

The 4-parameter Compressible Packing Model (CPM) including a critical cavity size ratio

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Abstract. The 4-parameter Compressible Packing Model (CPM) has been developed to predict the packing density of mixtures constituted by bidisperse spherical particles. The four parameters are: the wall effect and the loosening effect coefficients, the compaction index and a critical cavity size ratio. The two geometrical interactions have been studied theoretically on the basis of a spherical cell centered on a secondary class bead. For the loosening effect, a critical cavity size ratio, below which a fine particle can be inserted into a small cavity created by touching coarser particles, is introduced. This is the only parameter which requires adaptation to extend the model to other types of particles. The 4-parameter CPM demonstrates its efficiency on frictionless glass beads (300 values), spherical particles numerically simulated (20 values), round natural particles (125 values) and crushed particles (335 values) with correlation coefficients equal to respectively 99.0%, 98.7%, 97.8%, 96.4% and mean deviations equal to respectively 0.007, 0.006, 0.007, 0.010.

1 Introduction

The packing density is often an optimized physical value for high performing and sustainable materials. Packing density models call upon to geometrical interaction coefficient functions of the size ratio: the *wall effect* b and the *loosening effect* a [1-7] to which can be added [2-3, 5] a *wedging effect* c . Because the packing density depends on the packing process efficiency and on the particle shape, two strategies can be adopted.

The first strategy [2-3, 5] consists in taking that into account in the mathematical expressions used for the interaction coefficients. It means that an adjustment by regression analysis is each time necessary to find the best functions of these parameters. Three sets have been proposed by the authors: the first one for the spherical particles, the second and the third ones for respectively the compacted and uncompacted angular particles. Recently, the interaction function for the *wedging effect* c has been expressed in two separated parts [5].

The second strategy [4, 7] consists in taking into account of the packing process efficiency by a *compaction index* K and of the particle shape and of the surface texture through a *critical cavity size ratio* x_0 . Below this value, a small particle can be inserted into a small cavity created by touching coarser particles without disturbed their arrangement. Beyond this value, the local volume of the coarse class decreases: a *loosening effect* occurs. Thus, the *4-parameter Compressible Packing Model (CPM)* [4, 7] has been developed by incorporating an original theory about the *wall effect* b and the *loosening effect* a . The 4 parameters are: a , b , K and x_0 . Even if the expressions of a and b are more sophisticated, the adjustment of K and x_0 has proven effective for spherical

particles for which $x_0 = 0.2$, in harmony with the tetrahedral cavern theory ($x_0 \approx 0.224$). It remains to be able to adjust easily x_0 for other types of particles, especially angular ones.

2 Wall effect, loosening effect, interference effect

Fig. 1 represents the "void ratio" ($e=1/\phi^*-1$) of a binary mixture vs the "volume fraction of fine particles". In the "*coarse dominant*" field, the *insertion mechanism* is represented by AM and, in the "*fine dominant*" field, the *substitution mechanism* by FM. The *wall effect* is highlighted by FH which is located above FO. In the AK section, the *loosening effect* is only localized around a particle insufficiently fine to insert into a cavity created by some touching coarser particles. The more global *interference effect* is materialized by the AKJI section.

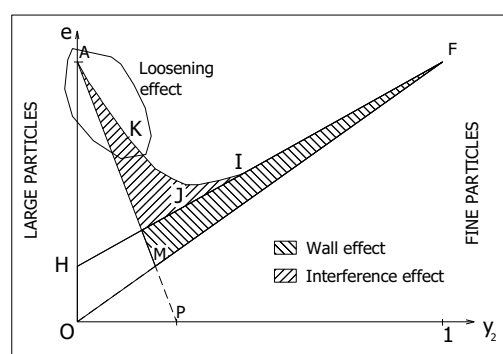


Fig. 1. Illustration of wall effect, loosening effect, interference effect.

3 The 4-parameter CPM

The reference frame is constituted by the CPM [1]. The real packing density corresponds to that obtained for a randomly dense packing. It depends on the process used to fill and to compact the material inside the mould. If we consider a perfect placing process where each particle is placed one by one in its ideal location, the *compaction index* K tends to infinity and the packing density reaches the virtual packing density. For a binary mixture "with geometrical interactions", the virtual packing density γ obtained by the CPM can be written:

$$\gamma = \text{Min} \left(\gamma_1 = \frac{\beta_1}{1 - \left(1 - \frac{\beta_1}{\beta_2}\right) a_{12}} y_2 ; \gamma_2 = \frac{\beta_2}{1 - \left(1 - \beta_2 + b_{21} \beta_2 \left(1 - \frac{1}{\beta_1}\right)\right) y_1} \right) \quad (1)$$

where γ_1 and γ_2 are respectively the virtual packing densities of the binary mixture in the "coarse dominant" field and in the "fine dominant" field, β_1 and β_2 the virtual packing densities of the coarse size class 1 and of the fine size class 2, y_1 and y_2 their volume fractions by reference of the total solid volume, a_{12} the *loosening effect* coefficient, b_{21} the *wall effect* coefficient. When the size ratio $x = d_2/d_1$ (fine/coarse) is equal to 1, a total interaction occurs:

$$b_{21}(1) = 1 \text{ and } a_{12}(1) = 1 \quad (2)$$

The real packing density ϕ^* is calculated by introducing the *compaction index* K . The expression (3) is an implicit equation of ϕ^* , with one and only one positive root.

$$K = \frac{y_1}{\frac{1}{\phi^*} - \frac{1}{\gamma_1}} + \frac{y_2}{\frac{1}{\phi^*} - \frac{1}{\gamma_2}} \quad (3)$$

3.1. Wall effect theory for spheres

The *wall effect* is the further amount of voids in the packing of small particles against the wall of a coarser one. The disturbance is studied in a portion delimited by two concentric spheres. The first one has a diameter d_1 . The second hypothetical one is chosen in such a way that the boundary condition $b_{21}(1)=1$ is respected. That's why we call this theory: ROAD (ROund ADjustable cell) theory.

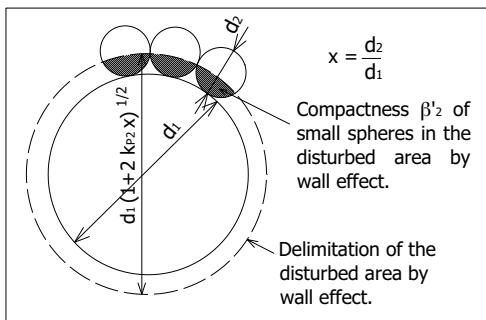


Fig. 2. Definition of the spherical reference cell for studying the wall effect.

The diameter of the sphere delimiting the disturbed volume by the *wall effect* around a coarse sphere is:

$$d_{hyp} = d_1 \sqrt{1 + 2 k_{p2} x} \quad (4)$$

Let α_2 be the real packing density of the fine class 2. We can deduce its virtual packing density β_2 by involving K :

$$\beta_2 = \alpha_2 \frac{(1 + K)}{K} \quad (5)$$

The number of fine spheres against a coarse one is calculated from the Spherical Square Model (SSM), not presented here, as part of the dense virtual packings:

$$N_{12,SSM}^{dense}(x) = \frac{\pi(1+x)}{x \arcsin\left(\frac{x}{1+x}\right)} \quad (6)$$

The packing density β'_2 of the fine particles in the disturbed volume by the *wall effect* is then deduced:

$$\beta'_2(x) = \frac{\pi(1+x)}{4x \left((1+2k_{p2}x)^{\frac{3}{2}} - 1 \right) \arcsin\left(\frac{x}{1+x}\right)} \times \left(2(1+2k_{p2}x)^{\frac{3}{2}} - 3(1+2k_{p2}x) \left(1 + \frac{k_{p2}x}{1+x}\right) + \left(1 + \frac{k_{p2}x}{1+x}\right)^3 - x^3 + \frac{3k_{p2}x^3}{1+x} + x^3 \left(1 - \frac{k_{p2}}{1+x}\right)^3 \right) \quad (7)$$

In the case where $\beta_1 = \beta_2$, the *wall effect* coefficient is:

$$b_{21}(x) = \frac{(\beta_2 - \beta'_2(x)) \left[(1+2k_{p2}x)^{\frac{3}{2}} - 1 \right]}{(1-\beta_2)} \quad (8)$$

It remains to be determined k_{p2} by respecting the boundary condition $b_{21}(1)=1$, leading to the numerical solution, of the following equation:

$$(\beta_2 - 6) \left(1 + 2k_{p2}\right)^{\frac{3}{2}} + \frac{9}{2} k_{p2}^2 + 18k_{p2} + 5 = 0 \quad (9)$$

The equation to be solved being of the 3rd degree, the value of k_{p2} which is coherent with those presented in the following table should be kept:

Table 1. Examples of values of k_p as a function of the virtual packing density β .

β	0.65	0.68	0.70	0.72	0.734
k_p	0.4466	0.7543	0.9253	1.1476	1.4729

An example of the *wall effect* coefficient b_{21} as a function of the size ratio x is given Fig. 3. The results from the ROAD theory have been achieved for a *compaction index* $K=9$ corresponding to a packing process "vibration + compression" (Table 2) and a real packing density of the fine class 2 $\alpha_2=0.641$ corresponding to the *Random Close Packing RCP* state.

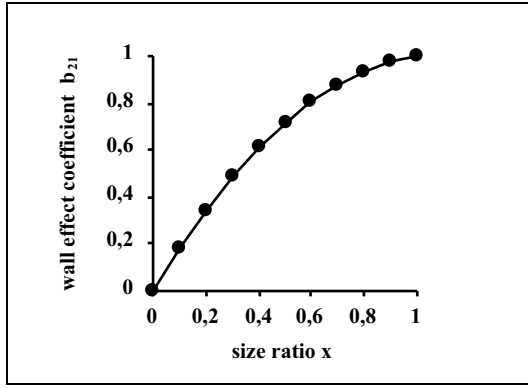


Fig. 3. ROAD theory: $\alpha_2=0.641$ – Vibration + compression.

3.2 Loosening effect and interference effect theory for spheres

The concept of *critical cavity size ratio* x_0 is now introduced: below this value, the intrusion of small quantities of finer spheres does not disturb the bed of coarser spheres; beyond this value, the fine spheres cannot be placed without disturbing this one. The concept of *local isotropic expansion* of the coarse particle skeleton around a small one is adopted when $x > x_0$. To be consistent with the *wall effect* theory, a spherical reference cell is used.

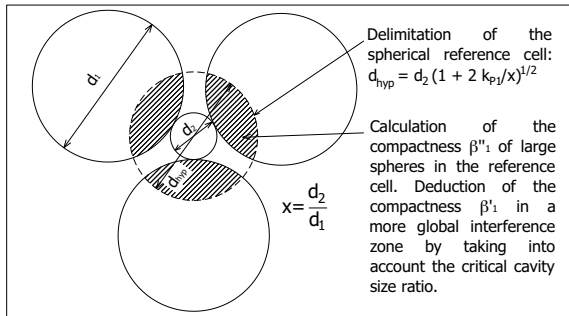


Fig. 4. Definition of the spherical reference cell for studying the loosening effect.

The diameter of the sphere delimiting the disturbed volume by the *loosening effect* is:

$$d_{hyp} = d_2 \sqrt{1 + \frac{2k_{p1}}{x}} \tag{10}$$

Let α_1 be the real packing density of the large class 1. We can deduce its virtual packing density β_1 :

$$\beta_1 = \alpha_1 \frac{(1+K)}{K} \tag{11}$$

The number of coarse spheres against a fine one is calculated from the Spherical Square Model (SSM):

$$N_{21,SSM}^{dense}(x) = \frac{\pi(1+x)}{\arcsin\left(\frac{1}{1+x}\right)} \tag{12}$$

The packing density β_1'' of the coarse spheres in the reference cell is:

$$\beta_1''(x) = \frac{\pi(1+x)}{4\left(1 + \frac{2k_{p1}}{x}\right)^{\frac{3}{2}} \arcsin\left(\frac{1}{1+x}\right)} \times \left(2\left(1 + \frac{2k_{p1}}{x}\right)^{\frac{3}{2}} - 3\left(1 + \frac{2k_{p1}}{x}\right)\left(1 + \frac{k_{p1}}{1+x}\right) + \left(1 + \frac{k_{p1}}{1+x}\right)^3 - \frac{1}{x^3} + \frac{3k_{p1}}{x^2(1+x)} + \left(\frac{1}{x} - \frac{k_{p1}}{1+x}\right)^3\right) \tag{13}$$

The packing density of the large spheres in the more global interference zone is:

$$\beta_1'(x) = \frac{\beta_1''(x)}{\beta_1''(x_0) \left(1 + \frac{(x-x_0)}{(1-x_0)} \left(\sqrt[3]{\frac{2\beta_1''(1)}{\beta_1''(x_0)}} - 1\right)\right)^3} \beta_1 \tag{14}$$

When $x=x_0$, $\beta_1'(x_0)/\beta_1=1$: the *loosening effect* and the *interference effect* do not occur. When $x=1$, $\beta_1'(1)/\beta_1=0.5$: the continuity between the "coarse dominant" field and the "fine dominant" field is provided when the volume fractions of each class are equal and when all particles become identical. The volume fraction of the fine particles at the "eutectic" point, corresponding to the crossing point into the "fine dominant" field, is:

$$\phi_2^*(x) = \beta_2 + ((1-\beta_2)(1-b_{21}(x)) - 1)\beta_1'(x) \tag{15}$$

The *loosening effect* coefficient is deduced from this expression:

$$a_{12}(x) = \frac{\beta_1 - \beta_1'(x)}{\phi_2^*(x)} \text{ if } x \geq x_0 \text{ and } a_{12}(x) = 0 \text{ if } x \leq x_0 \tag{16}$$

k_{p1} remains to be determined by respecting the boundary condition $a_{12}(1)=1$ leading to solve the equation (9) by replacing k_{p2} by k_{p1} and β_2 by β_1 .

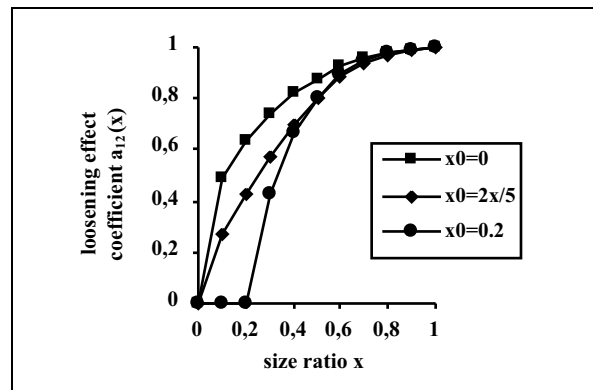


Fig. 5. ROAD theory: $\alpha_1=0.641$ – Vibration + compression.

3.3 Compaction index K and critical cavity size ratio x_0

K (Table 2) and x_0 (Table 3) need now to be calibrated by an analysis on 780 values on binary mixtures.

x_0 depends on the shape and on the finish surface of the materials. For equivalent sizes, the *loosening effect* created by a small particle in a matrix of coarse ones is more important when this particle is irregularly shaped.

Table 2. Compaction index K .

Packing process	Pouring	Vibration or compaction	Vibration + compression	Optimized filling + vibration + compression
K	4.7	5.6	9	15

Table 3. Critical cavity size ratios x_0 (x : size ratio).

Type of particles	Crushed particles			Round natural particles	Frictionless glass beads
	Very Angular / very rough	Angular / rough	Angular		
x_0	0	0.02	0.1	$x_0=2x/5$	0.2

The *4-parameter CPM* has been evaluated [7] by comparison with 300 values of packing densities on frictionless glass beads, 20 on spherical particles numerically simulated, 125 on round natural particles and 335 on crushed particles (Table 4).

Table 4. Mean deviation ξ and correlation coefficient r . Comparison between different models.

Models / Materials (780 values)	Original CPM [1]		3PPM [2, 3]		4-parameter CPM [4, 7]	
	ξ	r	ξ	r	ξ	r
Glass beads	0.012	0.9754	0.009	0.9863	0.007	0.9904
Simulation	0.012	0.8783	0.008	0.9598	0.006	0.9877
Round natural	0.009	0.9619	0.012	0.9534	0.007	0.9788
Crushed	0.013	0.9408	0.013	0.9455	0.010	0.9642

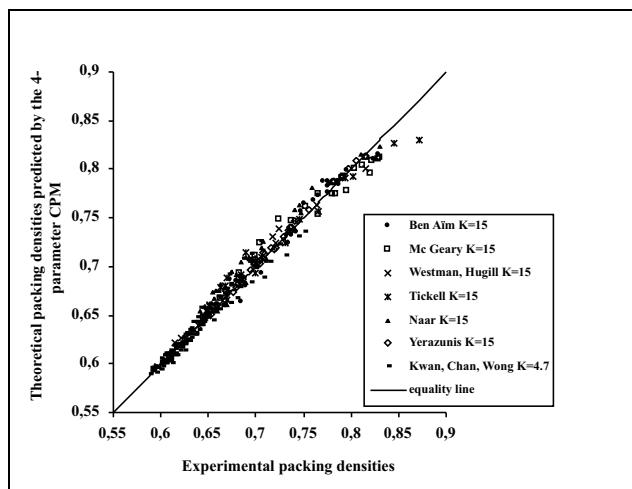


Fig. 6. Theoretical predictions by the 4-parameter CPM for frictionless glass beads with $x_0=0.2$.

The *4-parameter CPM* has been compared with the *original CPM* [1] and with the *3-parameter particle packing model (3PPM)* [2, 3] (Table 4). It provides reliable predictions, as well as for low and high packing densities, especially at the optimum.

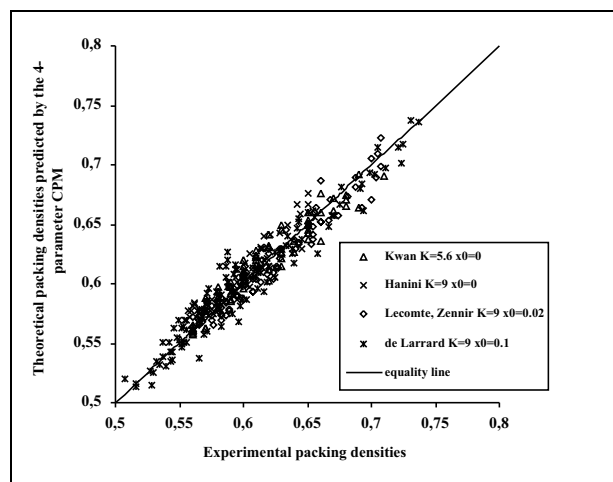


Fig. 7. Theoretical predictions by the 4-parameter CPM for crushed particles.

4 Conclusion

The *4-parameter CPM* offers a new theory on the *wall effect* and on the *loosening effect* studied by considering elementary juxtaposed cells. The effect of a coarse or a fine particle is then examined on the basis of a foreign sphere surrounded by dominant class neighbours. It includes the introduction of a *critical cavity size ratio* depending on the type of materials. This reveals that there would not be a single curve for the geometrical interactions. The limitations of the proposed model are of two types. Firstly, the interaction coefficients are determined by considering the same virtual packing density for each granular class. Secondly, it remains to characterize the shape, the angularity and the surface texture of the particles by reference to test methods. However, the results obtained suggest that efforts will not have been in vain.

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