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## Classical and quantum parts in Madelung variables

### Splitting the source term of the Einstein equation into classical and quantum parts

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**Abstract.** Postulating a particular quantum correction to the source term in the classical Einstein equation we identify the conformal content of the above action and obtain classical gravitation for massive particles, but with a cosmological term representing off-mass-shell contribution to the energy-momentum tensor.

## 1 Introduction

In this short paper we report an interesting observation based on a conformal treatment of the Schrödinger equation: as if the quantum mechanical problem of obtaining wave functions could be splitted to a massive and a conformal part in line with a classical – quantum partition. Identifying the quantum part as belonging to a traceless relativistic energy-momentum tensor, we suggest to modify its Bohm-Takabayashi form [1] and connect the remaining classical part to Einstein's gravity equation in form of a dust matter source of massive point particles moving on Bohm trajectories. In this scenario the quantum nature of the wave function reveals itself in deviations from the classical on-mass-shell relation  $P_\mu P^\mu = (mc)^2$ , and our suggested natural coupling to gravity makes a simple conformal transformation of the full Einstein tensor expedient. The details of the outlined calculations are given in [2].

## 2 Energy-momentum tensors of the Klein-Gordon equation

In the following Greek letters denote four-indices and  $g_{\mu\nu}$  is the flat Lorentzian spacetime metric.  $\hbar$ ,  $c$  and  $m$  denotes the Planck constant, the speed of light and the mass of a particle.

The textbook form of the Klein-Gordon energy-momentum tensor reads as

$$T_{\mu\nu} = mcR^2 u_\mu u_\nu + \frac{\hbar^2}{2mc} \left( 2\partial_\mu R \partial_\nu R - g_{\mu\nu} (\partial_\alpha R \partial^\alpha R + R \square R) \right). \quad (1)$$

Here the wave function is represented by the real fields  $R$  and  $\alpha$ , according to the following decomposition:

$$\psi = \frac{\hbar}{\sqrt{mc}} R e^{\frac{i}{\hbar}\alpha}, \quad (2)$$

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and the  $u_\mu = \partial_\mu \alpha / (mc)$  vector field is analogous to a four-velocity with an off-mass shell dispersion relation. In a metric view this dispersion relation is:

$$g_{\mu\nu} u^\mu u^\nu = 1 + \left( \frac{\hbar}{mc} \right)^2 \frac{\square R}{R}. \quad (3)$$

This formula, multiplied by  $R^2/2$  can be rewritten as a full divergence of the *Bohm-Takabayashi* energy-momentum tensor

$$\mathcal{T}_{\mu\nu} = mcR^2 u_\mu u_\nu - \frac{\hbar^2}{2mc} \left( R \partial_\mu \partial_\nu R - \partial_\mu R \partial_\nu R \right). \quad (4)$$

The full energy-momentum tensor can be obtained by the linear combination of the above Klein-Gordon and Takabayashi forms, requiring, that only the classical dust contributes to the trace and the part proportional to  $\hbar^2$  – the quantum contribution – is traceless. Introducing dilaton field like variable [3–6] by  $\sigma = \ln(\sqrt{V}R)$ , where  $V$  is constant, it can be written as:

$$\mathfrak{T}_\nu^\mu = \frac{mc}{V} e^{2\sigma} u^\mu u_\nu + \frac{\hbar^2}{2mcV} e^{2\sigma} \left( 2\lambda \partial^\mu \sigma \partial_\nu \sigma + (\lambda - 1) \partial^\mu \partial_\nu \sigma - \lambda \delta_\nu^\mu (2\partial_\alpha \sigma \partial^\alpha \sigma + \square \sigma) \right). \quad (5)$$

### 3 General relativistic quantum mechanics

The previous formula is interpreted in a general relativistic framework by a conformal transformation of a slightly modified Einstein equation:

$$G_\nu^\mu + \Lambda \delta_\nu^\mu = \left[ \bar{G}_\nu^\mu + (2\square s + \partial_\alpha s \partial^\alpha s + \Lambda) \delta_\nu^\mu - 2\partial^\mu \partial_\nu s + 2\partial^\mu s \partial_\nu s \right] = \frac{8\pi G}{c^3} e^{-2\sigma} \mathfrak{T}_\nu^\mu. \quad (6)$$

Where the modification is in the source term multiplied by  $e^{-2\sigma}$ .

The term by term comparison of (5) and (6) reveals that  $s = \sigma$  leads to a classical Einstein equation governed by the dust moving on Bohmian trajectories

$$\bar{G}_\nu^\mu = \frac{8\pi L_S}{V} u^\mu u_\nu. \quad (7)$$

We obtain also a cosmological term proportional to the off-mass-shell effect. That is, we have splitted the slightly modified Einstein equation with a cosmological term (6) in the form

$$G_{\mu\nu}^{flat} - 3 \frac{\square R}{R} g_{\mu\nu}^{flat} = \frac{8\pi G}{c^3} e^{-2\sigma} \left( T_{\mu\nu}^{classical} + T_{\mu\nu}^{quantum} \right), \quad (8)$$

and obtained a quantum interpretation for the cosmological term in flat metric. In this scenario the - on the Planck scale surprisingly small - cosmological constant stems from quantum binding with a Bohr radius  $a$  as being  $\Lambda = 3/a^2$ . This is the same relation as for the de Sitter cosmological horizon.

### References

- [1] T. Takabayasi, *Progress of Theoretical Physics*, **8**(2), 143–182 (1952).
- [2] T.S. Biró and P. Ván, arXiv:1312.1316.
- [3] C.H. Brans and R.H. Dicke, *Physical Review*, **124**(3), 925–935 (1961).
- [4] D.H. Delphenic, arXiv:gr-qc/0211065.
- [5] S. Capozziello and V. Faraoni, *Beyond Einstein Gravity*, Springer, 2011.
- [6] Y. Fujii and K.-I. Maeda, *The Scalar-Tensor Theory of Gravitation*. Cambridge University Press, 2004.