

Search for rare processes with ZnWO₄ crystal scintillators

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Abstract. Radiopure ZnWO₄ crystal scintillators with mass (0.1–0.7 kg) have been developed and put in measurement in the Gran Sasso National Laboratories of the INFN to search for rare processes. The radioactive contamination of the crystals have been estimated and the double beta decay of zinc and tungsten isotopes was searched for, reaching a sensitivity at the level of 10¹⁸ – 10²¹ yr; in addition a new half-life limit on alpha transition of ¹⁸³W to the metastable excited level of ¹⁷⁹Hf has also been obtained. The achieved radiopurity of the ZnWO₄ crystals make them very promising detectors for $\beta\beta$ decay investigations while their anisotropic features make them very interesting detectors to investigate dark matter particles directionality.

1 ZnWO₄ crystal scintillators

In recent years ZnWO₄ crystal scintillators were considered in the search for rare processes [1–5]. In particular four clear, slightly colored ZnWO₄ crystal scintillators have been produced: i) crystals ZWO-1 (117 g) and ZWO-2 (699 g) were produced in the Institute for Scintillation Materials (ISMA, Kharkiv, Ukraine) from crystal ingots grown in platinum crucibles by the Czochralski method; ii) the crystal ZWO-3 (141 g) was obtained by recrystallization from the sample ZWO-2 at the ISMA; iii) the ZWO-4 (239 g) was produced in the Nikolaev Institute of Inorganic Chemistry (Novosibirsk, Russia) by the low-thermal gradient Czochralski technique also in platinum crucible. The radioactive contamination of these ZnWO₄ crystal scintillators has been investigated deep underground in the Gran Sasso laboratory, by using the low background facility DAMA/R&D [6]. In the measurement the considered ZnWO₄ crystals were fixed inside a cavity of $\varnothing 49 \times 59$ mm in the central part of a polystyrene light-guide 66 mm in diameter and 312 mm in length. The cavity was filled up with high purity silicone oil. The light-guide was optically connected, on the opposite sides, to two low radioactive EMI9265B53/FL 3 inches photomultipliers (PMT). The light-guide was wrapped by PTFE reflection tape. Some upgradings of the detector system have been performed in the different periods of measurements; in particular, in the final stages of the experiment two polished quartz light-guides ($\varnothing 66 \times 100$ mm) were installed between the polystyrene light-guide and the PMTs to suppress γ ray background from the PMTs. The detector was surrounded by Cu bricks and sealed in a low ra-

dioactive air-tight Cu box continuously flushed with high purity nitrogen gas. The Cu box was surrounded by a passive shield made of 10 cm of high purity Cu, 15 cm of low radioactive lead, 1.5 mm of cadmium and 4/10 cm polyethylene/paraffin to reduce the external background. The whole shield has been closed inside a Plexiglas box, also continuously flushed by high purity nitrogen gas. An event-by-event data acquisition system accumulates the amplitude, the arrival time, and the pulse shape of the events.

The energy scale and the energy resolution of the ZnWO₄ detectors have been measured with γ sources ²²Na, ⁶⁰Co, ¹³³Ba, ¹³⁷Cs, ²²⁸Th, and ²⁴¹Am. The energy resolution of the detectors (full width at half of maximum) was in the range of (8.8 – 14.6)% for 662 keV γ line of ¹³⁷Cs.

As an example of the energy distribution accumulated in the measurements, in Fig. 1 the spectra collected in some runs are reported [4]. Some peaks in the spectra can be ascribed to γ lines of naturally occurring radionuclides such as ⁴⁰K, ²¹⁴Bi (²³⁸U chain) and ²⁰⁸Tl (²³²Th) from the materials of the set-up. A detailed investigation of the measured spectra has been performed by using various data analysis strategies and Monte Carlo simulations. In particular, the technique of the time-amplitude analysis (described in details in [7]) has been used in order to estimate the ²²⁸Th and ²²⁷Ac (²³⁵U family) activities. The pulse shape discrimination (PDS) capability of the ZnWO₄ scintillators has been considered in order to study the measured α spectra and the so called BiPo events. The radioactive contamination of the ZnWO₄ crys-

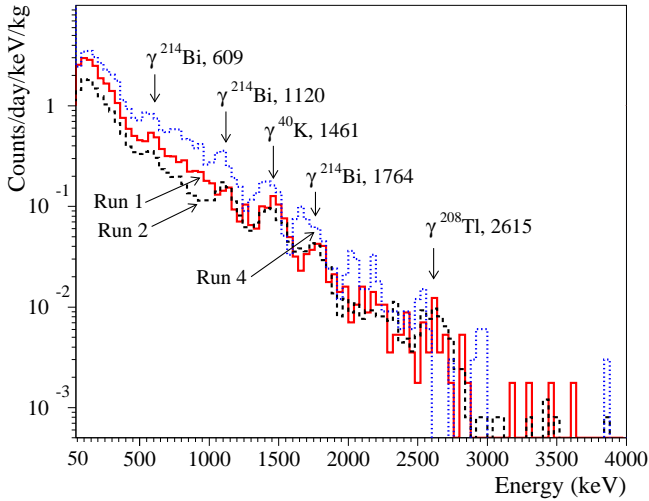


Figure 1. Energy distributions of the ZnWO_4 scintillators measured in the low background set-up during Runs 1 (with ZWO-1 in 2906 h) (solid red line), 2 (with ZWO-2 in 2130 h) (dashed black line), and 4 (with ZWO-4 in 834 h) (dotted blue line). Energies of γ lines from residual contaminations are in keV.

tals is on the level of 0.002 – 0.025 mBq/kg; the total α activity is in the range 0.2 – 2 mBq/kg. Moreover, particular contaminations associated with the composition of ZnWO_4 detector were observed [4]: the EC active cosmogenic nuclide ^{65}Zn – that can be also produced by neutrons – ($T_{1/2} = 244.26$ d [8]) with activity 0.5 – 0.8 mBq/kg (depending on the ZnWO_4 sample) and the α active tungsten isotope ^{180}W (with half-life: $T_{1/2} \approx 10^{18}$ yr [4, 9–11], and energy of the decay: $Q_\alpha = 2508(4)$ keV [12]) with activity 0.04 mBq/kg. In Fig. 2 it is shown the energy distribution of the $\beta(\gamma)$ events (identified by the PSD) accumulated in the low background set-up with the ZWO-4 crystal scintillator over 4305 h (Run 5) together with the model of the background. The main components of the background are shown: spectra of internal ^{65}Zn , ^{90}Sr - ^{90}Y , daughters of ^{238}U , and the contribution from the external γ quanta from PMTs and Cu box in these experimental conditions.

A summary of the radioactive contamination of the ZnWO_4 crystal scintillators can be found in ref. [4].

1.1 Results of the search for $\beta\beta$ decay in Zn and W isotopes and on rare α decay of W isotopes

The collected data have also been considered to search for double beta decay processes [5]. Zinc tungstate scintillators contain four potentially 2β active isotopes: ^{64}Zn , ^{70}Zn , ^{180}W and ^{186}W . It is worthwhile mentioning that ^{64}Zn and ^{186}W have comparatively large natural abundance that allows to apply ZnWO_4 detectors without using high cost enriched isotopes. Moreover, the $2\nu 2\beta^-$ decay of ^{186}W is expected to be strongly suppressed which could provide favorable conditions to search for neutrinoless $2\beta^-$ decays, including processes with emission of majoron(s) which have broad energy spectra, somewhat similar to that of the two-neutrino mode. The ^{180}W isotope is also an interesting 2β nuclide because in the case of the capture

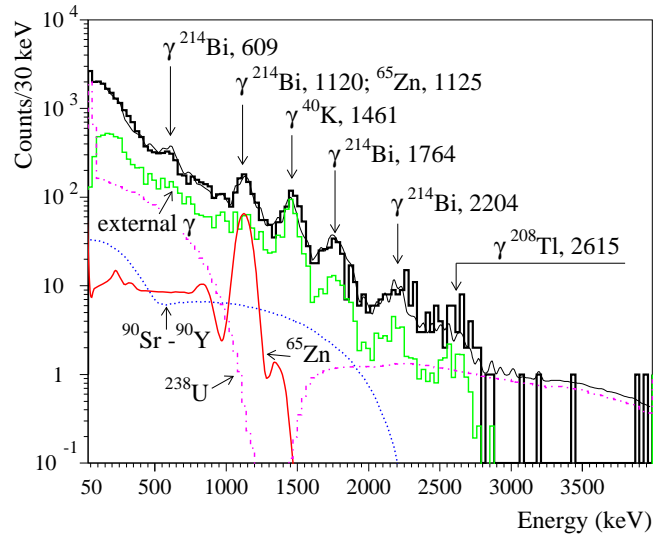


Figure 2. Energy distribution of the $\beta(\gamma)$ events (identified by the PSD) accumulated in the low background set-up with the ZWO-4 crystal scintillator over 4305 h (Run 5) together with the model of the background. The main components of the background are shown: spectra of internal ^{65}Zn , ^{90}Sr - ^{90}Y , daughters of ^{238}U , and the contribution from the external γ quanta from PMTs and Cu box in these experimental conditions.

of two electrons from the K shell ($E_K = 65.4$ keV), the decay energy is rather small (13 ± 4) keV; such a coincidence could give a resonant enhancement of the 0ν double electron capture to the corresponding level of the daughter nucleus.

The response functions of the ZnWO_4 detectors for the 2β processes in Zn and W isotopes were simulated with the help of the GEANT4 package [13] with the Low Energy Electromagnetic extensions. The initial kinematics of the particles emitted in the decays was generated with the DECAY0 event generator [14]. The background models included the internal contamination of the ZnWO_4 scintillators (^{40}K , ^{60}Co , ^{65}Zn , ^{87}Rb , ^{90}Sr - ^{90}Y , ^{137}Cs), and the external γ rays from radioactive contamination of the PMTs and the copper box (^{40}K , ^{232}Th , ^{238}U). Comparing the simulated response functions with the measured energy spectra of the ZnWO_4 detectors, no clear peculiarities, which can be evidently attributed to double beta decay of zinc or tungsten isotopes, have been found. Therefore only lower half-life limits have been set.

As an example, in Fig.3 the energy spectrum of the ZnWO_4 crystal scintillator $\varnothing 41 \times 27$ mm measured over 2798 h, corrected for the energy dependence of detection efficiency is reported, together with the $2\nu 2K$ peak of ^{64}Zn with $T_{1/2}^{2\nu 2K} = 1.1 \times 10^{19}$ yr excluded at 90% C.L.

Several $\beta\beta$ decay modes have been studied and competitive limits have been set. The main results are shortly summarized.

The previous limits on the $T_{1/2}$ of the $\beta\beta$ decay modes of ^{64}Zn , ^{70}Zn , ^{180}W and ^{186}W have been improved up to 2 orders of magnitude. In particular, new improved half-life limits on double electron capture and electron capture with positron emission in ^{64}Zn have been set in the range:

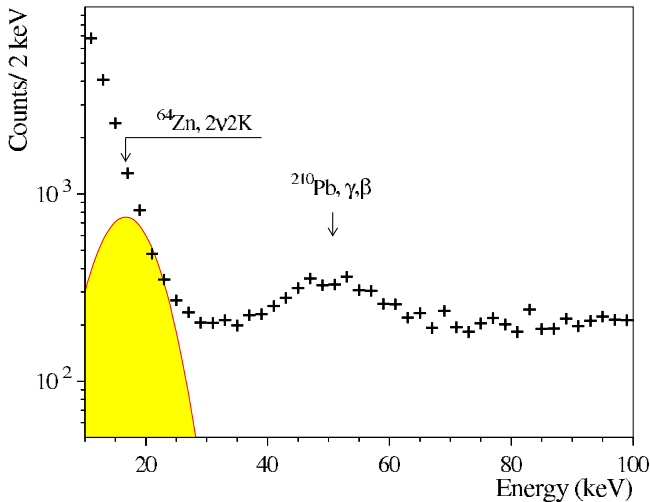


Figure 3. The energy spectrum of the ZnWO_4 crystal scintillator $\varnothing 41 \times 27$ mm measured over 2798 h, corrected for the energy dependence of detection efficiency, together with the $2\nu 2K$ peak of ^{64}Zn with $T_{1/2}^{2\nu 2K} = 1.1 \times 10^{19}$ yr excluded at 90% C.L.

10^{19} yr to 10^{21} yr depending on the mode. Similar sensitivity have been also reached for the 2β processes in ^{70}Zn , ^{180}W , and ^{186}W ($10^{18} - 10^{21}$ yr); note that the $0\nu 2\varepsilon$ capture in ^{180}W is of particular interest due to the possibility of resonant process [5]. The possible preliminary indication on the $(2\nu + 0\nu)\varepsilon\beta^+$ decay of ^{64}Zn with $T_{1/2} = (1.1 \pm 0.9) \times 10^{19}$ yr suggested in [15] is disproved by the results of these measurements. It is worth noting that to date only six nuclides (^{40}Ca , ^{78}Kr , ^{96}Ru , ^{106}Cd , ^{112}Sn , and ^{120}Te) among 34 candidates to 2ε , $\varepsilon\beta^+$, and $2\beta^+$ processes were studied at 10^{21} yr level of sensitivity in direct experiments. However, the limits are still far from theoretical predictions.

Rare α decay of ^{180}W has also been investigated in the same data by obtaining a new measurement of the half-life $T_{1/2} = 1.3_{-0.5}^{+0.6} \times 10^{18}$ yr. Also a new half-life limit on α transition of ^{183}W to the $1/2^-$ 375 keV metastable level of ^{179}Hf has been set as $T_{1/2} \geq 6.7 \times 10^{20}$ yr [5].

2 ZnWO_4 scintillator to investigate Dark Matter particle directionality

The above mentioned measurements and the correlated R&D works have shown that the ZnWO_4 scintillators can offer suitable features to investigate some Dark Matter (DM) particle candidates by exploiting the directionality technique. The use of anisotropic scintillators to study the directionality signature was proposed for the first time in ref. [16] and revisited in [17]. The directionality technique is effective only for those DM candidate particles able to induce just nuclear recoils. This approach studies the correlation between the arrival direction of the DM particles and the Earth motion in the galactic rest frame. In fact, the dynamics of the rotation of the Milky Way galactic disc through the halo of DM causes the Earth to experience a wind of DM particles apparently flowing along a direction opposite to that of the solar motion relative to the DM halo.

However, because of the Earth's rotation around its axis, their average direction with respect to an observer fixed on the Earth changes during the sidereal day. The possible nuclear recoils induced by the DM particles are expected to be strongly correlated with their impinging direction, while the background events are not; therefore, the study of the nuclear recoils direction can offer a way for pointing out the presence of the considered DM candidate particles.

The main advantages of ZnWO_4 scintillators to study the directionality are: i) high level of radio-purity reachable in future development, considering the very good results already achieved; ii) an energy threshold at level of few keV reachable (room for further improvement is possible); iii) light output of heavy particles (p , α , nuclear recoils) depending on the impinging direction of the particles with respect to the crystal axes, while the response to γ radiation being isotropic; iv) the scintillation decay time showing the same property.

The anisotropic effect has been ascribed to preferred directions of the excitons' propagation in the crystal lattice affecting the dynamics of the scintillation mechanism [18]. The anisotropic features of the ZnWO_4 detectors can provide two independent ways to exploit the directionality approach. In particular, the presence of heavy ionizing particles with a preferred direction (like recoil nuclei induced by the DM candidates considered here) could be discriminated from the electromagnetic background by comparing the low energy distributions measured by using different orientations of the crystal axes along the day. Moreover, the directionality technique can also be explored at some extent studying the time behaviour of the induced nuclear recoil pulses.

Thus, in the case of the ZnWO_4 detector, the anisotropy of the light output for recoiling nuclei induced by DM candidates could be discriminated from the electromagnetic events because of the expected variation of their detected energy distribution during the day [17]. In fact because of the Earth's daily rotation around its axis the preferential impinging direction of DM particle change during the sidereal time and the expected counting rate in a defined energy interval is expected to have a diurnal variation. This peculiarity of the rate can be considered to investigate the presence DM particle candidate inducing just nuclear recoils in the Galactic halo. Detailed discussion can be found in ref. [19].

3 Conclusion

In the last years radiopure ZnWO_4 crystal scintillator have been realized and have been measured in the DAMA/R&D facility at the Gran Sasso National Laboratory over more than 19 thousands hours. The measurements allowed us to study in details the radioactive residual contaminations of these crystals. The reached very good level of radiopurity and the study of the procedures followed to grow these crystals give us confidence for future developments of new ZnWO_4 crystal scintillators with higher level of radiopurity.

A low background experiment to search for 2β processes in ^{64}Zn , ^{70}Zn , ^{180}W , and ^{186}W was also carried out

with a total exposure of $0.5295 \text{ kg} \times \text{yr}$. New improved half-life limits on double beta decay modes in Zn and W isotopes in the range 10^{19} yr to 10^{21} yr have been set. The indication on the $(2\nu + 0\nu)\epsilon\beta^+$ decay of ^{64}Zn with $T_{1/2} = (1.1 \pm 0.9) \times 10^{19} \text{ yr}$ suggested in [15] is not confirmed. Note that to date only four nuclides (^{40}Ca , ^{78}Kr , ^{112}Sn , and ^{120}Te) among 34 candidates to 2ϵ , $\epsilon\beta^+$, and $2\beta^+$ processes were studied at 10^{21} level of sensitivity in direct experiments. However, it is worth noting that the limits are still far from theoretical predictions.

A search for rare α decay in W isotopes have also been performed. The rare α decay of ^{180}W with a half-life $T_{1/2} = 1.3^{+0.6}_{-0.5} \times 10^{18} \text{ yr}$ has been observed and new half-life limit $T_{1/2} \geq 6.7 \times 10^{20} \text{ yr}$ on α transition of ^{183}W to the $1/2^-$ 375 keV metastable level of ^{179}Hf has been set.

The anisotropic features of ZnWO_4 crystals make them very interesting target to investigate dark matter particles directionality.

Thus the ZnWO_4 crystal scintillators can be very good detectors for future developments and experiments either to study $\beta\beta$ decay modes and to investigate the directionality for some dark matter candidate particles.

References

- [1] F.A. Danevich et al., Nucl. Instr. Meth. A **544**, 553 (2005)
- [2] P. Belli et al., Phys. Lett. B **658**, 193 (2008)
- [3] P. Belli et al., Nucl. Phys. A **826**, 256 (2009)
- [4] P. Belli et al., Nucl. Instr. Meth. A **626-627**, 31 (2011)
- [5] P. Belli et al., J. Phys. G: Nucl. Part. Phys. **38**, 115107 (2011)
- [6] R. Bernabei et al., Astropart. Phys. **7**, 73 (1997); Nuovo Cim. A **110**, 189 (1997); P. Belli et al., Astropart. Phys. **10**, 115 (1999); Nucl. Phys. B **563**, 97 (1999); R. Bernabei et al., Nucl. Phys. A **705**, 29 (2002); P. Belli et al., Nucl. Instr. Meth. A **498**, 352 (2003); R. Cerulli et al., Nucl. Instr. Meth. A **525**, 535 (2004); R. Bernabei et al., Nucl. Instr. Meth. A **555**, 270 (2005); Ukr. J. Phys. **51**, 1037 (2006); P. Belli et al., Nucl. Phys. A **789**, 15 (2007); Phys. Rev. C **76**, 064603 (2007); Eur. Phys. J. A **36**, 167 (2008); Nucl. Instr. Meth. A **615**, 301 (2010); J. Phys. G: Nucl. Part. Phys. **38**, 015107 (2011); Phys. Rev. C **85**, 044610 (2012); A.S. Barabash et al., J. Instr. **6**, P08011 (2011)
- [7] F.A. Danevich et al., Phys. Lett. B **344**, 72 (1995); F.A. Danevich et al., Nucl. Phys. A **694**, 375 (2001)
- [8] R.B. Firestone et al., *Table of Isotopes*, 8-th ed., John Wiley, New York, 1996 and CD update, 1998
- [9] F.A. Danevich et al., Phys. Rev. C **67**, 014310 (2003)
- [10] C. Cozzini et al., Phys. Rev. C **70**, 064606 (2004)
- [11] Yu.G. Zdesenko et al., Nucl. Instr. Meth. A **538**, 657 (2005)
- [12] G. Audi, O. Bersillon, J. Blachot, and A.H. Wapstra, Nucl. Phys. A **729**, 337 (2003)
- [13] S. Agostinelli et al., Nucl. Instr. Meth. A **506**, 250 (2003); J. Allison et al., IEEE Trans. Nucl. Sci. **53**, 270 (2006)
- [14] O.A. Ponkratenko et al., Phys. At. Nucl. **63**, 1282 (2000); V.I. Tretyak, to be published
- [15] I. Bikit et al., Appl. Radiat. Isot. **46**, 455 (1995)
- [16] P. Belli et al., Il Nuovo Cim. C **15**, 475 (1992)
- [17] R. Bernabei et al., Eur. Phys. J. C **28**, 203 (2003)
- [18] J.B. Birks, *The theory and practice of scintillation counting* (Pergamon Press, London 1964); P.H. Heckmann, Z. Phys. **157**, 10 (1959); P.H. Heckmann et al., Z. Phys. **162**, 84 (1961); W.F. Kienzle, A. Flammersfeld, Z. Phys. **165**, 1 (1961); K. Tsukada, S. Kikuchi, Nucl. Instrum. Meth. **17**, 286 (1962); K. Tsukada et al., Nucl. Instrum. Meth. **37**, 69 (1965); F.J. Kratochwill, Z. Phys. **234**, 74 (1970); F.D. Brooks, D.T. Jones, Nucl. Instrum. Meth. **121**, 69 (1974)
- [19] P. Cappella et al., Eur. Phys. J. C **73**, 2276 (2013)