Recent Results from Experiment D0

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Abstract. Brief summary of resent experimental results from experiment D0 at TEVATRON in FNAL.

1 Introduction

After 28 years (1983 - 2011) of TEVATRON work in FNAL, particle physics reach high level of understanding of elementary interactions in center of mass energy up to 2 TeV. Two big particle detector, CDF and D0 could analyze all interactions detected with complex of subdetectors in full 4π cover and with magnetic spectrometry for all charged particles.

Briefs summary of some of results achieved at TEVATRON:
1985 - First observation of proton-antiproton collisions by CDF at 1.8 TeV (June 8).
1992 - D0 observes first proton-antiproton collision with cms energy 1.8 TeV.
1995 - Experiment CDF ad D0 announce discovery of top quark (March 3th).
1996 - Observation of antihydrogen atoms.
1996 - Observation of exotic charm meson.
1998 - Discovery of \( B_c \) meson.
1999 - Fixed target experiment KTeV observes direct CP violation in the decay of neutral Kaons.
2000 - The DONuT experiment reports first evidence for the direct observation of the \( \tau \) neutrino.
2004 - Run II, which started in 2001, achieves a peak luminosity \( 10^{32} \, cm^2 \, sec^{-1} \).
2006 - Discovery of Bs matter-antimatter oscillations 3 trillion times per second.
2007 - Discovery of cascade b-baryon.
2009 - Discovery of single top quark production.
2010 - Tevatron achieves a peak luminosity \( 4 \times 10^{32} \, cm^2 \, sec^{-1} \)
2011 - Tevatron produces final pron-antiproton collisions (sept 30)

Experiments D0 and CDF have collected about \( 10 \, fb^{-1} \) of data each.

The result from both experiment enriched knowledge in particle physics and fulfilled some open questions in Standard Model.

The analysis of data will still continue for few next years.
Following pictures present resent results from experiment D0 divided according the subjects: B-physics, Electroweek physics, Quatum Chromodynamics, Top quark physics, Higgs boson and New Phenomena. Most of the results are obtained with statistics of events corresponding to 10.4 fb$^{-1}$.

2 B-physics

Several results in spectroscopy of hadrons from is coming from study of physics of b quark, which could be directly identified in 60% of events. We present a measurement of the semileptonic mixing asymmetry $a = (\Gamma(\bar{A})-\Gamma(A))/\Gamma(A) + \Gamma(\bar{A})$ for $B^0$ mesons, $a_{s,l}^d$. Using two independent decay channels: $B^0 \rightarrow \mu + D + X$, $D \rightarrow K^+ \pi^-$ and $B^0 \rightarrow \mu + D^* + X$, $B^0 \rightarrow \mu + D^* + X$ with $D^* \rightarrow D^0 + \pi^-$, $D^0 \rightarrow K^+ \pi^-$ (and charge conjugate processes), we have determined the semileptonic mixing asymmetry for $B^0$ mesons, $a_{s,l}^d$. We extract the charge asymmetries in these two channels as a function of the visible proper decay length (VPDL) of the $B^0$ meson, correct for detector-related asymmetries using data-driven methods, and account for dilution from charge-symmetric processes using Monte Carlo simulation. The final measurement combines four signal VPDL regions for each channel, yielding:

$$a_{s,l}^d = 0.68 \pm 0.45 \text{ (stat.)} \pm 0.14 \text{ (syst.)}$$

Combination of measurements of $a_{s,l}^d$ (D0 and existing world-average from B factories), $a_{s,l}^s$, and the two impact-parameter-binned constraints from the same-charge dimuon asymmetry $A_{s,l}^b$. The bands represent the ±1 standard deviation uncertainties on each measurement. The ellipses represent the 1, 2, 3, and 4 standard deviation twodimensional confidence level regions of the combination [1] in Fig.2.

We have measure the time-integrated flavor-specific semileptonic charge asymmetry in the decays of $B^0_s$ mesons that have undergone flavor mixing, $a_{s,l}^s$, using $B^0_s(\bar{B}^0_s) \rightarrow D^\pm + \mu^+ + X$ decays, with $D^\pm \rightarrow \phi \pi^\pm$ and $\phi \rightarrow K^+ K^-$. of proton-antiproton collisions.
The significance of this observation is 5.6 standard deviations. The mass of the state is centered at charged muons, and the photon is identified through its conversion into an electron-positron pair. By the ATLAS Collaboration Fig. 3 [5].

\[
\nu_1 \twoBodyDec{J/\psi K^0} \rightarrow \Lambda_b^0 \rightarrow J/\psi K^0 \rightarrow J/\psi K^+ K^-(892)
\]

Figure 2. Invariant mass distribution of \(B^0\) candidates with \(ct > 200 \mu m\) for events in the mass range \(1.01 < M(K^+ K^-) < 1.03\) GeV. A fit to a sum of a Gaussian \(B^0 \rightarrow J/\psi K^0\) signal (dashed-dotted) a quadratic combinatorial background (dotted), and the reflection of the decay \(B^0 \rightarrow J/\psi K^*\) (892) (dashed), is used to extract the \(B^0\) yield (left). Proper decay length distributions for \(B^0 \rightarrow J/\psi K^0\) candidates, with fit results superimposed (right).

A fit to the difference between the time-integrated \(D^-_s\) and \(D^+_s\) mass distributions of the \(B^0_s\) and \(\bar{B}^0_s\) candidates yields the flavor-specific asymmetry

\[
a_{D_s} = -1.08 \pm 0.72(stat) \pm 0.17(syst)
\]

which is the most precise measurement and in agreement with the standard model prediction. We have investigated the decay \(B^0_s \rightarrow J/\psi \rightarrow K^+ K^-\) for invariant masses of the \(K^+ K^-\) pair in the range \(1.35 < M(K^+ K^-) < 2\) GeV Fig. 2.

From the study of the invariant mass and spin of the \(K^+ K^-\) system, we find evidence for the two-body decay and measure the relative branching fraction of the decays to be [2] \(R_{J/\psi} = 0.22 \pm 0.05\) (stat) \(\pm 0.04\) (syst). We measure the \(\Lambda_b^0\) lifetime in the fully reconstructed decay \(\Lambda_b^0 \rightarrow J/\psi \Lambda^0\) [3] Fig. 2.

The lifetime of the topologically similar decay channel \(B^0 \rightarrow J/\psi K^0\) is also measured [4]. We obtain

\[
\tau(\Lambda_b^0) = 1.303 \pm 0.075(stat.) \pm 0.035(syst.) \text{ps}
\]

and \(\tau(B^0) = 1.508 \pm 0.025(stat.) \pm 0.043(syst.) \text{ps}\).

Using these measurements, we determine the lifetime ratio of

\[
\tau(\Lambda_b^0)/\tau(B^0) = 0.864 \pm 0.052(stat.) \pm 0.033(syst.)
\]

We present a measurement of the relative branching fraction, \(R_{f_0/\phi}\), of \(B^0_s\) to \(J/\psi f_0(980)\), with \(f_0(980) \rightarrow \pi^+ \pi^-\), to the process \(B^0_s \rightarrow J/\psi f_0\), with \(\phi \rightarrow K^+ K^-\) Fig. 2. The \(J/\psi f_0(980)\) final state corresponds to a CP-odd eigenstate of \(B^0\) that could be of interest in future studies of CP violation.

With \(8 fb^{-1}\) of data recorded with the D0 detector we find

\[
R_{f_0/\phi} = 0.275 \pm 0.041(stat.) \pm 0.061(syst.)
\]

Using data corresponding to an integrated luminosity of \(1.3 fb^{-1}\), we observe a narrow mass state decaying into \(\nu(1S) + \gamma\), where the \(\nu(1S)\) meson is detected by its decay into a pair of oppositely charged muons, and the photon is identified through its conversion into an electron-positron pair. The significance of this observation is 5.6 standard deviations. The mass of the state is centered at \(10.551 \pm 0.014(stat.) \pm 0.017(syst.)\) GeV/c\(^2\), which is consistent with that of the state recently observed by the ATLAS Collaboration Fig. 3 [5].

3 Electroweak Interactions

Mass measurement of W mass combined from experiments CDF and D0 are in Fig. 4 [6].
We present a measurement of the W boson mass using data corresponding to 4.3 fb$^{-1}$ of integrated luminosity collected with the D0 detector during Run II at the Fermilab Tevatron $p\bar{p}$ collider. With a sample of 1,677,394 decay $W \to e\nu$ or $\mu\nu$ candidate events, we measure $M_W = 80.367 \pm 0.026$ GeV. This result is combined with an earlier D0 result determined using an independent Run II data sample, corresponding to 1 fb$^{-1}$ of integrated luminosity, to yield $M_W = 80.375 \pm 0.023$ GeV [7].

We study the processes $p\bar{p} \to WZ \to l\nu, l^+l^-$ and $p\bar{p} \to ZZ \to l^+l^-\nu\bar{\nu}$ where $l = e$ or $\mu$. Using 8.6 fb$^{-1}$ of integrated luminosity. We measure the $WZ$ production cross section to be $4.50^{+0.63}_{-0.66}$ pb which is consistent with, but slightly above a prediction of the standard model. The $ZZ$ cross section is measured to be $1.64 \pm 0.46$ pb, in agreement with a prediction of the standard model. Combination with an earlier analysis of the $ZZ \to l^+l^-\nu\bar{\nu}$ channel yields a $ZZ$ cross section of $1.44^{+0.35}_{-0.34}$ pb [9].

We study $WW$ and $WZ$ production with $b\bar{q}q(l = e, \mu)$ final states using data corresponding to 4.3 fb$^{-1}$ of integrated luminosity. Assuming the ratio between the production cross sections $\sigma(WW)$ and $\sigma(WZ)$ as predicted by the standard model, we measure the total $WW$ ($V=W,Z$) cross section to be $\sigma(WW) = 19.6^{+3.2}_{-3.0}$ pb, and reject the background-only hypothesis at a level of 7.9 standard deviations. We also use b-jet discrimination to separate the $WZ$ component from the dominant $WW$ component. Simultaneously fitting $WW$ and $WZ$ contributions, we measure $\sigma(WW) = 15.9^{+3.7}_{-3.2}$ pb and $\sigma(WZ) = 3.3^{+4.1}_{-3.3}$ pb, which is consistent with the standard model predictions [8].

We present a measurement of $p\bar{p} \to Z\gamma \to ll + \gamma, l = e, \mu$ production with a data sample corresponding to an integrated luminosity of 6.2 fb$^{-1}$. The results of the electron and muon channels are combined, and we measure the total production cross section and the differential cross section $d\sigma/dp_T^\gamma$, where $p_T^\gamma$ is the momentum of the photon in the plane transverse to the beamline Fig. 5. The results obtained are consistent with the standard model predictions from next-to-leading order calculations. We use the transverse momentum spectrum of the photon to place limits on anomalous $ZZ\gamma$ and $Z\gamma\gamma$ couplings [9]. We measure the cross section and the difference in rapidities between photons and charged leptons for inclusive $W \to l\nu + \gamma$ production in $e\gamma$ and $\mu\gamma$ final states. Using data corresponding to an integrated luminosity of 4.2 fb$^{-1}$, the cross section multiplied by the branching fraction for the process $p\bar{p} \to W\gamma + X \to l\nu\gamma + X$, measured to be $15.8 \pm 0.8$ (stat.) $\pm 1.2$ (syst.) pb,

![Figure 3](image-url) Dimuon invariant mass spectrum for opposite-charge pairs passing the muon selection criteria. The solid curve is a fit to the data assuming three $\nu$ resonances and a combinatorial background. The relative contributions from the $\nu(1S), \nu(2S)$, and $\nu(3S)$ (left). The distribution of $M_{\mu\mu} - M_{\nu(1S)}$ fit with three signal functions and the mixed event background (right).
Mass of the W Boson

<table>
<thead>
<tr>
<th>Measurement</th>
<th>$M_W$ [MeV]</th>
</tr>
</thead>
<tbody>
<tr>
<td>CDF-0/I</td>
<td>80432 ± 79</td>
</tr>
<tr>
<td>DØ-I</td>
<td>80478 ± 83</td>
</tr>
<tr>
<td>DØ-II (1.0 fb$^{-1}$)</td>
<td>80402 ± 43</td>
</tr>
<tr>
<td>CDF-II (2.2 fb$^{-1}$)</td>
<td>80387 ± 19</td>
</tr>
<tr>
<td>DØ-II (4.3 fb$^{-1}$)</td>
<td>80369 ± 26</td>
</tr>
<tr>
<td>Tevatron Run-0/II</td>
<td>80387 ± 16</td>
</tr>
<tr>
<td>LEP-2</td>
<td>80376 ± 33</td>
</tr>
<tr>
<td>World Average</td>
<td>80385 ± 15</td>
</tr>
</tbody>
</table>

Figure 4. Summary of measurement of W mass as March 2012

and the distribution of the charge-signed photon-lepton rapidity difference are found to be in agreement with the standard model. These results provide the most stringent limits on anomalous $WW\gamma$ couplings for data from hadron colliders: $-0.4 < \Delta \kappa\gamma < 0.4$ and $-0.08 < \lambda\gamma < 0.07$ at the 95% C.L.

4 QCD Interactions

We present a measurement of the average value of a new observable at hadron colliders that is sensitive to QCD dynamics and to the strong coupling constant, while being only weakly sensitive to parton...
momentum of the inclusive jets, in di all jets in an inclusive jet sample is measured and the results are presented as a function of transverse integrated luminosity of 0 requirements for the neighboring jets. The measurement is based on a data set corresponding to an of a photon in association with a b-quark jet for photons with rapidities photon transverse momentum Fig. 8. The b-quark jets are required to have Δy in next-to-leading order in the strong coupling constant, corrected for non-perturbative effects. From these results, we extract the strong coupling and test the QCD predictions for its running over a range of momentum transfers of 50 to 400 GeV [10] Fig. 6.

We present measurements of the differential cross section dσ/dp_{Tγ} for the inclusive production of a photon in association with a b-quark jet for photons with rapidities |y_{γ}| < 1.0 and 30 < p_{Tγ} < 300GeV, as well as for photons with 1.5 < |y_{γ}| < 2.5 and 30 < p_{Tγ} < 200GeV, where p_{Tγ} is the photon transverse momentum Fig. 8. The b-quark jets are required to have pT > 15GeV and rapidity |y_{jel}| < 1.5. The results are based on data corresponding to an integrated luminosity of 8.7fb^{-1}. The measured cross sections are compared with next-to-leading order perturbative QCD calculations using different sets of parton distribution functions as well as to predictions based on the kT-factorization QCD approach, and those from the Sherpa and Pythia Monte Carlo event generators [11].

In special runs of D0 experiment with Proton Forward Detector we have measured elastic scattering of p\bar{p} at \sqrt{s} = 1.960TeV. Comparison with other experiments is displayed in Fig. refQCD-3ab [12]. We investigate the decay B_{c}^{0} \rightarrow J/ψK^{+}K^{-} for invariant masses of the K^{+}K^{-} pair in the range 1.35 < M(K^{+}K^{-}) < 2 GeV. The data sample corresponds to an integrated luminosity of 10.4 fb^{-1} of p\bar{p} collisions at \sqrt{s} = 1.96 TeV. From the study of the invariant mass and spin of the K^{+}K^{-} system, we find evidence for the two-body decay and measure the relative branching fraction of the decays to be R_{f_{\bar{c}}/φ} = 0.22 \pm 0.05 \text{ (stat)} \pm 0.04 \text{ (syst)}. 

Figure 5. A comparison of the measured WW and WZ signals (filled histograms) to background-subtracted data (points) in the dijet mass distribution (summed over electron and muon channels) for 0, 1, and 2-tag sub-channels after the combined fit to data using the dijet mass distribution [9] (left). Unfolded (points) in the dijet mass distribution (summed over electron and muon channels) for 0, 1, and 2-tag sub-channels.
We measure correct jets in simulated events to replicate the quark flavor dependence of the jet response in data. Correcting for detector acceptance we obtain the forward-backward asymmetry 

\[ \Delta R = 1.8 \text{ (stat)} \pm 0.7 \text{ (syst) GeV} \] 

We also correct jets in simulated events to replicate the quark flavor dependence of the jet response in data. Correcting for detector acceptance we obtain the forward-backward asymmetry 

\[ \Delta R = 1.8 \text{ (stat)} \pm 0.7 \text{ (syst) GeV} \] 

The summary of cross section production of \( t\bar{t} \) is presented in Fig. 9, with final combination: 

\[ \sigma_{t\bar{t}} = 7.65 \pm 0.42 \text{ pb} \] 

We measure lepton angular distributions in \( t\bar{t} \rightarrow W^+ b, W^- \bar{b} \rightarrow l^+ \nu bl^- \). Using data corresponding to an integrated luminosity of 5.4 GeV, we find that the angular distributions of \( l^+ \) relative to antiprotons and \( l^- \) relative to protons are in agreement with each other. Combining the two distributions and correcting for detector acceptance we obtain the forward-backward asymmetry 

\[ A_{FB}^l = (5.8 \pm 5.1 \text{ (stat)} \pm 1.3 \text{ (syst)}) \] 

compared to the standard model prediction of \( A_{FB}^l \text{(predicted)} = (4.7 \pm 0.1) \). This result is further combined with the measurement based on the analysis of the 1-jets final state to obtain
\begin{figure}[ht]
\centering
\includegraphics[width=\textwidth]{plot.png}
\caption{Inclusive jet cross section measurements as a function of jet $p_T$ in six $|y|$ bins. The data points are multiplied by 2, 4, 8, 16, and 32 for the bins $1.6 < |y| < 2.0, 1.2 < |y| < 1.6, 0.8 < |y| < 1.2, 0.4 < |y| < 0.8,$ and $|y| < 0.4$, respectively (left). MC closure test of the method used to extract the inclusive jet $p_T$ cross section for the jet $|y| < 0.4$ bin. The full analysis was repeated treating MC events as data and comparing the result to the input cross section. Good agreement is found within the statistical uncertainties of fits to jet energy scale and resolution, and unfolding present in MC (shaded band), which are much smaller than the systematic uncertainties in data (right).}
\end{figure}

$A_{FB}^t = (11.8 \pm 3.2)$ [14].

We present measurements of the $tWb$ coupling form factors using information from electroweak single top quark production and from the helicity of $W$ bosons from top quark decays in $t\bar{t}$ events Fig. 10 [15]. We set upper limits on anomalous $tWb$ coupling.

Combining measurements that simultaneously determine the fractions of $W$ bosons with longitudinal ($f^0$) and right-handed ($f^+$) helicities, we find

\begin{align*}
    f^0 &= 0.722 \pm 0.081 [\pm 0.062 (\text{stat.}) \pm 0.052 (\text{syst.})] \\
    f^+ &= -0.033 \pm 0.046 [\pm 0.034 (\text{stat.}) \pm 0.031 (\text{syst.})].
\end{align*}

Combining measurements where one of the helicity fractions is fixed to the value expected in the standard model, we find

\begin{align*}
    f^0 &= 0.682 \pm 0.057 [\pm 0.035 (\text{stat.}) \pm 0.046 (\text{syst.})] \\
    f^+ &= -0.015 \pm 0.035 [\pm 0.018 (\text{stat.}) \pm 0.030 (\text{syst.})] [16].
\end{align*}

The results are consistent with standard model expectations.

The $t\bar{t}$ spin correlation strength $C$ is defined by

\[d^2\sigma_{\bar{t}t}/d\cos\Theta_1d\cos\Theta_2 = \sigma_{\bar{t}t}(1 - C\cos(\Theta_1)\cos(\Theta_2))/4,\]

where $\Theta_1, \Theta_2$ are the angles between the spin-quantization axis and the direction of flight of the down-type fermion from the $W$ boson decay in the respective parent $t$ or $\bar{t}$ rest frame. The fractional difference in the number of events where the top and antitop quark spins are aligned and those where the top quark spins have opposite alignment $A$ is defined as $C = A[\alpha_1\alpha_2]$ where $\alpha_i$ is the spin analyzing power of the final state parton under consideration. In NLO QCD $\alpha_{l^+} = 1$ for the charged lepton in $t \to l + \nu_l b$ decays and $\alpha_d = 0.97$ for the anti-down quark in $t \to d\bar{u}b$ decays. To distinguish
Figure 8. MC closure test of the method used to extract the inclusive jet $p_T$ cross section for the jet $|y| < 0.4$ bin. The full analysis was repeated treating MC events as data and comparing the result to the input cross section. Good agreement is found within the statistical uncertainties of fits to jet energy scale and resolution, and unfolding present in MC (shaded band), which are much smaller than the systematic uncertainties in data (left). $d\sigma/dt$ differential cross section of elastic scattering measured by the D0 Collaboration and compared to the CDF and E710 measurements at $\sqrt{s} = 1.8$ TeV, and to the UA4 measurement at $\sqrt{s} = 0.546$ TeV (scaled by a factor of 100) (right).

Figure 9. The measurements of $\sigma(t\bar{t})$ cross section from the CDF and D0 experiments (left). Reconstructed mass of top quark from events with 2 leptons and 2 jets (right).
**Figure 10.** Likelihood contours at the 68% C.L. and the 95% C.L. as a function of W boson helicity fractions. Statistical uncertainties and systematic uncertainties that are uncorrelated with the single top quark measurement are included. The squares, triangles and upside-down triangles show \( f_R V \), \( f_L T \) and \( f_R T \) varying in fifty equal-size steps such that their ratio to \( f_L V \) goes from zero to ten-to-one. The dashed triangle denotes the physically allowed region. (left). Discriminant R for measurement of top-antitop spin correlation [17](right).

**Figure 11.** Search for \( T \bar{T} \) resonance in mass distribution.
between correlated and uncorrelated top quark spin hypotheses, we define a discriminant $R$, which is displayed in Fig. 10 together with MC events with and without spin correlation as well as background. Data support the hypothesis with top-antitop spin correlation $C = 0.85 \pm 0.29$ [17] which is in good agreement with the SM prediction.

The total width of top quark $\Gamma_t$ is extracted from the partial decay width $\Gamma(t \rightarrow Wb)$ and the branching fraction $\mathcal{B}(t \rightarrow Wb)$, $\Gamma(t \rightarrow Wb)$ is obtained from the $t$-channel single top quark production cross section and $\mathcal{B}(t \rightarrow Wb)$ is measured in $t\bar{t}$ events. For a top mass of 172.5 GeV, the resulting width is $\Gamma_t = 2.00^{+0.47}_{-0.43}$ GeV. This translates to a top-quark lifetime of $\tau_t = (3.29^{+0.90}_{-0.63}) \times 10^{-25}$ s [18].

We have also extract an improved direct limit on the CKM matrix element $0.81 < |V_{tb}| \leq 1$ at 95% C.L. and a limit of $|V_{tb}| < 0.59$ for a high mass fourth generation bottom quark assuming unitarity of the fourth generation quark mixing matrix. We searched for a narrow $t\bar{t}$ resonance that decays into a lepton+jets final state based on an integrated luminosity of $5.3 fb^{-1}$ of proton-antiproton collision. We set upper limits on the production cross section of such a resonance multiplied by its branching fraction to $t\bar{t}$ which we compare to predictions for a leptophobic topcolor $Z'$ boson. We exclude such a resonance at the 95% confidence level for masses below 835 GeV Fig. 11 [19]. We present new direct constraints on a general $Wtb$ interaction using data corresponding to an integrated luminosity of $5.4 fb^{-1}$. The standard model provides a purely left-handed vector coupling at the $Wtb$ vertex, while the most general, lowest dimension Lagrangian allows right-handed vector and left- or right-handed tensor couplings as well.

We obtain precise limits on these anomalous couplings by comparing the data to the expectations from different assumptions on the $Wtb$ coupling [20].

We present measurements of production cross sections of single top quarks in $p\bar{p}$ collisions at $\sqrt{s} = 1.96 TeV$ [21] in a data sample corresponding to an integrated luminosity of $5.4 fb^{-1}$. We select events with an isolated electron or muon, an imbalance in transverse energy, and two, three, or four jets, with one or two of them containing a bottom hadron. We obtain an inclusive cross section of $(p\bar{p} \rightarrow tb + X)$,

$$\sigma(tqb + X) = 3.43 \pm 0.73 \text{ pb}$$

and use it to extract the CKM matrix element $0.79 < |V_{tb}| \leq 1$ at the 95% C.L. We also measure $\sigma(p\bar{p}tb + X) = 0.68_{-0.35}^{+0.38} \text{ pb}$ and $\sigma(p\bar{p}tb + X) = 2.86_{-0.63}^{+0.69} \text{ pb}$

### 6 HIGGS boson

We combine searches by the CDF and D0 Collaborations for the associated production of a Higgs boson with a W or Z boson and subsequent decay of the Higgs boson to a bottom-antibottom quark pair.

The data correspond to integrated luminosities of up to $9.7 fb^{-1}$ Fig. 12 [22]. The searches are conducted for a Higgs boson with mass in the range 100 - 150 GeV. We observe an excess of events in the data compared with the background predictions, which is most significant in the mass range between 120 and 135 GeV/$c^2$. The largest local significance is 3.3 standard deviations, corresponding to a global significance of 3.1 standard deviations. We interpret this as evidence for the presence of a new particle consistent with the standard model Higgs boson, which is produced in association with a weak vector boson and decays to a bottom-antibottom quark pair [23].

### 7 New Phenomena

We present the first search for supersymmetry (SUSY) in $Z\gamma$ final states with large missing transverse energy using data corresponding to an integrated luminosity of $6.2 fb^{-1}$ [24]. This signature is predicted in gauge-mediated SUSY-breaking models, where the lightest neutralino is the next-to-lightest
Figure 12. Background-subtracted distribution of the reconstructed dijet mass $m_{jj}$, summed over all input channels. The $VZ$ signal and the background contributions are fit to the data, and the fitted background is subtracted. The fitted $VZ$ and expected SM Higgs, $m_H = 125\text{GeV}$ contributions are shown with filled histograms (left). The dark and light-shaded regions indicate the 1 s.d. and 2 s.d. measurement uncertainties, and the SM prediction is shown as the smooth, falling curve with a narrow band indicating the theoretical uncertainty. The expected cross section fit values assuming the SM Higgs boson is present at $m_H = 125\text{GeV}/c^2$ SM rate (dark blue) and the best fitted rate from data (light magenta) (right).

Figure 13. Search for Stop quark and Gauging-like Chargino masses (left).

supersymmetric particle (NLSP) and is produced in pairs, possibly through decay from heavier supersymmetric particles. The NLSP can decay either to a $Z$ boson or a photon and an associated gravitino that escapes detection. We exclude this model at the 95% C.L. for SUSY breaking scales of $\Lambda < 87\text{TeV}$, corresponding to neutralino masses of $< 151\text{GeV}$.

We present a search for Kaluza-Klein (KK) particles predicted by models with universal extra dimensions (UED) using a data set corresponding to an integrated luminosity of 7.3 fb$^{-1}$ [25]. The decay chain of KK particles can lead to a final state with two muons of the same charge. This signature is used to set a lower limit on the compactification scale of $R^{-1} > 260 \text{GeV}$ in a minimal UED model.
Figure 14. Search for Neutralino and new Z boson decaying into top-antitop pair.

The Tevatron has finished data taking, but the analysis of data with full statistics of reconstructed events will still continue and new results will come in the next year.

References

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