

Giant Dipole Resonance in $A \sim 144$ mass region

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Abstract. Exclusive measurement of giant dipole resonance (GDR) γ rays has been performed in ^{144}Sm nucleus which was populated at near barrier energy using the heavy ion reaction of ^{28}Si beam on ^{116}Cd target. GDR γ rays were detected in coincidence with low energy γ rays using 4π sum-spin spectrometer. The ^{144}Sm nucleus was populated at an excitation energy of 68 MeV in the temperature range of 1.1-1.3 MeV. The measured GDR widths in this temperature range are consistent with the Kusnezov's parametrization.

1 Introduction

The phenomenon of Giant Dipole resonance (GDR) built on excited states [1, 2], offers unique possibility to study the fundamental properties of hot rotating nuclei. GDR characteristics, viz. centroid, width, strength etc. are found to be dependent on the temperature (T) and angular momentum (J). Using high efficiency detector array setup, a few experiments [3, 4] have been carried exclusively to study the dependence of temperature and angular momentum on GDR observables. Heavy ion fusion reactions impart a large amount of energy and angular momentum to the internal degrees of freedom in hot rotating nuclei. The quadrupole deformation experienced by the compound nuclear system can be inferred from the splitting of GDR strength function. Though several theoretical models have been proposed to explain the effect of angular momentum on GDR widths, a parametrization proposed by Kusnezov et al. [5] has been successfully used in fitting the evolution of GDR widths in 1-2 MeV temperature range in different mass regions. Bhattacharya et al. [6] have used this phenomenological function to scale experimental widths in recent experiments spanning different mass regions. In these measurements, it is necessary to have high resolution of angular momentum selection to decouple the effect of J and T mixing at a given excitation energy. The multiplicity filter with high efficiency and granularity plays a crucial role in such experiments. Increase in solid angle coverage leads to better selection of J windows. Using the 4π spin-spectrometer and high energy γ spectrometer at Inter University Accelerator Centre (IUAC), we have studied the GDR characteristics for ^{144}Sm [7] at wide range of T and J for which the ground state systematics are already known [8]. In this work, we aim to extract GDR widths

from ^{144}Sm nucleus by populating it at near barrier energies. Thus, we could map the evolution of GDR widths as a function of angular momentum as well as temperature. In this paper, we present initial observations of the effect of angular momentum on GDR width in ^{144}Sm at near barrier energy and low temperatures.

2 Experiment details

The experiment was performed at the 15UD Pelletron accelerator facility at IUAC, New Delhi. The accelerated DC beam of ^{28}Si at 125 MeV energy was bombarded on ^{116}Cd target to form the compound nucleus (CN) ^{144}Sm . The target was prepared at target laboratory of IUAC by rolling to a thickness of ~ 1.8 mg/cm². The targets were self supporting with more than 98% enrichment. The ^{144}Sm nuclei were populated at an excitation energy of 68 MeV. The high energy γ rays were detected in High Energy Gamma Ray Spectrometer (HiGRaSp) [9] which consists of a single crystal of NaI(Tl) in the centre surrounded by four plastic scintillator detectors. Plastic detectors were used as active shield for cosmic background reduction. The NaI(Tl) has a dimension of 25.4 cm \times 30.5 cm and is coupled with seven Photo Multiplier Tubes (PMT) operating at a voltage of +800V. All the PMTs were gain matched and the dynode signals were fed to spectroscopic amplifiers. The summed output from these amplifiers was fed to Analog to Digital Converter (ADC) as energy signal. Anode signals from all the PMTs were amplified using Fast amplifiers and fed to Constant Fraction Discriminators (CFD) providing fast timing signals for starting various Time to Amplitude Converters (TAC).

HiGRaSp was positioned at a distance of 80 cm from the target centre at an angle of 90° with respect to the beam

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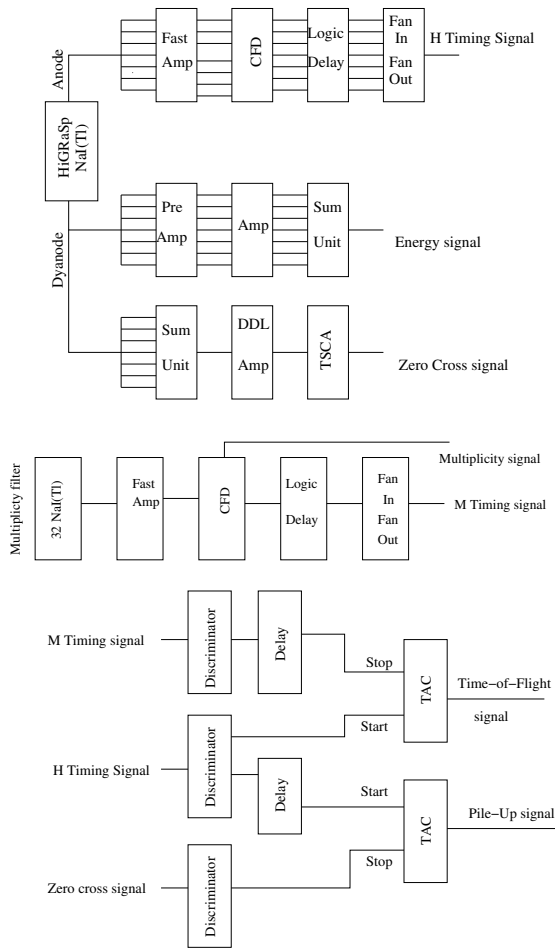


Figure 1. Block diagram of electronics circuit used in the experiment.

direction. Laboratory γ ray sources were used for calibration and extraction of experimental response function. The response function of detector for mono-energetic γ rays was generated using monte-carlo simulation package GEANT4 [10]. The dynamic range of HiGRaSp was set at 34 MeV. Energy gain was monitored in-beam using the neutron capture peak of 6.83 MeV.

Low energy γ rays from the fusion reaction were detected using the 32 element 4π sum-spin spectrometer [11]. In the present experiment, 27 detectors were used covering nearly 82% of 4π solid angle. An energy threshold of 100 keV was set for each detector. A Time-of-Flight (TOF), set between HiGRaSp and sum-spin spectrometer, separated neutron and γ events. To filter piled up events due to high count rate, a Pile-Up (PU) circuit was used implementing the zero cross-over technique. The block diagram of electronics used for processing signals from HiGRaSp and multiplicity filter is shown in fig. 1. The online data was recorded and analysed using the software package CANDLE [12], developed at IUAC. The raw 2D spectra of Energy vs TOF and Energy vs PU are shown in figures 2 and 3, respectively.

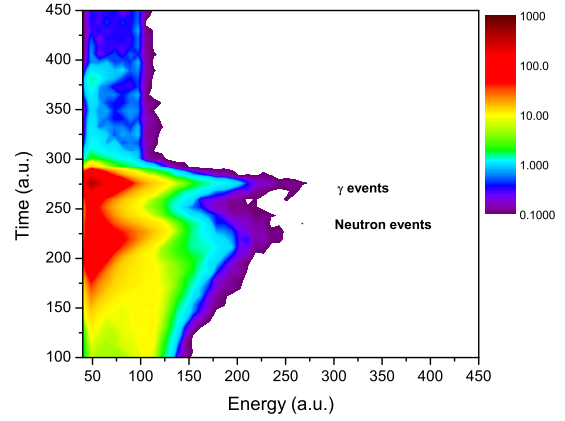


Figure 2. A time of flight (TOF) 2D spectrum showing prompt γ events and other delayed events for 125 MeV beam energy. The spectrum is gated with the condition of fold 6 and above.

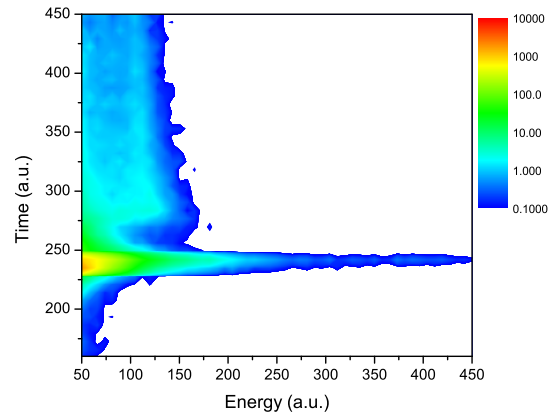


Figure 3. A Pile-Up 2D spectrum showing piled up events and good events. A tight gate was used on good events for cleaning the high energy γ spectra.

3 Data Reduction & Analysis

Various 1D histograms were generated offline using CANDLE. The fold distribution at beam energy of 125 MeV is shown in fig. 4. Average angular momentum for different fold bins were calculated by extracting multiplicity distributions for those fold distributions. This was accomplished in the following way. The response function of multiplicity was generated using the iterative algorithm prescribed in Ref. [13]. According to this method, the probability of M number of γ rays hitting F number of detectors out of total N number of detectors is given by

$$S(F, M) = aS(F, M-1) + bS(F-1, M-1) + cS(F-2, M-1)$$

$$a = 1 - (N - F)\omega(1 + \epsilon F/N - 1)$$

$$b = (N - F + 1)\omega(1 - \epsilon(N - 2F + 1)/(N - 1))$$

$$c = (N - F + 2)\omega\epsilon((N - F + 1)/(N - 1))$$

where ω is the cross talk probability and ϵ is the efficiency of each detector. Simulated fold distributions were generated using following equation.

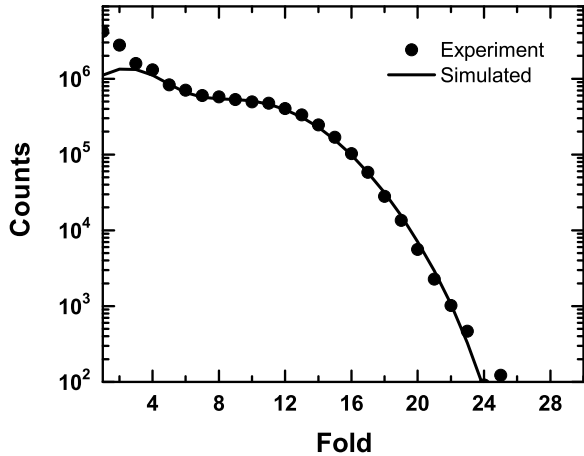


Figure 4. Experimental fold distribution at 125 MeV beam energy data. Solid line is the best fit of simulated fold distribution.

$$P(F) = \sum_{M=0}^{M_{max}} S(F, M)P(M)$$

The form of multiplicity distribution used in this method was assumed to be a modified fermi function and was given by

$$P(M) = \frac{2M + 1}{1 + \exp\left(\frac{M - M_0}{\delta M}\right)}$$

where M_0 and δM are free parameters. The simulated fold distribution was compared with the experimental fold distribution by varying the M_0 and δM and minimizing the χ^2 . For 125 MeV beam data, $M_0 = 17$ and $\delta M = 3.0$ were obtained. The experimental and simulated fold distributions are shown in fig. 4. The multiplicity distributions corresponding to different fold bins were generated as $P(M) = S(F, M)P(F)$. Finally, the average of these distributions can be related to average angular momentum as $\langle J \rangle = 2 \langle M \rangle + C$, with $C = 4$.

High energy γ ray spectra were extracted offline by using γ gate in TOF 2D spectra (fig. 2) and PU gates in PU-2D spectra (fig. 3). The resulting spectra were gated with different fold bins. These spectra were also cleaned by plastic detector signal in anti-coincidence for cosmic rays rejection. The final resulting spectra were Doppler corrected and were binned in energy bin of 1 MeV.

A modified version of statistical model code CASCADE [14] incorporating two component lorentzian function for GDR was used for the statistical model calculations. Asymptotic level density parameter was fixed at the value of $\bar{a} = A/8.5 \text{ MeV}^{-1}$. The code has an option of six parameter fitting for two component lorentzian distribution of GDR, viz. centroid energy (E_1, E_2), strength (S_1, S_2) and width (Γ_1, Γ_2). The CASCADE output, after varying GDR parameters, was folded with detector response and compared with the experimental spectra. Fit parameters were obtained by minimizing the χ^2 in the

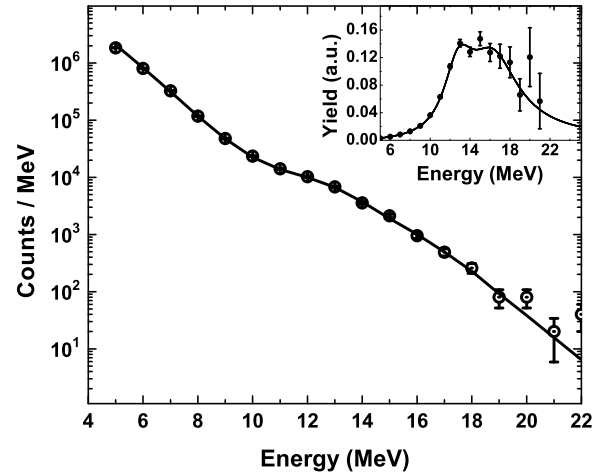


Figure 5. A cleaned high energy γ ray spectra with a condition of fold 11-12 at beam energy of 125 MeV. Solid line is the CASCADE fit incorporating the two component GDR strength function. GDR strength in linear scale is shown in inset. Linearized plot is generated as $Y_{exp}/Y_{th} \times F_{2L}$. See text for details.

energy range of 12-20 MeV. From these parameters, linearized GDR strength functions [15] were generated as $Y_{exp}/Y_{th} \times F_{2L}$, where Y_{exp} is the experimental yield, Y_{th} is the CASCADE calculation folded with detector response and F_{2L} is the two component lorentzian function with each component given by

$$F_L(E_\gamma, E_C, \Gamma) = \frac{\Gamma^2 E_\gamma^2}{(E_\gamma^2 - E_C^2) + \Gamma^2 E_\gamma^2}$$

A fold gated high energy γ ray spectrum fitted with CASCADE output is shown in fig. 5. The GDR widths were extracted as FWHM of the GDR strength function as depicted in inset of fig. 5.

4 Results

The experimental GDR widths ($\Gamma(T, J)$) at temperature (T) and angular momentum (J) have been extracted for ^{144}Sm and presented as a function of J at the average temperature bins of 1.1 and 1.3 MeV. The experimental width has been compared with calculated GDR widths at given T using the phenomenological function proposed by Kusnezov et al. which is based on thermal shape fluctuation model (TSFM). The GDR width at a finite T and J in liquid drop regime can be calculated as

$$\Gamma(T, J) = \Gamma(T, 0) \left(1 + \frac{1.8}{1 + \exp\left(\frac{1.3 - \xi}{0.2}\right)} \right)^{\frac{4T_0}{T+3T_0}}$$

where

$$\Gamma(T, 0) = \left(6.45 - \frac{A}{100} \right) \ln \left(1 + \frac{T}{T_0} \right) + \Gamma_0(A)$$

Here $\Gamma_0(A)$ is the width for the spherical shape of the nucleus and the value of T_0 used is 1 MeV. ξ is the reduced angular momentum defined as $\xi = J/A^{5/6}$, which is

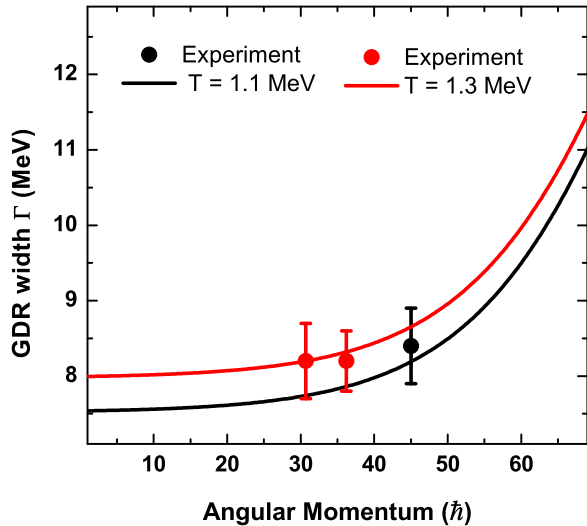


Figure 6. GDR width as a function of average angular momentum. Circles represent experimental data, and Solid lines are theoretical calculations for different T values using the kusnezov's parametrization.

proportional to the rotational energy of the system. In these calculations, Γ_0 was adjusted to match the experimental width at low angular momentum and was fixed to be 3.8 MeV. The experimental GDR widths are consistent with the widths calculated using this parametrization, in the given temperature range as shown in fig. 6. The analysis of data at higher beam energies for the same reaction channel is under process.

5 Summary

Angular momentum gated GDR widths were studied in ^{144}Sm nucleus using heavy ion reaction of $^{28}\text{Si}+^{116}\text{Cd}$ at beam energy of 125 MeV. The compound nucleus was populated at near barrier energy to get the lowest temperatures. The experimentally extracted width is compared with theoretically calculated width using Kusnezov's parametrization. The theoretical predictions are con-

sistent with experimental results for this nucleus. More experiments are planned to scan wider range of temperature and angular momentum for this nucleus.

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