

Laser-plasma booster for ion post acceleration

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Abstract. A remarkable ion energy increase is demonstrated for post acceleration by a laser-plasma booster. An intense short-pulse laser generates a strong current by high-energy electrons accelerated, when this intense short-pulse laser illuminates a plasma target. The strong electric current creates a strong magnetic field along the high-energy electron current in plasma. During the increase phase in the magnetic field, a longitudinal inductive electric field is induced for the forward ion acceleration by the Faraday law. Our 2.5-dimensional particle-in-cell simulations demonstrate a remarkable increase in ion energy by several tens of MeV.

1. INTRODUCTION

By chirped pulse amplification, high laser intensities have been realized, and high intensity short pulse lasers are now available for experiments and applications. On the other hand, ion beams are useful for basic particle physics, medical ion therapy, controlled nuclear fusion, high-energy sources, and so on. The energy of ions, which are accelerated in an interaction between an intense laser pulse and a near-critical density target, are over a few tens of MeV [1–6]. The issues in laser ion acceleration include ion beam collimation, ion energy spectrum control, ion production efficiency, etc. [1–28]. Depending on ion beam applications, the ion particle energy should be controlled. For example, ion beam cancer therapy needs 100 ~ 150 MeV for proton energy. Therefore, in this paper we focus on a boost of ion beam energy by post-acceleration in laser plasma interaction. In this paper we propose a laser-plasma booster as an ion post acceleration scheme. In our study, we employ an intense short-pulse laser and a near-critical density plasma target, which consists of hydrogen. Figures 1(a) and (b) show the conceptual diagram of the laser-plasma booster. Figure 1(a) presents just a possible example scheme for a future multi-stage laser accelerator. In this paper we focus on one post acceleration stage. In this paper, we perform 2.5-dimensional particle-in-cell simulations to investigate the ion beam post-acceleration.

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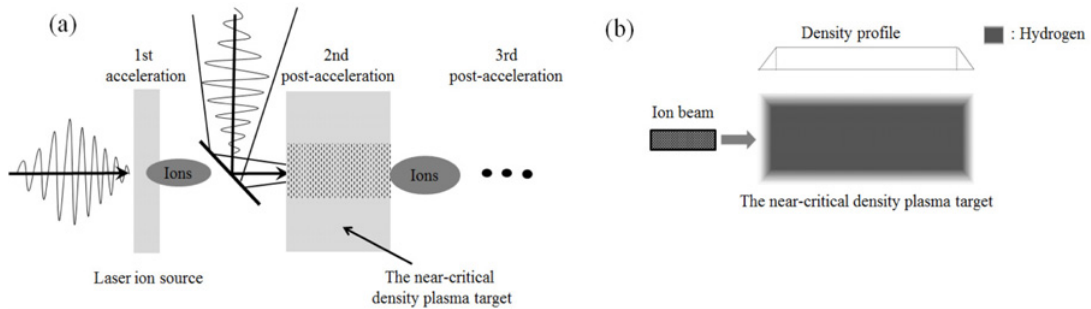


Figure 1. (a) An example concept for future ion post acceleration scheme and (b) the near-critical density hydrogen plasma target. The near-critical density plasma booster target provides a stable inductive acceleration field, for post acceleration of laser produced ions.

2. LASER-PLASMA BOOSTER FOR ION POST ACCELERATION

2.1 The near-critical density plasma target simulation

We perform 2.5-dimensional particle-in-cell simulations. The target model is shown in Fig. 1(b). The near-critical density plasma target is located in $9.5\lambda < x < 49.5\lambda$ and $19.0\lambda < y < 31.0\lambda$. The plasma target has a flat-top density profile. The edge region has a linear density gradient from $0n_c$ to the maximum density of $0.5n_c$, and has a linear density gradient in 2λ in the X and Y directions at the target edges. The laser intensity is $I = 1.0 \times 10^{20} \text{ W/cm}^2$, the laser spot diameter is 4.0λ , and the pulse duration is 40 fs. The laser transverse profile is Gaussian, and the laser temporal profile is also Gaussian. The laser wavelength is $\lambda = 1.053 \mu\text{m}$. The simulation box is 80λ in the longitudinal direction and 30λ in the transverse direction. The near-critical density plasma target density is $0.5n_c$ and the ion beam density is 10^{16} cm^{-3} . The pre-accelerated proton beam has the initial mean energy of 110 MeV, has the initial proton beam temperature of 2 MeV and is located at the left of the plasma target initially. The initial proton beam size is 2λ in X and 5λ in Y . The size and the initial temperature depend on the pre-acceleration mechanism. The size of the pre-accelerated ion beam and its mean energy are selected here to demonstrate the post-acceleration mechanism by the inductive acceleration.

2.2 Post-acceleration of the ion beam

Figures 2 show the distributions of acceleration electric field E_x in $\text{MV}/\mu\text{m}$ at (a) $t = 110 \text{ fs}$ and (b) $t = 230 \text{ fs}$ and the distributions of magnetic field B_z in Tesla at (c) $t = 110 \text{ fs}$ and (d) $t = 230 \text{ fs}$. The laser generates the high-energy electrons inside of the target. A magnetic field is also formed along the channel in the laser plasma interaction [26–28]. When the intense laser pulse propagates through the plasma, it accelerates a part of electrons. The electrons form a high current and generate the magnetic field. In the laser plasma interaction, the ion dynamics is affected directly by the electric field and the behavior of the electrons. The electrons form a strong magnetic field, and during the increase in the azimuthal magnetic field a strong inductive electric field is generated [6]. The ions are accelerated by the inductive electric field. The inductive acceleration field moves with a speed less than c , depending on the plasma density. At this target density of $0.5n_c$, the speed of the inductive electric field is about $0.66c$. Therefore, the inductive acceleration field is appropriate for ion post acceleration and is rather stable. Figure 3 shows the ion beam energy distribution at $t = 50 \text{ fs}$, 230 fs and 450 fs for the near-critical density plasma target. At $t = 50 \text{ fs}$, Fig. 3 shows the initial energy of the ion beam which is not yet accelerated. The maximum ion beam energy reaches 171.2 MeV from the initial energy of 110 MeV. The energy conversion efficiency from laser to the pre-accelerated protons is $4.67 \times 10^{-2}\%$, and the

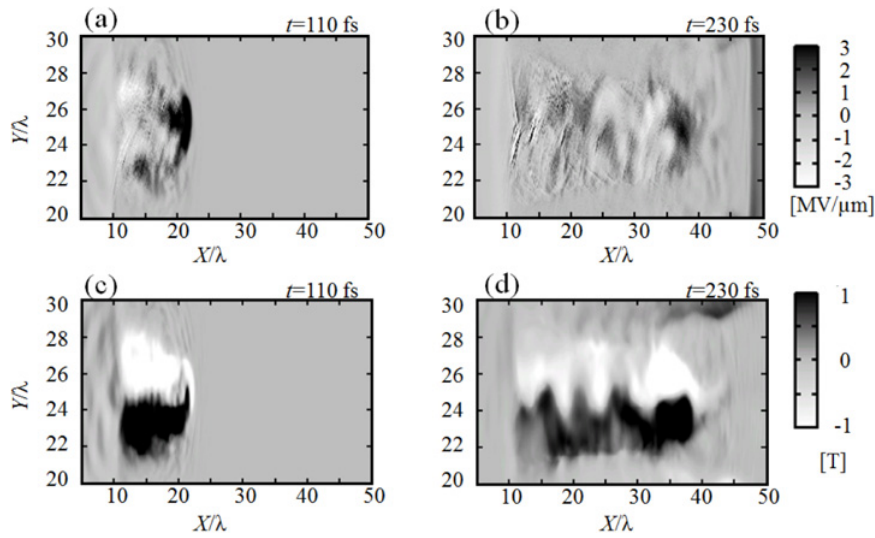


Figure 2. The distributions of acceleration electric field E_x in $MV/\mu m$ at (a) $t = 110$ fs and (b) $t = 230$ fs, and the distributions of magnetic field B_z in Tesla at (c) $t = 110$ fs and (d) $t = 230$ fs.

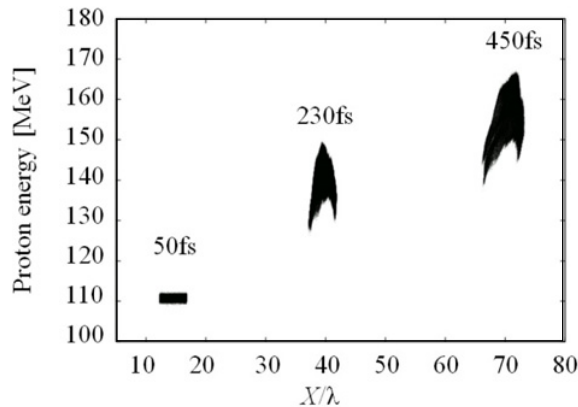


Figure 3. The ion beam energy distributions at $t = 50$ fs, 230 fs and 450 fs for the near-critical density plasma target. At $t = 50$ fs, the ion beam is not yet accelerated.

energy efficiency to the background target protons is 35.5%. The energy coupling efficiency to the pre-accelerated proton beam is rather low. However, the present result shows a possible post-acceleration mechanism with successive inductive accelerations. Figure 3 shows that the final energy spread of protons becomes large compared with the initial energy spread. The final energy spread influences the beam quality or particle selection device parameters depending on ion beam utilization purposes. The requirement of the final beam quality should be addressed in the future, when the application purposes are specified.

3. CONCLUSIONS

In this paper we have proposed a laser-plasma booster as an ion post-acceleration scheme. We succeeded to increase the energy of the ion beam by the inductive acceleration field in a near-critical density

plasma target illuminated by an intense short-pulse laser. The post-accelerated ion maximum energy increases from 110 MeV to 171.2 MeV by the strong inductive acceleration electric field. The work in this paper presents an important method for ion energy control in laser plasma acceleration. For practical applications of the laser ion accelerator the issues include the post-acceleration of the ion beam as studied in this paper, the ion beam quality improvement in the energy spectrum control including a mono-energetic ion beam generation, neutralized or unneutralized ion beam transportation for a long distance, etc. This study presents a new way to generate and to control the ion beam energy.

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