

Highlight of Charm Physics at BESIII

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Abstract. In this talk, we present the recent results of charm physics from the BESIII collaboration. This talk covers the studies of the leptonic and semi-leptonic decays of D mesons, and the measurement of the strong phase difference based on the quantum correlation.

1 Introduction

The BESIII [1] experiment at the BEPCII collider started data taking since 2008. For the study of charm meson decays, about 2.9 fb^{-1} data on the $\psi(3770)$ peak has been accumulated. $\psi(3770)$ dominantly decays into $D\bar{D}$, which provides an ideal place for studying the decays of D^0 and D^+ mesons.

2 Determination of $\mathcal{B}(D^+ \rightarrow \mu^+ \nu)$ and the decay constant f_{D^+}

In the Standard Model of particle physics, the D^+ meson can decay into $l^+ \nu_l$ via a virtual W^+ boson. The decay rate depends on the D^+ decay constant f_{D^+} , in which all of the strong interaction effects between the two initial-state quarks are absorbed. The decay width of $D^+ \rightarrow l^+ \nu_l$ is given by [2]

$$\Gamma(D^+ \rightarrow l^+ \nu_l) = \frac{G_F^2 f_{D^+}^2}{8\pi} |V_{cd}|^2 m_l^2 m_{D^+} \left(1 - \frac{m_l^2}{m_{D^+}^2}\right)^2, \quad (1)$$

where G_F is the Fermi coupling constant. So by measuring the branching ratio of $D^+ \rightarrow l^+ \nu_l$, the decay constant f_{D^+} can be determined to test the calculations from lattice QCD. In addition, $|V_{cd}|$ can be directly accessed. f_{D^+} is also helpful to constrain the CKM matrix element $|V_{td}|$ through its relation to the mixing parameter x_B and the ratio f_D/f_B [3, 4].

To determine the branching ratio of $D^+ \rightarrow \mu^+ \nu_\mu$, the D tagging method is used. Firstly, D^- mesons are reconstructed with 9 hadronic decay modes ($D^- \rightarrow K^+ \pi^- \pi^-, K_S^0 \pi^-, K_S^0 K^-, K^+ K^- \pi^-, K^+ \pi^- \pi^- \pi^0, \pi^+ \pi^- \pi^-, K_S^0 \pi^- \pi^0, K^+ \pi^- \pi^- \pi^- \pi^+$ and $K_S^0 \pi^- \pi^- \pi^+$), and then μ^+ is sought out in the remaining tracks. $M_{miss}^2 = (E_{beam} - E_{D^-} - E_{\mu^+})^2 - (\vec{p}_{D^-} - \vec{p}_{\mu^+})^2$ is used to extract the signal events which peak

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around 0. Finally, with 2.9 fb^{-1} $\psi(3770)$ data, the branching fraction of $D^+ \rightarrow \mu^+ \nu_\mu$ is measured to be [5]

$$\mathcal{B}(D^+ \rightarrow \mu^+ \nu_\mu) = (3.71 \pm 0.19 \pm 0.06) \times 10^{-4}.$$

With combination of G_F , the mass of μ^+ and D^+ , the lifetime of D^+ and $|V_{cd}|$ from the global fit in the SM [6], the decay constant f_{D^+} is obtained with the best precision in the world [5]

$$f_{D^+} = (203.2 \pm 5.3 \pm 1.8) \text{ MeV}.$$

3 Measurement of $\mathcal{B}(D^0 \rightarrow K^-/\pi^- e^+ \nu)$ and the Form-Factor $f_+(q^2)$

The differential decay rate of $D^0 \rightarrow K^-(\pi^-)e^+ \nu$ is written as

$$\frac{\Delta\Gamma(D \rightarrow K(\pi)e\nu)}{dq^2} = \frac{G_F^2 |V_{cs(d)}|^2}{24\pi^3} p^3 |f_+(q^2)|^2. \quad (2)$$

With the precisely determined $V_{cs(d)}$ from the unitarity of the CKM matrix, we can obtain the form-factor $f_+(q^2)$ to check the lattice QCD calculations.

This analysis has been done using one-third of the 2.9 fb^{-1} $\psi(3770)$ data, with a partial-blind method. Firstly, the singly tagged \bar{D}^0 is reconstructed with four hadronic modes ($\bar{D}^0 \rightarrow K^+ \pi^-, K^+ \pi^- \pi^0, K^+ \pi^- \pi^0 \pi^0$ and $K^+ \pi^- \pi^- \pi^+$). The signal candidates are searched by reconstructing two opposite charged tracks ($K^+/\pi^+ e^-$) left in the event, and $U_{miss} = E_{miss} - p_{miss}$ is used to extract the signal events which peak around 0. Here, $E_{miss} = E_{beam} - E_{hadron} - E_{electron}$, and $p_{miss} = |-\vec{p}_{hadron} - \vec{p}_{electron}|$. The branching fractions are measured to be

$$\mathcal{B}(D^0 \rightarrow K^- e^+ \nu) = (3.542 \pm 0.030 \pm 0.067)\%$$

$$\mathcal{B}(D^0 \rightarrow \pi^- e^+ \nu) = (0.288 \pm 0.008 \pm 0.005)\%.$$

These two results are consistent with the world averages [6].

To measure the form factor, the partial decay rates are measured in different q^2 bins, where q^2 is the squared invariant mass of $e^+ \nu$ system. The $f_+(q^2)$ distributions are shown in figure 1 with theoretical curves overlaid [7].

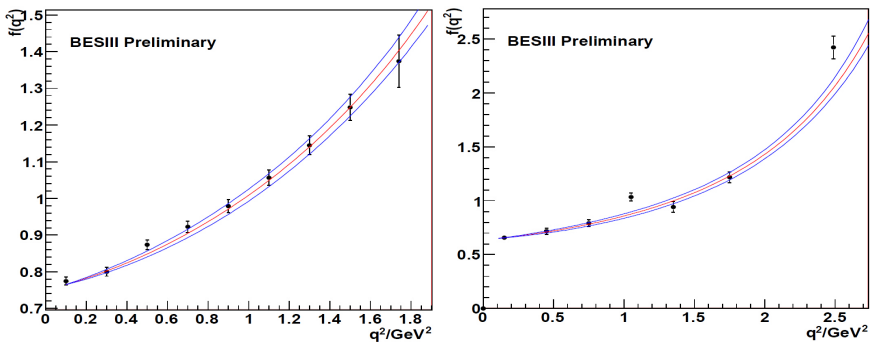


Figure 1. $f_+(q^2)$ distributions for the decay of $D^0 \rightarrow K^- e^+ \nu$ (left) and $D^0 \rightarrow \pi^- e^+ \nu$ (right). Points are measured from data, curves are the theoretical predictions varying within one statistical standard sigma.

4 Measurement of the strong phase difference $\delta_{K\pi}$

Studying the D^0 - \bar{D}^0 mixing is important for identifying the size of the long distance effect and searching for new physics [8], and also helpful for studying CP violation in charm physics. Charm mixing is described by two dimensionless parameters x and y . The measurement of the strong phase difference δ between the doubly Cabibbo-suppressed (DCS) decay $\bar{D}^0 \rightarrow K^- \pi^+$ and the corresponding Cabibbo-favored (CF) $D^0 \rightarrow K^- \pi^+$, allows x and y to be extracted from $x' \equiv x \cos \delta + y \sin \delta$ and $y' \equiv y \cos \delta - x \sin \delta$ [9]. Furthermore, finer precision of δ helps the γ/ϕ_3 angle measurement in CKM matrix according to the so-called ADS method [6].

Using the quantum-correlated technique, δ can be accessed using the following formula

$$2r \cos \delta + y = (1 + R_{\text{WS}}) \cdot \mathcal{A}_{CP \rightarrow K\pi}, \quad (3)$$

where $\mathcal{A}_{CP \rightarrow K\pi}$ is the asymmetry between the branching fractions of CP-odd and CP-even states decaying to $K^- \pi^+$. $\mathcal{B}(D^{CP^\pm} \rightarrow K^- \pi^+)$ are measured with D tagging method using 5 CP-even D^0 decay modes ($K^+ K^-$, $\pi^+ \pi^-$, $K_S^0 \pi^0 \pi^0$, $\pi^0 \pi^0$ and $\rho^0 \pi^0$) and 3 CP-odd modes ($K_S^0 \pi^0$, $K_S^0 \eta$ and $K_S^0 \omega$). With 2.9 fb^{-1} $\psi(3770)$ data, we obtain the asymmetry to be [10]

$$\mathcal{A}_{CP \rightarrow K\pi} = (12.7 \pm 1.3 \pm 0.7)\%.$$

By quoting the external inputs of $R_D = r^2 = (3.50 \pm 0.04)\%$, $y = (6.7 \pm 0.9)\%$ from HFAG 2013 [11] and $R_{\text{WS}} = (3.80 \pm 0.05)\%$ from PDG [6], we obtain [10]

$$\cos \delta_{K\pi} = 1.02 \pm 0.11 \pm 0.06 \pm 0.01,$$

where the third uncertainty is due to the errors introduced by the external input parameters. This result provides the world best constrain to $\delta_{K\pi}$.

5 Measurement of the Relative Strong-phase Difference between D^0 and \bar{D}^0 decay to $K_S^0 \pi^+ \pi^-$

The CKM angle γ/ϕ_3 can be measured from the decay $B^+ \rightarrow D(K_S^0 \pi^+ \pi^-) K^\pm$ through the GGSZ method [12, 13]. The Dalitz plot can be separated into $2N$ bins as shown in figure 2(a). The N bins lying below the symmetry axis are denoted by the index i , while the remaining bins are indexed with \bar{i} , and these square bins are symmetric under exchange of x and y ($x \equiv m_{K_S^0 \pi^-}^2$, $y \equiv m_{K_S^0 \pi^+}^2$). The two parameters c_i and s_i denote the weighted average of $\cos(\Delta\delta_D)$ and $\sin(\Delta\delta_D)$ in the i^{th} bin, respectively. And $\Delta\delta_D$ is the strong phase difference between the DCS decay $\bar{D}^0 \rightarrow K_S^0 \pi^+ \pi^-$ and the corresponding CF decay $D^0 \rightarrow K_S^0 \pi^+ \pi^-$.

The number of events in the i^{th} bin of a Dalitz plot originating from the B^\pm decay is

$$N_i^\pm = \frac{a_B}{a_D} [K_i + r_B^2 K_{\bar{i}} \pm 2r_B \sqrt{K_i K_{\bar{i}}} (\cos(\delta_B + \gamma) c_i + \sin(\delta_B + \gamma) s_i)]. \quad (4)$$

The normalization factors, r_B and δ_B could be determined at B factories. Using $\psi(3770)$, K_i can be obtained by flavor tagged $D \rightarrow K_S^0 \pi^+ \pi^-$ decays, c_i can be obtained from CP tagged $D \rightarrow K_S^0 \pi^+ \pi^-$ decays, and s_i is obtained from double Dalitz decays $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ and $\bar{D}^0 \rightarrow K_S^0 \pi^+ \pi^-$.

The Dalitz plot is split into 8 bins as shown in figure 2(b), which have been optimized, to get the highest sensitivity of γ . The results of c_i and s_i in each bin are given in figure 2(c). Our results will allow for increased precision in the measurement of γ/ϕ_3 .

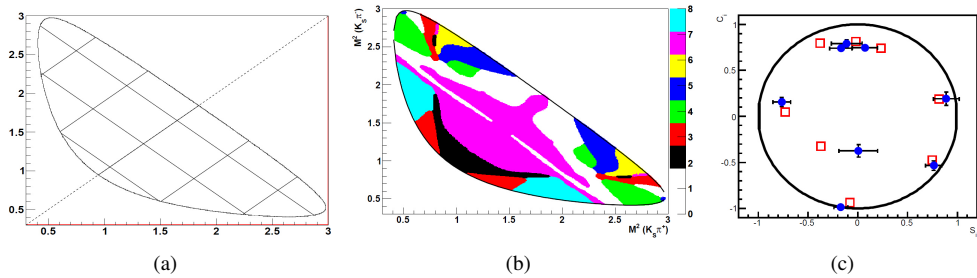


Figure 2. (a): Squared binned Dalitz plot with symmetric bins over an exchange of x and y . (b): Modified optimal binning. (c): Red squares mark the model predicted value of c_i and s_i [14], and the blue dot mark fitted value from data with statistical error.

6 Summary

Based on the 2.9 fb^{-1} data taken at the $\psi(3770)$ peak, the recent results on the study of charm physics are presented. The BESIII experiment aims at getting 20 fb^{-1} $\psi(3770)$, and more promising results are expected.

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References

- [1] M. Ablikim *et al.* [BESIII Collaboration], Phys. Rev. Lett. **110**, 252001 (2013)
- [2] D. Silverman and H. Yao, Phys. Rev. D **38**, 214 (1988)
- [3] C. Bernard *et al.*, arXiv:hep-ph/9709328
- [4] K. Hara *et al.* [Belle Collaboration], Phys. Rev. Lett. **110**, 131801 (2013)
- [5] M. Ablikim *et al.* [BESIII Collaboration], Phys. Rev. D **89**, 051104 (2014)
- [6] J. Beringer *et al.* [Particle Data Group], Phys. Rev. D **86**, 010001 (2012)
- [7] John A. Bailey *et al.*, arXiv:1111.5471
- [8] S. Bianco, F. L. Fabbri, D. Benson and I. Bigi, Riv. Nuovo Cim. **26N7**, 1 (2003)
- [9] R. Aaij *et al.* [LHCb Collaboration], Phys. Rev. Lett. **110**, 101802 (2013)
- [10] M. Ablikim *et al.* [BESIII Collaboration], Phys. Lett. B **734**, 227(2014)
- [11] Heavy Flavor Averaging Group: <http://www.slac.stanford.edu/xorg/hfag/charm/>
- [12] A. Giri, Y. Grossman, A. Soffer, and J. Zupan, Phys. Rev. D **68**, 054018 (2003)
- [13] A. Bondar, A. Poluektov, Eur. Phys.J.C **47**, 347 (2006)
- [14] B. Aubert *et al.* [BABAR Collaboration], Phys. Rev. D **78**, 034023 (2008)