

Evolution of heavy quark distribution function on quark-gluon plasma: Using the Iterative Laplace Transform Method

Sharareh Mehrabi Pari^{1,2*}, Kurosh Javidan^{1**}, Fatemeh Taghavi Shahri^{1***}

¹ Department of Physics, Ferdowsi University of Mashhad, 91775-1436
Mashhad, Iran

² Department of Physics and Technology, University of Bergen, 5007
Bergen, Norway

Abstract

The “Iterative Laplace Transform Method” is used to solve the Fokker-Plank equation for finding the time evolution of the heavy quarks distribution functions such as charm and bottom in quark gluon plasma. These solutions will lead us to calculation of nuclear suppression factor R_{AA} . The results have good agreement with available experiment data from the PHENIX collaboration.

1 Introduction

Heavy partons, which are produced in the early stage of heavy ion collisions could induce a large amount of energy loss for hard partons produced in the initial stage of the collision. So, the observed jet quenching in such collisions can indicate the formation of this strongly interacting dense matter [1, 2]. Charm and bottom quarks play crucial role in this endeavor because they are

*Email: sharareh.mehrabi.pari@gmail.com

**Email: Javidan@um.ac.ir

***Email: taghavishahri@um.ac.ir

the witness to the entire space-time evolution of the system and due to their large masses, a memory of their interaction history may be preserved [3].

The Fokker-Plank equation (FPE) provides an appropriate framework for studying the evolution of the heavy quark distribution functions [4].

Heavy quarks dissipate their energy during their propagation through QGP via two processes [5]: The first one is the collisional process and the second one is the radiative process [6]. These processes can be described well by the Fokker-Plank equation. In this paper our main task is to study the evolution of heavy quarks in quark gluon plasma by solving the Fokker-Plank equation with a new method named "Iterative Laplace Transform Method" [7]. This method gives numerical solutions in the form of convergent series. To verify our solutions, we calculate the nuclear suppression factor, $R_{AA}(p_T)$, of heavy quarks using simulated data of heavy ion collisions [8]. To do this we consider both the collisional and radiation parts of energy loss in FPE and finally we compare our results with experimental data from PHENIX [9].

2 The heavy quark momentum evolution in FPE dynamics

The Quark Gluon Plasma (QGP) is formed at the time τ_i after the impact. Time evolution of the momentum distribution of non-equilibrated heavy quarks due to the interaction with the equilibrated QGP during the time interval $\tau_i < \tau < \tau_{HQ}$ can be calculated by the Fokker-Plank equation as follow:

$$\frac{\partial f(p, t)}{\partial t} = A(p) \frac{\partial}{\partial p} (pf(p, t)) + D(p) \left[\frac{\partial}{\partial p} \right]^2 f(p, t) \quad . \quad (1)$$

Where A is drag and D is diffusion coefficients. The drag coefficient is an important quantity carrying information about the dynamics of elastic collisions. It is expected that the drag coefficient should be determined by the properties of the bath and not by the nature of the HQ. In this work we use the drag coefficient as follows:

$$A(p, t) = -\frac{1}{p} \frac{dE}{dL} \quad . \quad (2)$$

If coupling between HQs and background QGP is weak, then the diffusion coefficient can be calculated using the Einstein relation, $D(p) = ET A(p)$, where M is the mass of the heavy quark and T is the thermal bath temperature [10]. In our calculations, the collisional energy loss of HQs in QGP by

considering "Hard and Soft Thermal Loops" are given as follows [11]. For $E \ll \frac{M^2}{T}$:

$$-\frac{dE}{dL} = \frac{8\pi\alpha_s^2 T^2}{3} \left(1 + \frac{n_f}{6}\right) \left[\frac{1}{v} - \frac{1-v^2}{2v^2} \ln\left(\frac{1+v}{1-v}\right)\right] \ln\left[2^{\frac{n_f}{6+n_f}} B(v) \frac{ET}{m_g M}\right], \quad (3)$$

for $E \gg \frac{M^2}{T}$:

$$-\frac{dE}{dL} = \frac{8\pi\alpha_s^2 T^2}{3} \left(1 + \frac{n_f}{6}\right) \ln\left[2^{\frac{1}{2}(\frac{n_f}{6+n_f})} 0.92 \frac{\sqrt{ET}}{m_g}\right], \quad (4)$$

where v is the HQ speed, α_s is the strong coupling constant, n_f is the number of quark flavors in the medium, E and M are energy and mass of the propagated HQ respectively, $m_g = \sqrt{(1+n_f/6)g^2 T^2/3}$ is the thermal gluon mass, $g = \sqrt{4\pi\alpha_s}$ is the gauge coupling parameter and $B(v)$ is a smooth velocity function, which can be taken approximately as 0.7 [11]. Finally the total energy loss due to both collision and radiation processes is:

$$-\frac{dE_{total}}{dL} = \frac{10\alpha_s^2 T^2}{3} \ln\left(\frac{3E}{8\pi\alpha_s T}\right) \left[\ln\left(\xi + \sqrt{1+\xi^2}\right) + \xi \ln\left(\frac{1}{\xi} + \sqrt{1+\frac{1}{\xi^2}}\right) + \frac{3}{8\pi}\right] \quad (5)$$

where $\xi = \frac{9E}{2\pi^3 T}$. Using equation (4), (5), we will calculate the drag and diffusion coefficients for HQ energy loss during the interaction with QGP medium. In the next section we will find the solution of the Fokker-plank equation, eq(1).

3 The Iterative Laplace Transform Method

In this section we want to introduce a new method for solving the FP equation, namely the "Iterative Laplace Transform Method" (ILTM) [7]. This new method was proposed by Daftardar-Gejji and Jafari to seek numerical solutions of nonlinear functional equations. This method is called the Iterative Laplace Transform Method, which is a combination of two powerful methods, namely, the Laplace transform method and the Iterative method. The method gives numerical solutions in the form of convergent series with easily computable components. The most outstanding feature of this method is that it provides an analytical solution using the initial conditions only, without any discretization or restrictive assumptions. In general, the Fokker-Plank equation can be written as [7]:

$$D_t^\alpha f = \left[D_p^\beta A(p, t, f) + D_p^{2\beta} B(p, t, f) \right] f(p, t) \quad . \quad (6)$$

Here $D_t^\alpha(\cdot)$, $D_p^\beta(\cdot)$, $D_p^{2\beta}(\cdot)$ are the Caputo fractional derivative with respect to t and p .

In our calculations we use the following definitions:

I. The Caputo fractional derivative of function (p, t) is defined as

$$D_t^\alpha f(p, t) = \frac{1}{\Gamma(m - \alpha)} \int (t - \eta)^{m-\alpha-1} f^m(p, \eta) d\eta \quad (7)$$

for $m - 1 < \alpha \leq m, m \in N$.

II. Laplace transform of $D_t^\alpha f(p, t)$ is given by

$$\mathcal{L}[D_t^\alpha f(p, t)] = x^\alpha \mathcal{L}[f(p, t)] - \sum_{k=0}^{m-1} f^{(k)}(p, 0) x^{\alpha-1-k}; \quad m - 1 < \alpha \leq m, \quad (8)$$

where $f^{(k)}(p, 0)$ is the k th derivative of $f(p, t)$ at $t = 0$.

To illustrate how the iterative Laplace transform method works, the general space-time fractional partial differential equation is considered

$$D_t^\alpha f = \mathcal{A}(f, D_p^\beta f, D_p^{2\beta} f, \dots) \quad m - 1 < \alpha \leq m \\ n - 1 < \beta \leq n \quad m, n \in N, \quad (9)$$

where $\mathcal{A}(f, D_p^\beta f, D_p^{2\beta} f, \dots)$ is a linear or nonlinear operator of $f, D_p^\beta f, D_p^{2\beta} f, \dots$

With the initial condition

$$f^{(k)}(p, 0) = h_k(p); \quad k = 0, 1, \dots, m - 1, \quad (10)$$

and $f(p, t)$ will be determined later.

First of all we take the Laplace transform of both sides of (9)

$$x^\alpha \mathcal{L}[f(p, t)] - \sum_{k=0}^{m-1} x^{\alpha-1-k} f^{(k)}(p, 0) = \mathcal{L}[\mathcal{A}(f, D_p^\beta f, D_p^{2\beta} f, \dots)]. \quad (11)$$

After simplification, we have

$$\mathcal{L}[f(p, t)] = \sum_{k=0}^{m-1} x^{-1-k} f^{(k)}(p, 0) + x^{-\alpha} \mathcal{L}[\mathcal{A}(f, D_p^\beta f, D_p^{2\beta} f, \dots)]. \quad (12)$$

By inverse Laplace transforms of both sides of Eq. (12) we have:

$$f(p, t) = \mathcal{L}^{-1} \left[\sum_{k=0}^{m-1} x^{-1-k} f^{(k)}(p, 0) \right] + \mathcal{L}^{-1} \left[x^{-\alpha} \mathcal{L}[\mathcal{A}(f, D_p^\beta f, D_p^{2\beta} f, \dots)] \right], \quad (13)$$

Then we set

$$f_0 = \mathcal{L}^{-1} \left[\sum_{k=0}^{m-1} x^{-1-k} f^{(k)}(p, 0) \right],$$

$$f_1 = \mathcal{L}^{-1} \left[x^{-\alpha} \mathcal{L}[\mathcal{A}(f, D_p^\beta f, D_p^{2\beta} f, \dots)] \right]$$

If we want to calculate the next term in the series we have to consider f_1 as initial condition and putting it in the equation (9) and repeat all the above steps. Therefore the solution of (9) with initial condition (10) is given by:

$$f(p, t) \cong f_0(p, t) + f_1(p, t) + \dots + f_m(p, t), m = 1, 2, \dots \quad (14)$$

If we set $\alpha = \beta = 1$ in Equation (9) we will reach the FPE. We solved the FPE using both finite difference method and ILTM with different types of initial conditions and compared their results. For initial conditions with non vanishing values at the borderer, the ILTM successfully solves the FPE while results of finite difference method diverges and therefore it fails to find the solution. We have taken transverse momentum distribution functions of charm and bottom quarks at $\sqrt{s} = 200\text{GeV}$ in Au-Au and also P-P collisions from [12]. ILTM has been performed up to $m = 2$.

4 Calculating the nuclear suppression factor $R_{AA}(P_T)$

The dissipation coefficients of the thermal bath and initial momentum distribution of HQs, are basic inputs required to solve the FPE. In our calculation the time dependence of temperature is taken from [13]. Here we considered a fixed value for coupling constant as $\alpha_s = 0.3$.

Figure 1 demonstrates the drag coefficient for b and c quarks due to collision and radiation separately. This figure clearly shows that the energy loss for c quark is greater than that for b quark.

One of the most important experimental observables to quantify the nuclear medium is the depletion of high p_T particles (D and B mesons

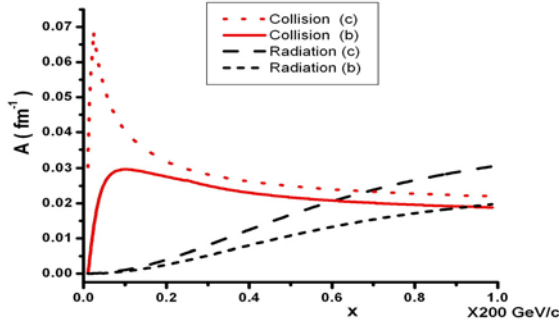


Figure 1: The drag coefficient due to collision and radiation as functions of momentum for Bottom and Charm quarks in QGP.

or non-photonic single e^- spectra resulting from semi leptonic decays of hadrons containing charm and bottom quarks) produced in Nucleus-Nucleus collisions with respect to those produced in proton-proton collisions which is formulated in nuclear suppression factor R_{AA} .

Here the nuclear modification factor in relativistic heavy ion collisions is calculated for demonstrating the power of ILTM applied to solve the FP equation. The solution of the FP equation have been convoluted with the fragmentation functions of the heavy quarks to obtain the P_T distribution of the heavy mesons [13]. Similar simulations also have been performed to find

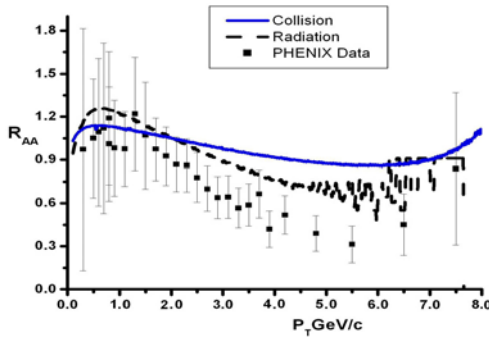


Figure 2: Nuclear suppression factor R_{AA} due to collisional and radiation energy loss separately as functions of P_T . Used parameters are: $\tau_0 = 0.33$ fm/c, $T_0 = 0.375$ GeV, $\alpha_s = 0.3$, $T_c = 0.175$ GeV [13].

the electron spectrum created in the P-P collisions. For this purpose HQ

distribution functions of proton at $\sqrt{s} = 200\text{GeV}$ have been taken as initial conditions to solve FPE again. The ratio of these two quantities is proportional to the nuclear suppression factor, R_{AA} measured in experiments:

$$R_{AA}(P_T) \propto \frac{\frac{dN^e}{dP_T d\eta} \text{ Au-Au}}{\frac{dN^e}{dP_T d\eta} \text{ P-P}} \quad (15)$$

A small value of R_{AA} indicates a strong suppression and therefore, large energy loss of heavy quarks in the medium. It is clear that this ratio is equal to one in the absence of any medium.

Figure 2 presents contribution of collision and radiation energy loss in suppression factor R_{AA} separately. Displayed experimental data has been obtained from the PHENIX collaborations for $Au+Au$ collisions at $\sqrt{s_{NN}} = 200\text{GeV}$. This figure shows that suppression at higher values of P_T can not be explained by collisional energy loss. In other words, contribution of radiation in energy loss at higher P_T is more important than collisional energy loss only.

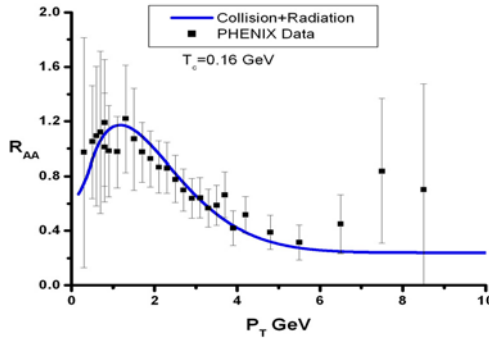


Figure 3: Nuclear suppression factor R_{AA} as function of P_T . Energy loss due to collision and radiation have been considered in calculating the drag coefficient.

5 Conclusions and remarks

The “Iterative Laplace Transform Method” has been introduced as an effective method to solve the Fokker-Plank equation for initial conditions with non-zero values at the boundaries. Our numerical algorithm shows that this method is able to solve time evolution of distribution function of HQs successfully. To verify the ability of ILTM and demonstrating its application,

we calculated the nuclear suppression factor R_{AA} . It is shown that there is a good agreement between simulation results and reported experimental data (fig. 3).

References

- [1] W. Scheid, M. Muller and W. Greiner, PRL **32**, 741 (1974). G. F. Chapline, M. H. Johnson, E. Teller, M. S. Weiss PRD **8**, 135 (1973). L. P. Csernai, et al., PRC **25**, 2482 (1982).
- [2] J.Y. Ollitrault, Eur. J. Phys. **29**, 275 (2008).
- [3] B. Svetitsky Phys. Rev. D **37**, 2484 (1988).
- [4] M. L. Mangano, P. Nason and G. Ridolfi, Nucl. Phys. B **538**, 282 (2002).
- [5] S. Mazumder, T. Bhattacharyya, J. Alam and S.K. Das, Phys. Rev. C **84**, 044901 (2011).
- [6] B. Z. Kopeliovich, I. K. Potashnikova, I. Schmidt, Phys. Rev. C **82**, 037901 (2010).
- [7] L. Yan, Hindawi Publishing Corporation, Volume 2013, Article ID 465160.
- [8] S.K. Das, F. Scardina, S. Plumari, and V. Greco, Phys. Rev. C **90**, 044901 (2014).
- [9] S. S. Adler et al., (PHENIX Collaboration), Phys. Rev. C **84**, 044905 (2011).
- [10] A.K. Chaudhuri, V.P. Goncalves, and L. F. Mackedanz, Brazilian Journal of Physics, 37,no 2B (2007).
- [11] M. B. Gay Ducati, V. P. Goncalves and L. F. Mackedanz, Brazilian Journal of Physics, 37, 656-660 (2007).
- [12] M. Cacciari, P. Nason, R. Vogt, PRL 95, 122001 (2005)
- [13] S. K. Das, J. Alam, and P. Mohanty, Phys. Rev. C **82**, 014908 (2010).