one may hope to discover new chemical elements, not existing in nature but synthesized in the laboratory. Theoretical predictions guiding such experiments are of major interest, especially information about spherical and deformed shell closures and related islands of stability. Given the ambiguity in models based on nucleonic degrees of freedom, alternative approaches have been sought.

The idea of modeling the effective nuclear interaction in-medium using quark degrees of freedom originated in the late 1980s with work by Guichon and collaborators [1–3]. They suggested that the origin of nuclear many body forces and of the saturation of nuclear forces can be found

in the modification of the structure of a nucleon when it is imbedded in a medium consisting of other nucleons. The model, referred to as the Quark-Meson-Coupling (QMC) model, has been applied to finite nuclei [4, 5], nuclear matter and neutron stars (see e.g. [6, 7] as well as hypernuclei [8] and a variety of other problems [9]). The model is based on the idea that, instead of the usual treatment of modeling nuclear forces through exchange of mesons coupled to nucleons, taken as point-like particles, this exchange takes place directly between quarks in different nucleons, taken as a cluster of confined valence quarks having a structure in the form, for example, of the MIT bag. When the quarks in one nucleon interact *self-consistently* with the quarks in the surrounding nucleons by exchanging a  $\sigma$  meson, the effective mass  $M_N^*$  of the nucleon is no longer linear in the scalar mean field ( $\sigma$ ). It is expressed as

$$M_N^* = M_N - g_{\sigma N}\sigma + (d/2)(g_{\sigma N}\sigma)^2, \quad (1)$$

where  $g_{\sigma N}$ , the  $\sigma$  nucleon coupling constant in free space, is a parameter of the model. By analogy with electromagnetic polarizabilities, the coefficient  $d$ , calculated in terms of the nucleon internal structure, is known as the “scalar polarizability” [1]. The appearance of this term in the nucleon effective mass is sufficient to lead to nuclear saturation. This demonstrates a clear link between the internal structure of the nucleon and fundamental properties of atomic nuclei.

Models of nuclei in the region above  $Z = 100$  can be divided into two classes [10], mean field models with effective density dependent interactions in the Hartree-Fock approximation [11–17], and macro-microscopic models based on the liquid-drop model combined with Strutinsky-type shell corrections (see e.g. [18–21]). Although these models yield similar results for some bulk properties of superheavy nuclei, they differ in important details. Su-

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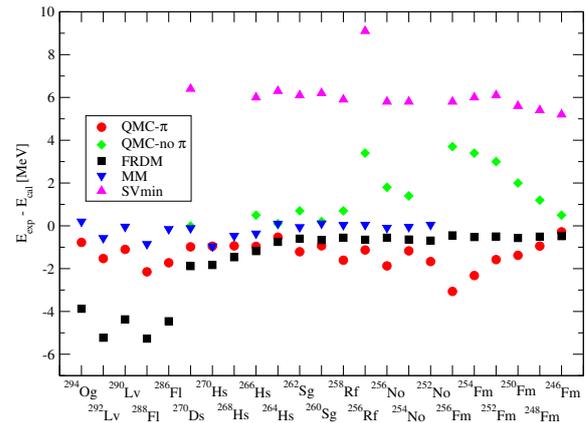
**Table 1.** Neutron numbers corresponding to proton and neutron drip lines, derived from the Fermi energy for isotopes of elements  $96 < Z < 136$

Z	N(p)	N(n)	Z	N(p)	N(n)
96	132	224	118	174	278
98	134	226	120	180	286
100	138	230	122	184	290
102	138	236	124	188	296
104	146	240	126	192	298
106	146	242	128	196	302
108	154	246	130	202	306
110	158	250	132	208	310
112	164	256	134	214	314
114	168	260	136	218	314
116	170	268			

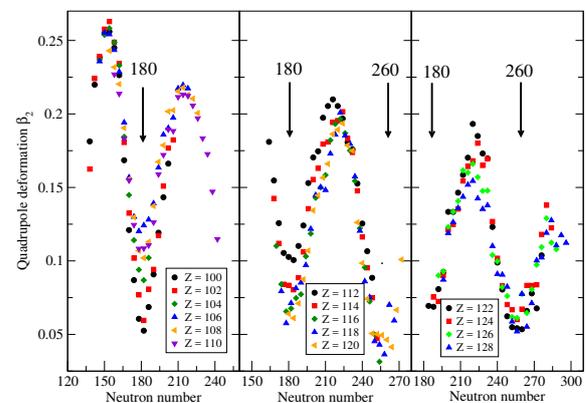
perheavy nuclei exist mainly because of their shell structure, in turn, dependent on level densities close to the Fermi level. It is therefore very important to understand their single-particle spectrum, in particular the location and magnitude of proton and neutron shell gaps. In addition, it may be expected that the competition between the Coulomb repulsion and the surface tension in many nucleon systems would play an important role and the sharp shell closures, familiar in lighter nuclei, may be diluted to form regions of a distinct spherical or deformed nature around a particular (N,Z) shell closure.

The various models predict different proton and neutron shell closures above  $^{208}\text{Pb}$ . The macroscopic–microscopic models yield a ‘canonical’ island of stability at  $Z = 114$ ,  $N = 184$  [10, 11], while the self-consistent mean-field models (SC) non-relativistic models predict  $Z = 126$ ,  $N = 172 - 184$  [11, 12] and the relativistic mean-field models (RF) favour  $Z = 120$ ,  $N = 172$  [11, 16]. The most recent comparison of traditional models applied to superheavy nuclei may be found in [14] but no way to reconcile these differences is offered.

It is therefore interesting to examine the predictions of the QMC model which belongs to the class of mean field models but differs in key ways from both non-relativistic models with the Skyrme or Gogny interaction and classical relativistic field models in important details [5]. In this paper we present a selection of the first results obtained in a comprehensive study of the ground state properties of spherical and axially symmetrical even-even superheavy nuclei in the region  $96 < Z < 136$  and  $118 < N < 320$ . A new version of the QMC model,  $\text{QMC}\pi$ , is used in this work, where the published model [5] has been extended to include pion-exchange. The calculational procedure is described in section 2 and the resulting ground state binding energies of selected isotopes and the evolution of the deformation parameter  $\beta_2$  for bound isotopes of elements with  $100 < Z < 128$  and isotones with  $148 < N < 204$ , along with the  $Q_\alpha$  values are presented in section 3. Section 4 contains predictions of the  $\text{QMC}\pi$  model for properties of superheavy nuclei considered in planning future experiments, followed by section 5 summarizing the current results and outlining future work.



**Figure 1.** (Color online). Ground state binding energies of selected ‘benchmark’ even-even superheavy nuclei. The experimental data were taken from [27, 28].

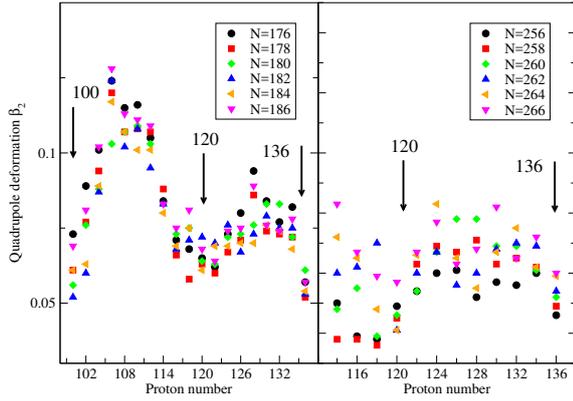


**Figure 2.** (Color online). Quadrupole deformation calculated in  $\text{QMC}\pi$  for isotopes with proton number  $100 < Z < 128$ .

## 2 Method of calculation

The  $\text{QMC}\pi$  energy density functional (EDF) has been constructed in the same way as described in [5], except that the explicit pion exchange has been included in the Hamiltonian. We used HF+ BCS code SKYAX allowing for axially symmetric and reflection-asymmetric shapes, adapted by P.-G. Reinhard [22] for use with the QMC EDF. The best parameter set was sought using the experimental data set by Klüpfel et al. [23] and the fitting package POUNDERS [24, 25]. The volume pairing in the BCS approximation as been adopted with proton and neutron pairing strength fitted to data in [23]. We note that the addition of the explicit pion exchange in the model did not increase the number of parameters beyond the four used in the previous work, namely  $G_\sigma$ ,  $G_\omega$ ,  $G_\rho$  and  $M_\sigma$ , but its addition was reflected in slight changes (less than 5%) from the values reported in [5]. The new parameter set is compatible with nuclear matter properties  $E_0 = -15.8$  MeV,  $\rho_0 = 0.153$  fm $^{-3}$ ,  $K_0 = 319$  MeV,  $S_0 = 30$  MeV and  $L = 27$  MeV.

The selection of isotopes studied was limited by the requirement that only bound single particle states are occupied; that is, the Fermi level does not increase to the



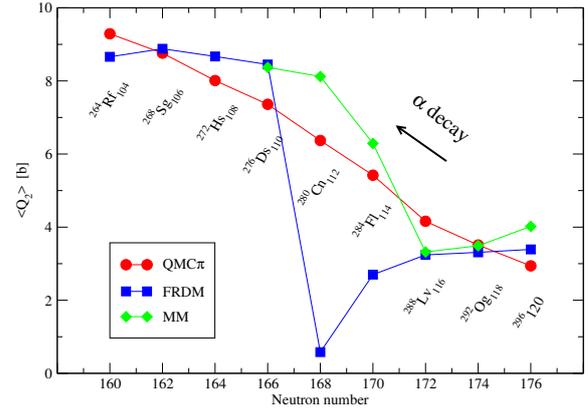
**Figure 3.** (Color online). The same as figure 2 but for isotones with neutron number  $176 < N < 186$  (left panel) and  $256 < N < 266$  (right panel).

continuum. This condition limited the range of neutron number for each element as given in Table 1.

### 3 Ground state properties and decay of selected superheavy nuclei

#### 3.1 Binding energies and shapes

Figure 1 illustrates the difference between experimental and calculated ground state binding energies of selected superheavy nuclei for which the experimental data is known well enough that they may be regarded as 'benchmark' data [26] to calibrate theoretical models for subsequent use in uncharted areas of the nuclear landscape, notably in the transuranic region. These results illustrate a rather dramatic change from a significant underbinding, yielded by the mean field model with the Skyrme SVmin parameter set [23] and, to a lesser extent, predicted by QMC-without-pion model in the  $Z=100 - 104$  region [5], to a slight overbinding with the present QMC $\pi$ , microscopic-macroscopic models Finite-Range-Droplet-Model (FRDM) [18] and of Muntian et al. (MM) [20, 21]. It is interesting to observe that the inclusion of pion-exchange in the QMC model has a significant effect on the ground state binding energies. As mentioned already in the introduction, one of the main motivations for modeling the superheavy region is to locate regions of relative stability. These regions would be the best targets for synthesis of either new elements or unknown isotopes for elements already known. The deformation parameters,  $\beta_2$ , determined in the QMC $\pi$  model by minimizing the ground state binding energies, are plotted with respect to neutron (proton) number in Figs. 2 and 3, respectively. The results in Fig. 2 clearly identify two regions of low deformation. Values of  $\beta_2$  less than about 0.1 are predicted for isotopes of all elements with  $100 < Z < 128$  with neutron number in the region around  $N \sim 180$  and, for  $Z > 114$  also in the vicinity of  $N = 260$ . The calculation of isotones with  $176 < N < 186$  reveals a strong indication for low deformation close to  $Z = 100, 120$  and  $136$  (or possibly higher), as illustrated in Fig. 3 (left panel). The



**Figure 4.** (Color online). Intrinsic quadrupole moments of even-even nuclei in the  $\alpha$ -decay chain of  $^{296}120$ , as calculated in the QMC $\pi$ , FRDM and MM models. See Fig. 1 in [29] and Ref. [30] for a comparison with other models and the text for more explanation.

isotones with  $256 < N < 266$  generally exhibit low deformation in the region of  $116 < Z < 136$  (right panel). The effect is less prominent at  $Z = 120$ , but still clearly present at  $Z = 136$ . The analysis of the quadrupole deformation in QMC $\pi$  shows no evidence for the  $Z = 114$  shell closure predicted in the FRDM and MM models nor at  $Z = 126$ , as suggested in SC models. However, it does agree with the results of RF models finding a shell closure at  $Z = 120$ . The neutron shell closures at  $N = 172$ , found in the SC and RF models is also not confirmed in this work, while the  $N=180$  region is somewhat close to the value  $N = 184$  found in the FRDM, MM and SC models. The region of sphericity we predict at  $N = 260$  and  $Z = 136$  is outside the reach of current experimental techniques but may potentially lead to a new island of stability. Further calculation, in particular of spectra of single-particle states, is needed to take this problem further.

Next we examine in more detail the development of the quadrupole deformation of isotopes along  $\alpha$  decay chains of special interest, as they are a prime tool for identification of new elements (keeping in mind that spontaneous fission is in competition with  $\alpha$  decay). In Fig. 4 we show the predicted mean intrinsic quadrupole moment calculated in this work. This is related to  $\beta_2$  as  $\langle Q_2 \rangle = (3/(4\pi))AR_0^2\beta_2$ , with  $R_0=1.2 A^{1/3}$ . We note that this definition [13] is consistent with the procedure used to calculate the deformation parameters in this work. There are alternative definitions in the literature, used for example in [29]. Therefore, only the qualitative neutron number dependence of  $\langle Q_2 \rangle$  can be compared between different calculations. We see a remarkably smooth decrease in deformation with increasing neutron number along the decay chain in the QMC $\pi$  model prediction, similar to that found by Shi et al. [29] in calculations using UNEDF1 and UNEDF1<sup>SO</sup> Skyrme forces. In contrast, the FRDM microscopic-macroscopic model predicts a distinctive spherical shape at  $N = 168$ . The MM model more or less shows a decrease in deformation with increasing neutron number with a slight tendency to sphericity at  $N = 172$ .

**Table 2.**  $Q_\alpha$  in MeV calculated in the QMC $\pi$ , FRDM and MM models. Experimental data without errors have been taken from nuclear mass tables [42], while entries with errors are from original papers [33, 35–37].

Z	N	$Q_\alpha^{\text{exp}}$	$Q_\alpha^{\text{QMC}\pi}$	$Q_\alpha^{\text{FRDM}}$	$Q_\alpha^{\text{MM}}$
102	150	8.548	8.20	8.82	8.53
102	152	8.226	8.10	7.97	8.06
102	154	8.581	8.00	8.57	8.36
104	152	8.926	9.40	8.75	8.93
104	154	9.190	9.10	9.32	9.29
106	154	9.901	10.10	9.93	9.95
106	156	9.600	10.00	9.61	9.49
108	156	10.591	11.00	10.57	10.59
108	158	10.346	10.60	9.69	10.04
108	160	9.62± 0.16	10.30	9.00	9.49
108	162	9.02± 0.03	10.20	8.69	8.78
110	160	11.117	11.10	10.30	11.36
112	172	10.230	10.40	8.69	9.76
114	172	10.33±0.08	11.30	9.39	9.76
114	174	10.08±0.06	11.00	9.16	10.32
116	174	11.00±0.08	11.60	11.12	11.08
116	176	10.80±0.07	11.40	10.82	11.06
118	176	11.81±0.06	12.20	12.28	12.11

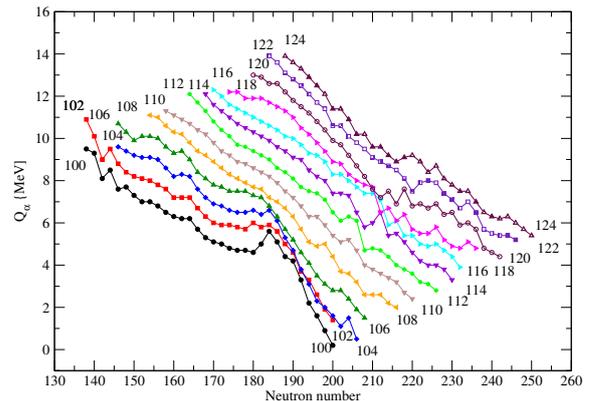
It is important to realize that deformation is treated differently in mean-field and micro-macro models (for details see [14]). Thus experimental data on quadrupole moments of superheavy elements may provide important information, leading to some distinction between the two classes of model to be used in the superheavy region. The recent progress in laser resonance ionization spectroscopy [31, 32] offers a prospect that electromagnetic moments, charge radii and isotope shifts of superheavy isotopes will be accessible experimentally in the near future.

### 3.2 $Q_\alpha$ values

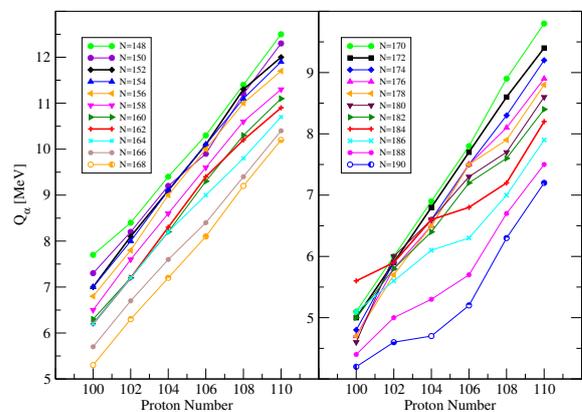
Knowledge of  $\alpha$  decay life-times is crucial to predicting properties of the  $\alpha$  decay chains of superheavy elements, which, as pointed out above, are crucial in the search for and detection of new elements and their isotopes. The  $\alpha$ -decay life-times are exponential functions of the energy release,  $Q_\alpha$ , in the decay, which, in turn, depends on the mass difference between the parent and daughter states. This means that while the absolute values of the nuclear masses are not crucial in this context, the differences are essential. We adopt here the expression for the  $\alpha$ -decay half-life given in Ref. [19] for an even-even parent with the atomic number  $Z$

$$\log(T_{1/2}/s) = (aZ + b)(Q_\alpha/\text{MeV})(cZ + d) \quad (2)$$

with  $a=1.66175$ ,  $b=-8.5166$ ,  $c=-0.20228$  and  $d=-33.9069$ . It follows, for example, that a difference in  $Q_\alpha$  of the order of 1 MeV in a nucleus with  $Z = 118$  would make a difference in  $T_{1/2}$  of three orders of magnitude. Thus the calculation of  $Q_\alpha$  as close to reality as possible is vital for planning experiments. The neutron number dependence of  $Q_\alpha$  obtained in this work is illustrated in Fig. 5. It is



**Figure 5.** (Color online). Values of the  $\alpha$  particle separation energy,  $Q_\alpha$ , calculated in QMC $\pi$  for isotopes with  $100 < Z < 120$  in the region of neutron numbers  $138 < N < 252$ .



**Figure 6.** (Color online). The same as Fig. 5 but for isotones with  $148 < N < 190$  in the region of proton numbers  $100 < Z < 110$ . The  $N=152$ ,  $162$ ,  $172$  and  $184$  systems, corresponding to the suggested shell closures in different models, are highlighted by thicker lines in both panels.

beyond the scope of this work to compare the QMC $\pi$  predictions for  $Q_\alpha$  with all the models in the literature. We refer the reader to, for example, [14, 39–41] and references therein. The important result, which has not been observed in any of the other models, is that the (weak) effects of the  $N = 152$  and  $162$  shell closures disappear in nuclei with  $Z > 108$ , while the effects are enhanced for  $N \sim 180$ . We predict a smooth neutron number dependence of  $Q_\alpha$  for  $N < 200$  for all elements with  $Z$  up to 124, not showing any effects of shell structure. Some variations may be indicated for higher  $N$  but at this stage we can draw no systematic conclusion. To our knowledge such behaviour of  $Q_\alpha$  in this region has not been predicted before and a follow up would be very interesting.

Values of  $Q_\alpha$  in the region  $Z \leq 110$  are given in Fig. 6, showing the proton number dependence of  $Q_\alpha$  for  $148 < N < 190$ . We see a relatively smooth increase in  $Q_\alpha$  with increasing  $Z$  for  $N \leq 168$ , with a minor sensitivity to  $N = 152$  and  $N = 162$  (highlighted by thicker lines in the left panel of Fig. 6. This pattern prevails to  $N = 170$ – $174$  (top part of the right panel) before a notable development of

**Table 3.** Predicted properties of selected isotopes of Rf, Hs, Fl, Og, El-120 and El-124, as calculated in the QMC $\pi$ , FRDM and MM models. Experimental binding energies are taken from [27, 28]. All energies are in MeV.

Isotope	Z	N	$E_{bin}^{exp}$	$Q_{\alpha}$			$Q_{\alpha}$			$Q_{\alpha}$		
				$E_{bin}$	$\beta_2$	$Q_{\alpha}$	$E_{bin}$	$\beta_2$	$Q_{\alpha}$	$E_{bin}$	$\beta_2$	$Q_{\alpha}$
				QMC $\pi$			FRDM			MM		
<sup>252</sup> Rf	104	148		1860.20	0.245	9.4	1860.29	0.251	9.54	1859.29	0.245	9.85
<sup>268</sup> Sg	106	162	1963.37	1965.30	0.229	9.4	1965.39	0.232	7.59	1964.48	0.233	7.89
<sup>262</sup> Hs	108	154		1911.80	0.228	11.1	1912.32	0.243	10.88	1911.21	0.244	11.00
<sup>272</sup> Hs	108	164	1981.79	1983.80	0.236	9.8	1984.48	0.220	9.20	1982.98	0.225	9.80
<sup>272</sup> Ds	110	162	1973.36	1974.2	0.238	10.90	1975.56	0.221	10.04	1973.90	0.226	10.74
<sup>284</sup> Fl	114	170	2034.01	2035.40	0.129	11.6	2037.98	0.064	9.44	2033.85	0.149	11.53
<sup>296</sup> Og	118	178		2095.80	0.095	11.9	2099.12	-0.063	12.29	2094.67	0.039	12.06
<sup>296</sup> 120	120	176		2082.90	0.065	13.2	2085.70	0.075	13.69	2081.83	0.085	13.23
<sup>298</sup> 120	120	178		2097.30	0.063	13.0	2100.16	0.040	13.35	2095.48	0.054	13.44
<sup>306</sup> 124	124	182		2124.80	0.076	15.0	2126.39	0.000	13.43			

lower  $Q_{\alpha}$  values, varying systematically with Z (bottom part of the right panel).

The outcome of the QMC $\pi$  model indicates that there is a subtle interplay between proton and neutron degrees of freedom in developing regions of nuclei with increased  $\alpha$ -decay half life. As already discussed in the introduction and in the literature (e.g. [11]), it is likely that the sharp shell closures and shape changes observed in lighter nuclei, will instead be manifest as smoother patterns around the expected “shell closures”. These patterns have their origin in the competition between the Coulomb repulsion and surface tension of the large nuclear systems in which the single-particle structure is only one of the critical ingredients.

## 4 Future experiments

Following discussions during the Fusion2017 meeting [26, 38] we provide predictions for binding energies, quadrupole deformations and  $Q_{\alpha}$  values for selected nuclei summarized in Table 3. Corresponding results from the FRDM and MM models have been added for comparison. The reason for interest in these nuclei comes from their calculated shell structure in some traditional models, for example <sup>260</sup>Hs (N = 152, Z = 108 both supposed to be deformed shell closures). The presence of the N = 162 shell closure should manifest itself in <sup>268</sup>Sg and hence influence the  $Q_{\alpha}$  value of <sup>272</sup>Hs. The Z = 108 shell closure could be studied by examining the properties of <sup>272</sup>Ds. The existence of an  $\alpha$ -branch from the decay of <sup>272</sup>Ds is expected but has not yet been seen, the likelihood depending on the  $Q_{\alpha}$  value. The prospect of observing <sup>296</sup>Og, <sup>296,298</sup>120 and, possibly, <sup>306</sup>124, is considered to be realistic using available beams and targets.

Although all three models in Table 3 agree reasonably well in their predictions of binding energies and deformations, the differences in  $Q_{\alpha}$  are significant (see section 3.2). The most striking result is that QMC $\pi$  gives a much higher value of  $Q_{\alpha}$  than those yielded by the FRDM and MM models. For example, the calculated  $\alpha$ -decay half-life of <sup>268</sup>Sg in the QMC $\pi$ , FRDM and MM models are 0.21 s, 3.6 d and 5.9 h, respectively. One of the reasons for these variations is the presence of the N=162 shell closure in the

models. The higher  $Q_{\alpha}$  value is consistent with existence of the shell closure. This result also supports the suggestion of a higher  $Q_{\alpha}$  for <sup>272</sup>Hs. The calculated values of  $Q_{\alpha}$  for <sup>284</sup>Fl, <sup>296</sup>Og, <sup>296,298</sup>120 and <sup>306</sup>124 lead to very short  $T_{1/2}$ ,  $1.4 \times 10^{-5}$  s,  $4.0 \times 10^{-4}$  s,  $2.3 \times 10^{-6}$  s,  $5.8 \times 10^{-6}$  s and  $1.0 \times 10^{-8}$  s, respectively. More model predictions in this regions can be found in [30].

## 5 Summary and conclusions

We have presented the first results of the application of the QMC $\pi$  model to superheavy nuclei. Having in mind that the model is dependent on only four, well constrained variable parameters with a clear physical meaning, it is encouraging to observe that predictions for ground state binding energies, axially symmetrical shapes, regions of shell closure and  $Q_{\alpha}$  values are in good agreement with experimental data where available. The results presented in this work are of a very similar quality as the outcomes of other models, mean-field and macroscopic-microscopic which depend on many more parameters. Of course, the most interesting observation is that some of the predictions of the QMC $\pi$  model differ significantly from those in the traditional models and only experimental data can decide which model is closest to reality. Naturally there are some limitations in this work. As the energy surfaces as a function of quadrupole deformation may be complicated in the superheavy region (see e.g. [14]), it would be desirable to confirm the present results for nuclear shapes and the corresponding binding energies by repeating the minimization of the total binding energy with a constraint on the quadrupole moment. This procedure would help to avoid local minima and make a decision in cases of coexisting shapes with very similar binding energies. In addition, the predicted regions of shell closures, deduced in the present work from the neutron and proton number dependence of quadrupole deformation may be made more precise by examination of the spectrum of proton and neutron energy levels. These extensions of the current calculation are in progress.

## 6 Acknowledgement

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