

## Highlight of Charm Physics at BESIII

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**Abstract.** In this talk, we present the recent results of charm physics from the BESIII collaboration. This talk covers the studies of the leptonic and semi-leptonic decays of  $D$  mesons, and the measurement of the strong phase difference based on the quantum correlation.

### 1 Introduction

The BESIII [1] experiment at the BEPCII collider started data taking since 2008. For the study of charm meson decays, about  $2.9 \text{ fb}^{-1}$  data on the  $\psi(3770)$  peak has been accumulated.  $\psi(3770)$  dominantly decays into  $D\bar{D}$ , which provides an ideal place for studying the decays of  $D^0$  and  $D^+$  mesons.

### 2 Determination of $\mathcal{B}(D^+ \rightarrow \mu^+ \nu)$ and the decay constant $f_{D^+}$

In the Standard Model of particle physics, the  $D^+$  meson can decay into  $l^+ \nu_l$  via a virtual  $W^+$  boson. The decay rate depends on the  $D^+$  decay constant  $f_{D^+}$ , in which all of the strong interaction effects between the two initial-state quarks are absorbed. The decay width of  $D^+ \rightarrow l^+ \nu_l$  is given by [2]

$$\Gamma(D^+ \rightarrow l^+ \nu_l) = \frac{G_F^2 f_{D^+}^2}{8\pi} |V_{cd}|^2 m_l^2 m_{D^+} \left(1 - \frac{m_l^2}{m_{D^+}^2}\right)^2, \quad (1)$$

where  $G_F$  is the Fermi coupling constant. So by measuring the branching ratio of  $D^+ \rightarrow l^+ \nu_l$ , the decay constant  $f_{D^+}$  can be determined to test the calculations from lattice QCD. In addition,  $|V_{cd}|$  can be directly accessed.  $f_{D^+}$  is also helpful to constrain the CKM matrix element  $|V_{td}|$  through its relation to the mixing parameter  $x_B$  and the ratio  $f_D/f_B$  [3, 4].

To determine the branching ratio of  $D^+ \rightarrow \mu^+ \nu_\mu$ , the  $D$  tagging method is used. Firstly,  $D^-$  mesons are reconstructed with 9 hadronic decay modes ( $D^- \rightarrow K^+ \pi^- \pi^-, K_S^0 \pi^-, K_S^0 K^-, K^+ K^- \pi^-, K^+ \pi^- \pi^- \pi^0, \pi^+ \pi^- \pi^-, K_S^0 \pi^- \pi^0, K^+ \pi^- \pi^- \pi^- \pi^+$  and  $K_S^0 \pi^- \pi^- \pi^+$ ), and then  $\mu^+$  is sought out in the remaining tracks.  $M_{miss}^2 = (E_{beam} - E_{D^-} - E_{\mu^+})^2 - (\vec{p}_{D^-} - \vec{p}_{\mu^+})^2$  is used to extract the signal events which peak

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around 0. Finally, with  $2.9 \text{ fb}^{-1}$   $\psi(3770)$  data, the branching fraction of  $D^+ \rightarrow \mu^+ \nu_\mu$  is measured to be [5]

$$\mathcal{B}(D^+ \rightarrow \mu^+ \nu_\mu) = (3.71 \pm 0.19 \pm 0.06) \times 10^{-4}.$$

With combination of  $G_F$ , the mass of  $\mu^+$  and  $D^+$ , the lifetime of  $D^+$  and  $|V_{cd}|$  from the global fit in the SM [6], the decay constant  $f_{D^+}$  is obtained with the best precision in the world [5]

$$f_{D^+} = (203.2 \pm 5.3 \pm 1.8) \text{ MeV}.$$

### 3 Measurement of $\mathcal{B}(D^0 \rightarrow K^-/\pi^- e^+ \nu)$ and the Form-Factor $f_+(q^2)$

The differential decay rate of  $D^0 \rightarrow K^-(\pi^-)e^+ \nu$  is written as

$$\frac{\Delta\Gamma(D \rightarrow K(\pi)e\nu)}{dq^2} = \frac{G_F^2 |V_{cs(d)}|^2}{24\pi^3} p^3 |f_+(q^2)|^2. \quad (2)$$

With the precisely determined  $V_{cs(d)}$  from the unitarity of the CKM matrix, we can obtain the form-factor  $f_+(q^2)$  to check the lattice QCD calculations.

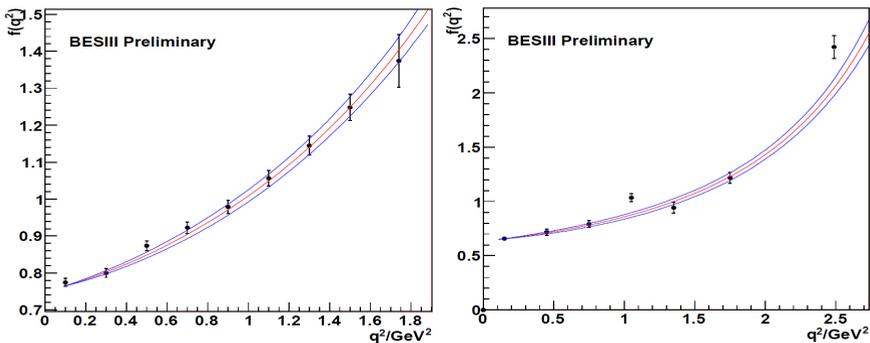
This analysis has been done using one-third of the  $2.9 \text{ fb}^{-1}$   $\psi(3770)$  data, with a partial-blind method. Firstly, the singly tagged  $\bar{D}^0$  is reconstructed with four hadronic modes ( $\bar{D}^0 \rightarrow K^+ \pi^-, K^+ \pi^- \pi^0, K^+ \pi^- \pi^0 \pi^0$  and  $K^+ \pi^- \pi^- \pi^+$ ). The signal candidates are searched by reconstructing two opposite charged tracks ( $K^+/\pi^+ e^-$ ) left in the event, and  $U_{miss} = E_{miss} - p_{miss}$  is used to extract the signal events which peak around 0. Here,  $E_{miss} = E_{beam} - E_{hadron} - E_{electron}$ , and  $p_{miss} = |-\vec{p}_{hadron} - \vec{p}_{electron}|$ . The branching fractions are measured to be

$$\mathcal{B}(D^0 \rightarrow K^- e^+ \nu) = (3.542 \pm 0.030 \pm 0.067)\%$$

$$\mathcal{B}(D^0 \rightarrow \pi^- e^+ \nu) = (0.288 \pm 0.008 \pm 0.005)\%.$$

These two results are consistent with the world averages [6].

To measure the form factor, the partial decay rates are measured in different  $q^2$  bins, where  $q^2$  is the squared invariant mass of  $e^+ \nu$  system. The  $f_+(q^2)$  distributions are shown in figure 1 with theoretical curves overlaid [7].



**Figure 1.**  $f_+(q^2)$  distributions for the decay of  $D^0 \rightarrow K^- e^+ \nu$  (left) and  $D^0 \rightarrow \pi^- e^+ \nu$  (right). Points are measured from data, curves are the theoretical predictions varying within one statistical standard sigma.

## 4 Measurement of the strong phase difference $\delta_{K\pi}$

Studying the  $D^0$ - $\bar{D}^0$  mixing is important for identifying the size of the long distance effect and searching for new physics [8], and also helpful for studying CP violation in charm physics. Charm mixing is described by two dimensionless parameters  $x$  and  $y$ . The measurement of the strong phase difference  $\delta$  between the doubly Cabibbo-suppressed (DCS) decay  $\bar{D}^0 \rightarrow K^- \pi^+$  and the corresponding Cabibbo-favored (CF)  $D^0 \rightarrow K^- \pi^+$ , allows  $x$  and  $y$  to be extracted from  $x' \equiv x \cos \delta + y \sin \delta$  and  $y' \equiv y \cos \delta - x \sin \delta$  [9]. Furthermore, finer precision of  $\delta$  helps the  $\gamma/\phi_3$  angle measurement in CKM matrix according to the so-called ADS method [6].

Using the quantum-correlated technique,  $\delta$  can be accessed using the following formula

$$2r \cos \delta + y = (1 + R_{\text{WS}}) \cdot \mathcal{A}_{CP \rightarrow K\pi}, \quad (3)$$

where  $\mathcal{A}_{CP \rightarrow K\pi}$  is the asymmetry between the branching fractions of CP-odd and CP-even states decaying to  $K^- \pi^+$ .  $\mathcal{B}(D^{CP^\pm} \rightarrow K^- \pi^+)$  are measured with  $D$  tagging method using 5 CP-even  $D^0$  decay modes ( $K^+ K^-$ ,  $\pi^+ \pi^-$ ,  $K_S^0 \pi^0 \pi^0$ ,  $\pi^0 \pi^0$  and  $\rho^0 \pi^0$ ) and 3 CP-odd modes ( $K_S^0 \pi^0$ ,  $K_S^0 \eta$  and  $K_S^0 \omega$ ). With  $2.9 \text{ fb}^{-1}$   $\psi(3770)$  data, we obtain the asymmetry to be [10]

$$\mathcal{A}_{CP \rightarrow K\pi} = (12.7 \pm 1.3 \pm 0.7)\%.$$

By quoting the external inputs of  $R_D = r^2 = (3.50 \pm 0.04)\%$ ,  $y = (6.7 \pm 0.9)\%$  from HFAG 2013 [11] and  $R_{\text{WS}} = (3.80 \pm 0.05)\%$  from PDG [6], we obtain [10]

$$\cos \delta_{K\pi} = 1.02 \pm 0.11 \pm 0.06 \pm 0.01,$$

where the third uncertainty is due to the errors introduced by the external input parameters. This result provides the world best constrain to  $\delta_{K\pi}$ .

## 5 Measurement of the Relative Strong-phase Difference between $D^0$ and $\bar{D}^0$ decay to $K_S^0 \pi^+ \pi^-$

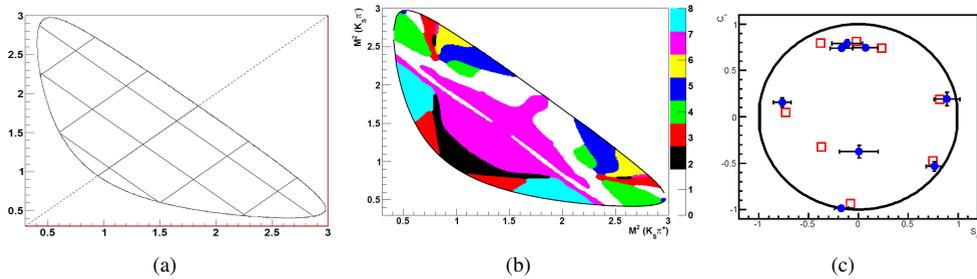
The CKM angle  $\gamma/\phi_3$  can be measured from the decay  $B^+ \rightarrow D(K_S^0 \pi^+ \pi^-) K^\pm$  through the GGSZ method [12, 13]. The Dalitz plot can be separated into  $2N$  bins as shown in figure 2(a). The  $N$  bins lying below the symmetry axis are denoted by the index  $i$ , while the remaining bins are indexed with  $\bar{i}$ , and these square bins are symmetric under exchange of  $x$  and  $y$  ( $x \equiv m_{K_S^0 \pi^-}^2$ ,  $y \equiv m_{K_S^0 \pi^+}^2$ ). The two parameters  $c_i$  and  $s_i$  denote the weighted average of  $\cos(\Delta\delta_D)$  and  $\sin(\Delta\delta_D)$  in the  $i^{\text{th}}$  bin, respectively. And  $\Delta\delta_D$  is the strong phase difference between the DCS decay  $\bar{D}^0 \rightarrow K_S^0 \pi^+ \pi^-$  and the corresponding CF decay  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ .

The number of events in the  $i^{\text{th}}$  bin of a Dalitz plot originating from the  $B^\pm$  decay is

$$N_i^\pm = \frac{a_B}{a_D} [K_i + r_B^2 K_{\bar{i}} \pm 2r_B \sqrt{K_i K_{\bar{i}}} (\cos(\delta_B + \gamma) c_i + \sin(\delta_B + \gamma) s_i)]. \quad (4)$$

The normalization factors,  $r_B$  and  $\delta_B$  could be determined at  $B$  factories. Using  $\psi(3770)$ ,  $K_i$  can be obtained by flavor tagged  $D \rightarrow K_S^0 \pi^+ \pi^-$  decays,  $c_i$  can be obtained from CP tagged  $D \rightarrow K_S^0 \pi^+ \pi^-$  decays, and  $s_i$  is obtained from double Dalitz decays  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  and  $\bar{D}^0 \rightarrow K_S^0 \pi^+ \pi^-$ .

The Dalitz plot is split into 8 bins as shown in figure 2(b), which have been optimized, to get the highest sensitivity of  $\gamma$ . The results of  $c_i$  and  $s_i$  in each bin are given in figure 2(c). Our results will allow for increased precision in the measurement of  $\gamma/\phi_3$ .



**Figure 2.** (a): Squared binned Dalitz plot with symmetric bins over an exchange of  $x$  and  $y$ . (b): Modified optimal binning. (c): Red squares mark the model predicted value of  $c_i$  and  $s_i$  [14], and the blue dot mark fitted value from data with statistical error.

## 6 Summary

Based on the  $2.9 \text{ fb}^{-1}$  data taken at the  $\psi(3770)$  peak, the recent results on the study of charm physics are presented. The BESIII experiment aims at getting  $20 \text{ fb}^{-1}$   $\psi(3770)$ , and more promising results are expected.

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