Persistence of the Polarization in a Fusion Process

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Abstract. We propose an experiment to test the persistence of the polarization in a fusion process, using a petawatt laser hitting a polarized HD target. The polarized protons and deuterons heated in the plasma induced by the laser have a significant probability to fuse producing a $^3$He and a $\gamma$ ray or a neutron in the final state. The angular distribution of the radiated $\gamma$ rays and the change in the corresponding total cross section are related to the polarization persistence, but the resulting signal turns out to be weak. By comparison, the neutrons are produced hadronically with a larger cross section and are much easier to detect experimentally. A significant reduction of the cross section by parallel polarization of the deuterons is reliably predicted by the theory. Therefore, it is expected that the corresponding signal on the neutron counting rate could be seen experimentally.

1 Introduction

The polarization of $D$ and $T$ nuclei should increase their reactivity when used as fuel material in fusion processes induced either by magnetic or by inertial confinement.

The fusion reaction:

$$D + T \rightarrow \alpha + \text{neutron} + 17.6 \text{MeV} \quad (1)$$

goes mainly hrough the excitation of an $^3$He 3/2+ intermediate state, resulting from the coupling of the spins 1/2 and 1/2 of the $D$ and $T$ nuclei to a total spin $S = 3/2$. Without polarization of $D$ and $T$, the statistical distribution of the six possible states gives four $S = 3/2$ and two $S = 1/2$ states. Only the $3/2$ states can produce the intermediate 3/2 resonance. With 100% parallel polarization of $D$ and $T$, all states would contribute to the fusion, increasing the reactivity by 50%. In addition, the polarization allows to control the direction in which the reaction products are emitted, in particular the neutrons have a $\sin^2 \theta$ distribution. This can be very useful to reduce damages or activation of costly equipments [1]. The question is to know if the polarization will persist in a fusion plasma. We propose to investigate the polarization persistency using the reactions:

$$P + D \rightarrow ^3 H_e + \gamma + 5.5 \text{MeV} \quad (2)$$

$$D + D \rightarrow ^3 H_e + n + 3.267 \text{MeV} \quad (3)$$

induced by fusion of polarized protons and deuterons heated in a plasma. It is anticipated that the angular distribution of the radiated $\gamma$ rays or significant changes in the fusion rates can be related to the persistence of the polarization.

2 Magnetic versus Inertial confinement

The idea of inertial confinement is to compress tiny amounts of DT - simultaneously with heating - to such an extent that sufficient fuel burn is achieved within the time interval the fuel keeps together inertially. It turns out that the plasma density $n$ and confinement time $\tau$ required for inertial fusion are very different from those for magnetic fusion (11 orders of magnitude):

<table>
<thead>
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<th></th>
<th>Inertial</th>
<th>Magnetic</th>
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</thead>
<tbody>
<tr>
<td>confinement $n(\text{cm}^{-3})$</td>
<td>$10^{26}$</td>
<td>$10^{14}$</td>
</tr>
<tr>
<td>$\tau$ (sec)</td>
<td>$10^{-10}$</td>
<td>$10^{-15}$</td>
</tr>
<tr>
<td>$n.\tau$ (\text{sec/cm}^3)</td>
<td>$10^{15}$</td>
<td>$10^{16}$</td>
</tr>
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In both cases, however, the product $n.\tau$ has to satisfy the Lawson criterion \(n.\tau \geq 10^{15} \text{ sec/cm}^3\) which is set by the DT fusion physics.

In a Tokamak like ITER, the confinement time is expected to be as large as 300 s, which makes it very difficult for the polarization to survive till the end of the cycle, while at MEGAOULE, the whole compression time of a tiny target is of the order of 35 ns, making it much easier for the survival of the polarization. Kulsrud [1] has investigated several depolarization mechanisms as: 1) inhomogeneous static magnetic fields, 2) binary collisions, 3) magnetic fluctuations, 4) atomic effects, and concluded that all of them are weak. Relaxation times can become very long, when the depolarization paths are suppressed, as for example for HD [2]. However, in this matter, an experimental verification is always needed. In US, there is a project [3] to inject polarized $D$ and $^3$He in the DIII-D tokamak of San Diego, in order to see a 15% increase of the reaction rate of emitted protons by the fusion reaction:

$$D + ^3 H_e \rightarrow ^4 H_e + p + 18.35 \text{MeV} \quad (4)$$
However, the injection of 55% polarized D and $^3$He into a tokamak is a problem in itself, requiring technical innovations which may take some time.

3 Tentative set-up

At IPN Orsay, we have developed the static polarization of HD molecules for samples as large as 25 cm$^3$ [4]. It has been demonstrated, that the distillation and the ageing technique [5] allow to get nuclear relaxation times larger than one week, even at 1.5 K and 1 T. Proton polarization in excess of 60% and deuteron vector polarization higher than 14% have been achieved. It is advocated that a petawatt laser hitting a piece of polarized HD ice will induce locally a plasma hot enough to allow the fusion reactions (2) and (3) to take place and to be measured. If both $^1$H and $^2$D, namely the proton and the deuteron of the HD molecules are polarized in the same direction and have kept their polarization in the fusion process, the 5.5 Mev $\gamma$ ray will be emitted with some angular distribution relative to the polarization axis, also the fusion rates will depend drastically on the initial state polarizations. A tentative sketch of the experimental set-up is displayed in Fig. 1. It should be mentioned that with a power of 200 mJ/shot, the laser repetition rate can be adjusted to prevent melting of the target. Without cooling power provided by the holding cryostat, 1,000 such laser shots would be necessary to melt completely 25 cm$^3$ of solid HD. The overall polarization will decrease with time, but is continuously monitored by the NMR coils.

Back in 1970, a french group of the "Commissariat à l’Énergie Atomique" in France [6] reported the observation of neutron emission from DD fusion, after focussing a 3 GW fast laser on a piece of $^2$D$_2$ ice 1 mm$^2$ in cross section. At that time, a rise time of 5 ns was considered as fast. Since then Petawatt lasers have been developed using the chirped pulse amplification, able to deliver several tens of J within 20 fs to 1 ps. Those lasers can be used for fast ignition in inertial confinement fusion [7] or to accelerate particles [8].

Pretzler [9] reports quantitative data resulting from the irradiation of $^1$C$_2$D$_4$ targets with laser pulses (200 mJ, 160 fs, 4.5 $\mu$m FWHM, 790 nm, $\sim 10^{18}$ W/cm$^2$, 10 Hz). A total rate of 140 neutrons per shot could be produced, through the fusion reaction (3).

The technique to produce ion beams with ultra high power lasers starts to be well documented [10]: scaling laws and models exist which can predict the number and energy distribution of accelerated particles. This in turn, allows to tentatively optimize experimental conditions, although in this field, large uncertainties remain concerning the details of the processes. In particular, it is not at all sure that in a block of ice, ion acceleration will take place:
acceleration of particles have been reported only with thin targets.

From the quantitative data of Ref. [9], given the measured cross sections of reaction (2): $\sigma_0 (10\text{keV}) = 18\text{mb}$ [11]; $\sigma_0 (10\text{MeV}) = 1\text{mb}$ [12], $1\text{ - 10}$ (radiative captures / laser shot) could be expected. The detection of the corresponding $\gamma$ rays is a serious experimental problem. Conventional Ge detectors cannot be used, because of the large number of energetic electrons and $\gamma$ rays emitted in an extremely short time, which will pile-up in the Ge detector. Pair spectrometers would perform better, in spite of their lower efficiency.

**4 The “Few-Body” problems**

For the radiative capture (2), the experiment is essentially based on the angular distribution of the radiated $\gamma$ ray with respect to the polarization axis. Assuming that $P$ and $D$ nuclei collide from all directions in a hot plasma, with a total spin $S = 3/2$ (quartet transitions: $\sigma_4$, namely 100% polarization), while an unpolarized plasma involves also transitions from a total spin $S = 1/2$ (doublet transitions: $\sigma_2$), the angular dependence has the form [13]:

$$d\sigma_4/d\omega \sim (1 + \cos^2 \theta)$$

At the energies of interest, tens of keV or so, the process proceeds via S and P wave capture, and is induced predominantly by magnetic (for S-wave) and electric (for P-wave) dipole transitions. Higher multipoles, at the low energies considered here, can be neglected [14]. In addition, the p-wave contributions are much smaller than the S-wave one, since they involve the small D-wave component of the $^3H_2$ bound state. Therefore, although P-wave contributions involve isotropic and $(3 \times \cos^2 \theta - 1)$ terms, one would expect a very small distortion of the $(1 + \cos^2 \theta)$ angular distribution due to a pure S-wave magnetic dipole transition, as given by Eq. (5). From experimental point of view, this angular distribution means that 1/3 of the $\gamma$ rays due to quartet transitions will be preferentially emitted in the direction of polarization, namely in the direction of the laser beam. However, it is out of question to put a $\gamma$ ray detector in the direction of the laser beam, because of the large number of energetic electrons produced at forward angles. A transverse position, typically 90 degrees, is the most convenient. There the effect of the polarization on the quartet $\gamma$ rays counting rate is reduced to 25% to be compared to values approaching 100% at forward angles. Taking into account the highest achievable $P$ and $D$ polarization rates of respectively 80% and 30% in HD by the static polarization method and the dominant $\gamma$ ray contribution coming from doublet transitions for unpolarized nuclei; typically $\sigma_2/\sigma_{\text{unpol}} \approx 0.2$ from theoretical estimates [13] and even much smaller from experimental results at low energies [15], one cannot expect a signal larger than 3% on the counting rates between polarized and unpolarized targets. This makes the radiative capture experiment fairly difficult to exploit. It should be mentioned here, that the HD polarization technology allows to polarize $H$ and $D$ in an antiparallel configuration [16] in order to enhance the dominant $\sigma_2$ contribution. So doing, an increase $\sigma_{\text{pol}}/\sigma_{\text{unpol}} \approx 1.07$, namely 7% could be expected and eventually measured.

There is an alternative possibility offered by the hadronic fusion reaction (3) producing 2.45 MeV neutrons which are much easier to detect than $\gamma$ rays in a surrounding background. Based on a partial wave analysis [17], it has been argued that the cross section should be significantly reduced if the interacting deuterons have parallel vector polarizations (i.e. with total spin $S = 2$, namely quintet transitions: $\sigma_5$) [1], but resonating-group calculations [18] found that polarized fusions are not suppressed. However, it is known that resonating-group calculations are not very reliable for weakly bound nuclei as deuterons [19]. On the other hand, DWBA calculations give a large Quintet Suppression Factor (QSF): $\sigma_5/\sigma_{\text{unpol}} \approx 0.08$ in the range $E_\gamma = 20 - 150 \text{keV}$ [20]. Large reduction factors are confirmed by recent calculations [21], with QSF going from 0.5 at 100 KeV to 0.2 at 4 MeV. In addition total cross sections are in the range of 100 mb [22], to be compared to 100 mb for the electromagnetic reaction (2) [11]. In view of those considerations, the $D + D \rightarrow ^3H_3 + n$ is the way to go. It should be noted that for a polarized HD target, it is possible to increase the $D$ polarization above 50% at the expense of the $P$ one, by transfer of the $P$ polarization to $D$, using adiabatic fast passage [16]. A decrease of the emitted neutron counting rate of 10-20% going from a unpolarized target to a polarized one, namely $\sigma_{\text{pol}}/\sigma_{\text{unpol}} \approx 0.85$ should be easily measurable. The corresponding effect is further increased by the fact that the neutrons produced by quintet transitions are preferentially emitted perpendicular to the polarization axis [22].

**5 CONCLUSION**

A considerable effort in under way to produce energy using controlled fusion either by magnetic or by inertial confinement. Polarized fusion fuel is of great interest, both to increase the fuel reactivity and to control the direction in which the reaction products are emitted. The question is to know if the polarization will persist in a fusion process. We propose a possibility to investigate this point using high power laser beams on polarized HD samples through fusion reactions like: $P + D \rightarrow ^3H_3 + \gamma$ or $D + D \rightarrow ^3H_3 + n$. Before undertaking the corresponding experimental venture, precise predictions of the cross sections and polarization observables at low and moderate energies were needed, which is precisely in the scope of the present Few-Body Conference. Those predictions were kindly and rapidly performed by the relevant specialists after the FB19 meeting. It turns out that the radiative capture, which was initially considered to demonstrate the persistence of the polarization in a fusion process is not the prefered way to go, because the $\gamma$ rays, not only are difficult to select, but they are emitted preferentially along the polarization axis, in a region of high electromagnetic background. In addition, the low cross sections attached to an electromagnetic process, make it very difficult to pin down a signal smaller than 3% on the counting rates, although significant change in the total
cross section (7%) could be exploited in a different polarization scheme: P and D in an antiparallel configuration.

By comparison, the hadronic fusion seems much better, having a cross section larger by 3 orders of magnitude and producing a signal as large as 10-20% on the neutron counting rates, further increased by a favorable angular distribution of the neutrons emitted by quintet transitions. Neutron counters can be shielded and can work in a high background environment. Polarized target preparation is more difficult, requiring high deuteron polarization, but the relevant techniques are now well established.

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References

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