

Isolation of kinetic and spatial properties of uni-axial dynamic tensile loading of OFHC copper

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Abstract. Materials performance is recognized as being central to many emergent technologies. Future technologies will place increasing demands on materials performance with respect to extremes in stress, strain, temperature, and pressure. In this study, the dynamic ductile damage evolution of OFHC Cu is explored as a test bed to understand the role of spatial effects due to loading profile and defect density as well as the role of the kinetics of tensile pulse evolution. Well-characterized OFHC Cu samples of 30 μm , 60 μm , 100 μm , and 200 μm grain sizes were subjected to plate impact uniaxial strain loading in spall geometry to produce early stage (incipient damage). Using 2D metallographic techniques, soft recovered samples were studied to statistically link mesoscale processes to continuum level observations of free surface particle velocity measured with VISAR. Based on these findings, mechanisms for the void nucleation/growth and coalescence are proposed.

1 Introduction

Materials are recognized as being central to many emergent technologies. Future technologies will place increasing demands on materials performance with respect to extremes in stress, strain, temperature, and pressure. Hence, it is not surprising that material failure and property degradation is a critical problem. This type of damage and failure is often linked to defects within the material and determines service lifetime. Dynamic material damage develops through an evolution of microstructural changes from the atomic scale to the macroscale damage event and is affected significantly by defects such as grain boundaries and inclusions within structural metals which typically serve as failure initiation sites. While for some defect types, their mere presence is enough to cause failure, for others, spatial considerations such as distribution or density can be critical [1]. Shock compression via plate impact is used to investigate dynamic damage processes because of the ability to control global one-dimensional, planar strain loading conditions and to recover the sample for metallurgical analysis. In a plate impact experiment, dynamic tensile damage occurs when rarefaction (decompression) waves within the target material interact to produce tensile stresses in excess of the yielding threshold required for damage initiation.

Extensive work to date has established that spall is a complex process strongly influenced by the dynamic loading profile or shock wave shape imparted to the specimen and microstructure [2–5]. However, while this is now generally accepted, the evolution of the tensile stress leading to yielding in the material is not well understood. Most of the cited work has focused on the effects on particle velocity pull-back measurements in experiments in which the sample experienced complete failure. This precludes any significant conclusion about the individual contributions of material characteristics from these measurements. The evolution of the tensile stress leading to damage is controlled by shock rise time, pulse duration, peak shock stress (compressive and tensile) and release rate. In addition,

as the shock wave shape is altered experimentally, one important parameter that is frequently unaddressed is the coupling of that changing shock wave shape to the changes in the volume of material that is subjected to tension during a dynamic experiment. Figure 1 illustrates the effect of changes in shock wave shape (resulting from changes in impactor and target geometry) on the characteristics of the region of tension developed in the material leading to damage. Changes in pulse shape that significantly effect the volume of material being sampled in tension are believed to couple with microstructural length scales in a specimen to yield changes in the damage observed in soft recovered flyer plate experiments.

Incipient spall experiments are performed on high-purity copper samples with a known grain boundary density/distribution (grain size) to examine the relationship between these defect characteristics and void growth. Soft recovery of samples provides characterization of the end state, quantification of microstructural state and incipient damage.

2 Experimental procedure

2.1 Initial material characterization

All target materials were prepared from fully annealed 99.999% pure oxygen-free high-conductivity (OFHC) copper and have the same pedigree as the samples used in one of our earlier studies [6], in which the quasi-static compressive yield strength of copper revealed a very weak grain size dependency (<5 MPa). The samples have average grain sizes of 30, 60, 100 and 200 μm following annealing under vacuum at 450 °C for 30 min., 600 °C for 1 hour, 850 °C for 1 hour and 900 °C for 35 min, respectively. It should be noted that not only was the defect density/distribution (grain boundary) strictly controlled, but the defect nature (grain boundary type) was, as well. In all cases, the textures and misorientation distributions are

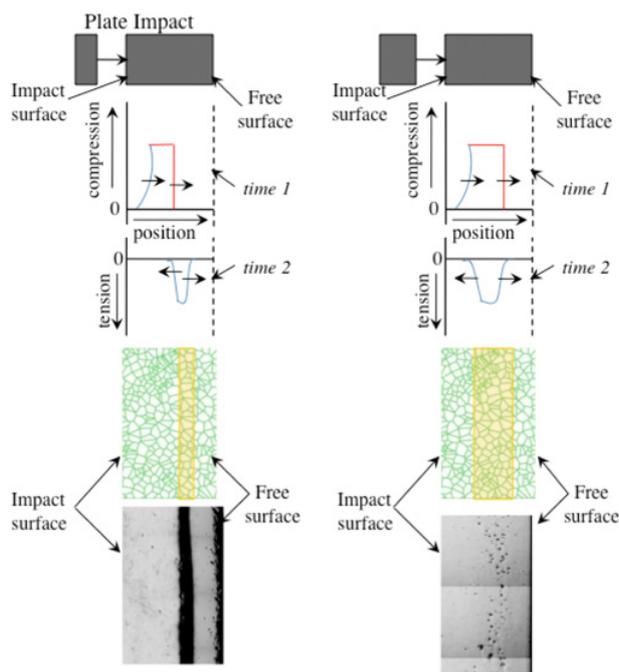


Fig. 1. Time evolution of the stress profile in a material during a plate impact experiment in spall geometry. The effect of changing the experimental geometry results in a tensile region of different size forming in the material. Microstructural considerations include number of defect sites subjected to the tension evolved in the volume defined by a given geometry.

consistent with those typically obtained in recrystallized copper.

Subsequent characterization of the damage in the recovered samples included optical and EBSD microscopy. In preparation for the optical and EBSD analyses, each recovered specimen is diametrically sectioned. They are subsequently mounted in an epoxy resin and prepared following standard metallographic techniques up to a $0.05\ \mu\text{m}$ colloidal silica final polish, performed on a vibratory polisher. Finally, they are electrochemically polished using a solution of two parts of phosphoric acid to one part water at $\sim 1.9\ \text{V}$ for $\sim 10\ \text{sec}$. A similar procedure is followed on the undeformed, as-annealed samples. Optical microscopy was performed on a Zeiss microscope equipped with an automated stage. Image processing is done using Image J. EBSD is performed on a Phillips XL30 FEG SEM using a voltage of $20\ \text{kV}$ and a spot size of 4.

2.2 Plate impact experiments

Plate impact experiments have been conducted using a smooth bore $78\ \text{mm}$ or a $50\ \text{mm}$ light gas gun. For the spatial effect study, a single well characterized loading condition is chosen for all experiments to ensure the kinetic aspects of the tensile stress profile evolution remain consistent for all experiments. The microstructural length scale is altered by changing the number of grain boundaries subjected to the tensile volume evolved by a given shock wave profile (controls defect density per unit volume). This is accomplished by using different grain sizes and identical

shock loading profile (identical experimental geometry). Quartz impactors (z-cut, $2\ \text{mm}$ nominal thickness) are mounted on brass sabots and launched using Ar gas. Quartz impactors ensure that a single, elastic shock wave is produced in the stress range of interest to this study. Table 1 lists the experimental details. To ensure that the recovered samples are subjected to only a known uniaxial loading/unloading history, significant radial release is minimized through the use of momentum trapping rings (technique described elsewhere [6]).

The kinetic study is accomplished by altering experimental geometry. When two rarefactions (the unloading portion of a shockwave) travelling in opposite directions collide inside the target material, a region of tension is formed. As rarefactions propagate through the material they decrease in slope (or the unloading rate decreases). These experiments utilize this property by interacting rarefactions that have propagated through material to achieve controlled rates at which tension is produced within the target material. This manipulation of geometry also alters the volume of material sampled in tension for each rate and requires us to control defect density per unit volume in similar fashion to what is shown in the microstructural length scale study.

After assembly, the parallelism, of the sample to the target plate, and the height of each pin with respect to the target plate is measured. Velocimetry profiles are obtained using a VISAR [7] system (spot size $\sim 0.5\text{--}1\ \text{mm}$ diameter) mounted 90° off of the shock direction axis to prevent the target from impacting the probe. In all cases shots are performed between two and four times to test repeatability as well as generate statistically accurate measurements of damage distributions.

3 Results and discussion

Figure 2 shows the free surface velocity (FSV) histories for the microstructural length scale experiments. Experimental parameters are listed in Table I along with calculated spall strength. The peak free surface velocities range from $82\text{--}84\ \text{m/s}$, corresponding to peak compressive stresses of $1.50\text{--}1.56\ \text{GPa}$. These values are calculated using the Mie Grüneisen equation of state for copper: $\rho_o = 8.924\ \text{gr/cm}^3$ (density), $C_o = 3.94\ \text{mm}/\mu\text{s}$ (bulk sound speed), $s = 1.489$ and $\gamma = 1.96$. The spall strength (σ_{spall}) is calculated using the relationship for a material that exhibits an elasto-plastic behavior [8]:

$$\sigma_{\text{spall}} \equiv \rho_o C_L \Delta FSV \left(1 + \frac{C_L}{C_o}\right)^{-1} \quad (1)$$

where $C_L = 4.77\ \text{mm}/\mu\text{s}$ is the longitudinal sound speed and ΔFSV is the difference in free surface velocity from the peak state to the minima after release. The measured and calculated experimental parameters are listed in Table 1. It is found that by holding the loading profile constant (i.e peak stress and experimental geometry identical for each experiment), the calculated spall strength is $\sim 1.31\text{--}1.38\ \text{GPa}$ with no clear influence of the grain size.

However, differences are observed in the pull back response after the release. The magnitude of the first pull

Table 1. Experimental parameters for OFHC Cu experiments and calculated values from velocimetry measurements.

Experiment	Grain size, μm	Impactor thickness, mm	U_D km/s	Target thickness, mm	Spall Strength, Gpa	V_{sp-pk}/V_{comp}
1s-1430	30	2.027	0.134	3.998	1.38	0.214
1s-1440	60	2.027	0.133	4.030	1.36	0.349
1s-1476	100	2.056	0.138	4.034	1.31	0.345
1s-1434	200	2.025	0.131	3.899	1.38	0.231
56-11-35	100	1.595	0.090	3.038	1.21	0.434
56-11-36	100	1.041	0.090	3.046	1.31	0.218
56-11-39	60	1.564	0.091	4.027	1.47	0.211
56-11-41	60	1.048	0.091	1.996	—	—
56-11-42	150	2.517	0.091	5.017	1.43	0.561

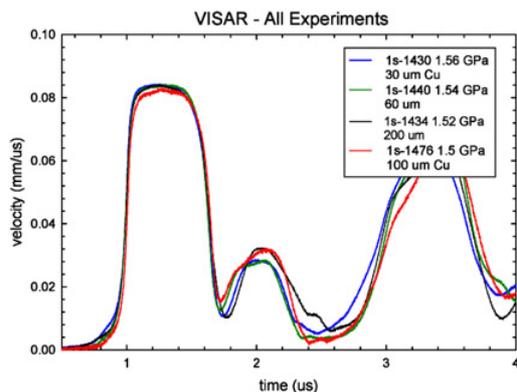


Fig. 2. VISAR results for the free surface velocity of OFHC Cu of 4 different grain sizes. While all shock wave shapes agree through the release, the response after the pullback shows differences in the yielding response and can be correlated to microstructural features in the recovered samples.

back and the rate at which the velocity rises to the first peak (i.e. spall peak) after the release minima are indicative of the mechanisms involved in the yielding process. Figure 3 shows the pull back signals for each experiment (from the minima reached after release to the spall peak velocity reached in the pull back). The times and velocities are shifted so that the minima coincides with the origin in the plot.

In a qualitative manner, Cochran and Banner [9] showed that the ratio of the peak particle velocity after the release minima (V_{sp-pk}) with respect to the peak particle velocity on compression (V_{comp}) is indicative of the observed void density in the recovered sample. Essentially, they observed that as the pull back signal increases in magnitude (as the difference between V_{comp} and V_{sp-pk} gets smaller) the amount of observed damage in also increased. Table II lists the ratio of these velocities for the current experiments. Figure 4 shows observed void area fraction from the recovered samples in these experiments plotted against grain size and it follows a trend similar to that expected by Cochran and Banner. The larger the difference between V_{comp} and V_{sp-pk} ; the larger the void area fraction observed in the recovered sample. The void area fraction in these experiments is calculated as the area of voids divided by the tensile area calculated using LASLO (a 1-D hydrocode developed at Sandia National Laboratory). In ductile materials, the rate at which the velocity

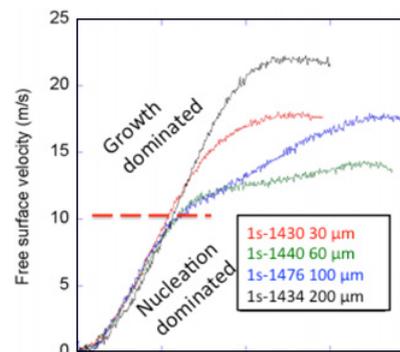


Fig. 3. VISAR pull back where all minima after the release have been normalized to 0. In all experiments the initial particle acceleration is similar. A mechanism change from nucleation dominated processes to growth dominated processes leads to a regime where microstructural length scale (grain size) controls growth mechanism (coalescence dominated results in high particle acceleration and individual void growth results in lower particle acceleration).

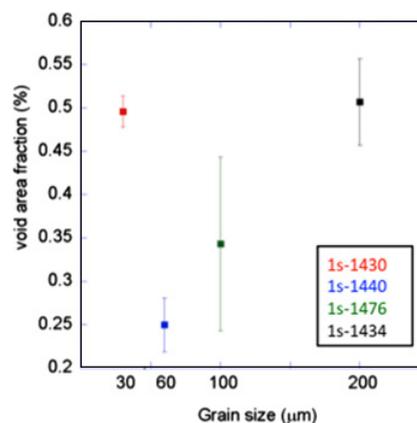


Fig. 4. Void area function measured from recovered samples indicates that identical loading for the cases of large grain size and small grain size the void field is larger or more developed, whereas for mid size grains the void field is smaller or less developed.

pull back rises to the spall peak after the minima is related to the void growth rate. The acceleration in the free surface velocity can be correlated with the velocity gradient of the wave reflected off the spall plane [8]. The velocity of this wave varies from C_L , in cases where there is a

Table 2. Experimental data for soft recovered OFHC Cu measured with 2D metallographic techniques.

Experiment	Grain size, μm	Void count	Void area fraction, %	Avg. void diameter, μm	Tensile area, mm^2	Tensile rate, Gpa/ μs
1s-1430	30	236	1.134	38.1	23.72	16.81
1s-1440	60	343	0.585	22.7	23.72	16.81
1s-1476	100	267	0.964	33.0	23.72	16.81
1s-1434	200	111	1.118	55.2	23.72	16.81
56-11-35	100	214	1.306	37.3	17.89	27.50
56-11-36	100	210	0.819	28.3	16.09	22.94
56-11-39	60	44	0.629	63.5	22.10	22.74
56-11-41	60	15	0.148	41.5	13.67	28.68
56-11-42	150	76	0.698	63.2	34.18	11.47

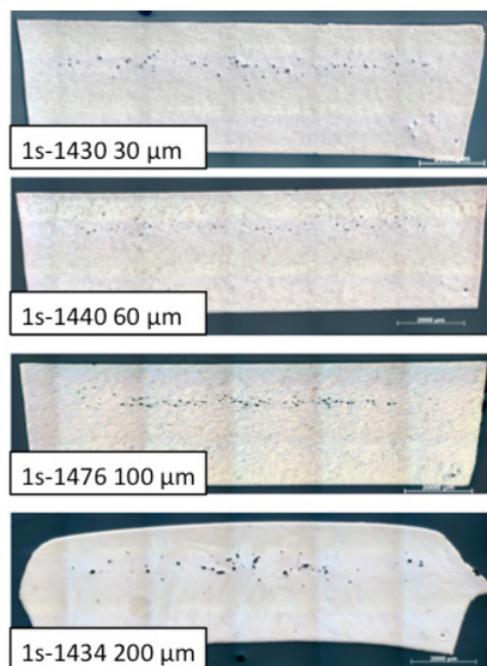


Fig. 5. Optical micrographs of the cross sections taken from the recovered Cu samples shows qualitatively the observed difference in damage field due only to differences in grain size under identical loading conditions.

rapid fracture process; to C_0 in the cases where there is a stress relaxation due to a resistance in void growth which leads to lower damage rates [10]. Thus, shorter rise times to the spall peak correlate with a more rapid completion of the damage. In the microstructural length scale experiments, the 30 and 200 μm samples exhibit a monotonic rise from the minima to the spall peak, indicative of a fast damage evolution process. Conversely, 60 and 100 μm samples experience a fast rise time initially, which subsequently displays a distinctive change in slope. This response indicates the dominate damage evolution mechanism undergoes a change. Figure 5 shows the optical micrographs of cross sections taken from the recovered samples. Figure 6 shows EBSD images from those same cross sections. Table 2 lists the void area fraction which is in agreement with the ratio of ratio of particle velocity for the compressive peak (V_{comp}) to the spall peak particle velocity after the minima ($V_{\text{sp-pk}}$). The statistics relating to the character of the damage indicates the dominate mech-

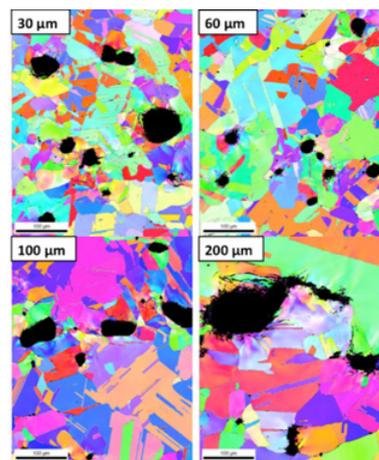


Fig. 6. EBSD images of the cross sections of the Cu show that for small and large grain sizes larger coalesced voids are observed. In the cases of the mid grain sizes voids are largely isolated no significant coalescence is observed.

anisms. For intermediate grain sizes of 60 and 100 μm , measurements reveal a larger count of primarily smaller sized voids which indicates a largely nucleation dominated damage growth mechanism with growth of individual isolated void (minimal coalescence). However, in the largest and smallest grain sizes of 30 and 200 μm , measurements reveal a smaller void count of primarily larger sized voids which is indicative of more rapid void coalescence and growth (minimal nucleation).

In the tensile pulse rate experiments, sample geometry was manipulated to alter the rate at which tension developed in 5 different experiments. For each case, defect density per unit volume was held constant. Figure 7 shows the VISAR results for free surface velocity for each experiments as well as a notional microstructural cross section indicating the region where tension would be formed for each wave shape. Note that experiment 56-11-41 shows only full thickness ringing in the particle velocity which indicates no spall plane was formed. Figure 8 shows the normalized spall peak after the minima for the 4 experiments displaying damage in the VISAR trace. Figure 9 shows the optical cross sections of the samples with EBSD misorientation maps. Statistics from the EBSD (shown in Table 2) show that the relationship between compressive peak particle velocity and the peak particle velocity of

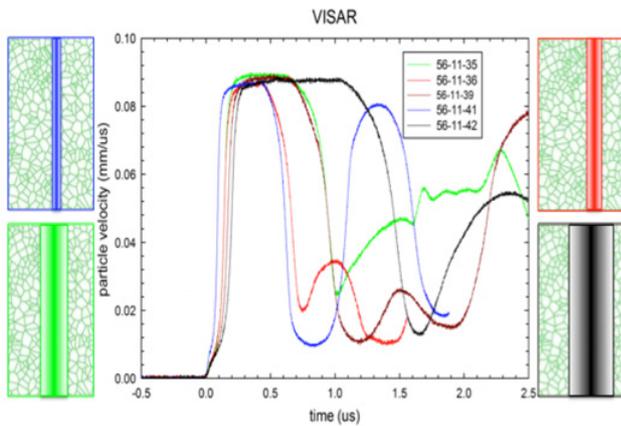


Fig. 7. VISAR free surface velocity measurements for tensile pulse rate experiments.

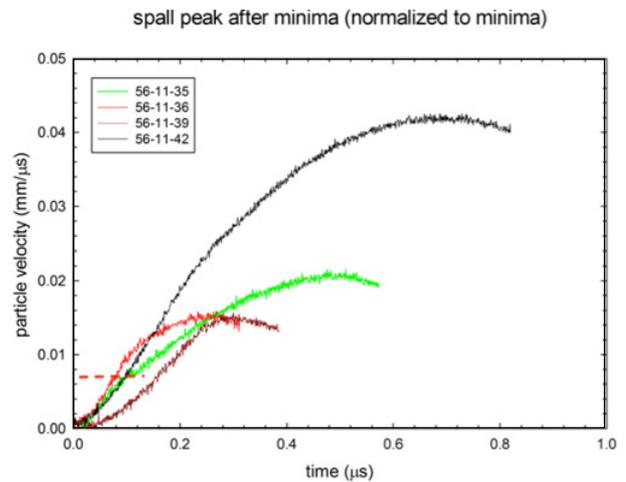


Fig. 9. Spall peak after the release minima is indicative of growth mechanism (rate of growth) as well as total damage content. The dotted red line indicates a possible nucleation dominated regime for all experiments except 56-11-39, which appears to be purely growth dominated.

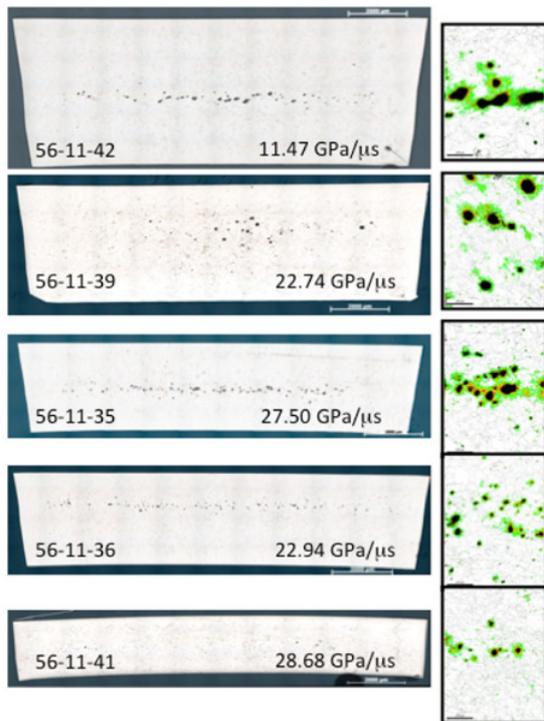


Fig. 8. Optical microscopy images of tensile pulse rate experiments with EBSD misorientation maps indicating retained plastic deformation in the microstructure due to yielding processes.

the spall peak after the minima (V_{comp}/V_{sp-pk}) does not go with void area fraction as it did in the microstructural length scale study. It is also observed that given different loading paths for tensile pulse evolution the transition from nucleation to growth dominated regimes as reflected in the spall peak after the minima may seem obvious for all experiments except 56-11-39 where the nucleation region is not obvious. As tensile rate increases (as seen in Fig. 9) a transition is observed from a small number of large voids (growth dominated) to a large number of small voids (nucleation dominated) to a rate at which both nucleation and growth appear to be overdriven (experiment 56-11-41, the small number of voids was not sufficient enough for

VISAR to detect a significant wave reflection and thus does not detect any damage in this experiment).

4 Conclusions

In the microstructural length scale experiments the samples are subjected to identical loading conditions (geometry, peak pressure, etc.). Drawing conclusions from continuum level measurements such as free surface velocity and correlating them to the mesoscale microstructural processes is accomplished through careful model driven experimental design and diagnostics coupled with metallographic analysis of recovered test materials. In particular, it has been found that in spite of identical loading conditions the calculated spall strength remains comparable in all experiments with no dependence on grain size while the character of the resultant damage shows significant differences. The free surface velocity acceleration after the pull-back minima correlates with void growth mechanism, in which higher rates (observed in the 30 μm and 200 μm experiments) rendered larger voids indicative of a coalescence dominated growth mechanism rather than growth of individual voids. Interestingly, the samples that show a change in slope of the free surface velocity after the minima, present higher average kernel misorientations (60 μm and 100 μm experiments). This is indicative of the competition between void growth and plastic dissipation processes that is the result of individual void growth in the absence of significant coalescence. In all experiments, nucleation and growth are likely occurring simultaneously. However, early in time during the yielding process the dominant mechanism is nucleation of voids until a change in mechanism leads to a growth dominated regime. During the growth dominated regime, microstructural length scale dictates the type of growth mechanism (coalescence or individual void growth) that results in the rate of total damage evolution observed in the particle velocity.

In the tensile pulse rate experiments, relationships between statistical observations of the resultant damage and the particle velocity begin to break down. This may be indicative of mechanisms active during compressive pulse duration having a significant effect on yielding processes during tension. However, metallographic analysis indicates that tensile pulse rate significantly influences the mechanisms dominating the damage evolution. For very low strain rates a growth dominated mechanism is present with very few voids which have significantly grown and coalesced. This is also observed in the monotonic rise to the spall peak after the minima. As tensile rate increases a transition to nucleation dominated mechanisms is observed where there are a large number of smaller voids present and the rise to the spall peak after the minima display two clear slopes. Interestingly, the highest rate experiments appear to overdrive both nucleation and growth mechanisms, resulting in almost no resultant voids. Further exploration into the mechanisms dominant in these regimes is ongoing.

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