

Excited lepton baryogenesis

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Abstract. The excited leptons that share the quantum numbers with the Standard Model leptons but have larger masses are widespread in many promising new physics theories. A subclass of excited leptons that at low energies interact with the SM fermions dominantly through the effective coupling to lepton and fermion-antifermion pair can be referred as leptomesons. I introduce possible generation of the baryon asymmetry of the universe using these new particles. The discussed baryogenesis mechanisms do not contradict to the small neutrino masses and the proton stability, and can be interesting for the collider experiments.

1 Introduction

In spite of the present success of the Standard Model (SM) several observations indicate its possible nonfundamentality: large number of the SM fermions, their arbitrary masses and mixings, fractional electric charge of quarks, etc. Among the diversity of new physics models the theories of compositeness [1–8] try to solve these problems by introducing a substructure of the SM particles, which subcomponents are commonly referred as preons [1]. The composite models may include in the particle content radial, orbital, topological, structural, and other excitations of the ground state particles, e.g., an *excited lepton* that shares leptonic quantum number with one of the existing leptons, has larger mass and no color charge.

Besides the outlined issues on the particles and their interactions that come from the laboratory studies, another opened questions (including the dark matter problem) arise from the astrophysical observations of the universe around us. In particular, our universe appears to be populated exclusively with baryonic matter rather than antimatter [9]. However this baryon asymmetry can not be explained within the Big Bang cosmology and the SM. Possible scenarios of dynamical generation of the baryon asymmetry during the evolution of the universe from a hot early matter-antimatter symmetric stage are referred as the baryogenesis (BG) mechanisms [10, 11], and include new physics.

In the next section we discuss one example of the composite models in question. The interactions and the mass bounds for the excited leptons are outlined in section 3, and the new BG scenarios that involve these particles are discussed in section 4. Finally, we conclude in section 5.

Table 1. The haplon quantum numbers

Haplon	Spin [\hbar]	Q [$[e]$]	C1	C2	SU(2) _{<i>h</i>}
α	1/2	+1/2	3	1	2
β	1/2	-1/2	3	1	2
ℓ	0	+1/2	$\bar{3}$	1	2
c_k	0	-1/6	3	3	2

2 Composite model example

Consider the haplon models [7, 12], which are similar to the earlier models with wakems and chroms [4, 13, 14], and are based on the symmetry SU(3)_{*c*} × U(1)_{*em*} × SU(*N*)_{*h*}, where the new haplon group SU(*N*)_{*h*} has the confinement scale of the order of 0.3 TeV, and denotes, e.g., SU(2)_{*L*} × SU(2)_{*R*}. These models contain the two categories of preons (haplons): the fermions $\alpha^{+1/2}$ and $\beta^{-1/2}$, and the scalars $\ell^{+1/2}$ and $c_k^{-1/6}$, where $k = r, g, b$ (“red, green, blue”). Their quantum number assignment is given in Table 1, where Q is the electric charge, C1 is the choice for the SU(3)_{*c*} representations in Ref. [7]¹ and C2 is an alternative choice [12]. Then the haplon pairs can compose the SM particles as $\nu = (\alpha\bar{\ell})$, $e^- = (\beta\bar{\ell})$, $u = (\alpha\bar{c}_k)$, $d = (\beta\bar{c}_k)$, $W^- = (\bar{\alpha}\beta)$, $W^3 = (\bar{\alpha}\alpha - \bar{\beta}\beta)/\sqrt{2}, \dots$, and the new particles, e.g., a scalar leptoquark $S_\ell^{+2/3} = (\ell\bar{c}_k)$, and the neutral scalars $S_\ell^0 = (\ell\bar{\ell})$ and $S_c^0 = (c_k\bar{c}_k)$. W^3 mixes with the photon γ similarly to the mixing between γ and ρ^0 -meson. H scalar can be a p-wave excitation of the Z , and the second and third generations can be dynamical excitations. Notice that $S_\ell^{+2/3}$, S_c^0 and S_ℓ^0 states (if their masses are small) may contribute to the low-energy observables, e.g., so-called, XYZ states [16].²

¹Notice that C1 assignment does not provide a spin-charge separation [15].

²Notice that the SM results can be reproduced in some composite models, at least at the tree level, due to the complementarity between Higgs phase and confining phase [17, 18].

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However, new questions arise: Where does this peculiar haplon picture come from? Can it, in turn, result from a substructure of haplons? Consider the two scalar “prepreons”³ π_k and $\bar{\pi}_k$, which are SU(3) triplets and have the electric charges of $-1/6$ and $+1/6$, respectively. Then the set of haplons with their electric and C2 color charges can be reproduced by the triples of “prepreons” ($\bar{\pi}_r \pi_g \bar{\pi}_b \rightarrow \{\alpha, \ell\}$, $\pi_r \pi_g \pi_b \rightarrow \{\beta, \bar{\ell}\}$, $\bar{\pi}_r \pi_r \pi_l \rightarrow c_k$), while additional mechanism of spin generation is required. One can think of possible relation of spin to a circular color currents similarly to some discussions in the context of the condensed matter [20] and gravity [21] theories, taking into account that the distribution of matter in a composite state can be imagined (in particular, in the SU(3) Yang-Mills theory) in terms of the wave functions or probability distributions for the effective subcomponents of a finite size [22–25]. Then a “spinning” and “nonspinning” states of the same preon (e.g., α and ℓ) may form a supersymmetric multiplet.

Notice that the possibility of multihaplon states such as $(\beta \bar{c}_k \bar{\ell} c_k)$, $(\alpha \bar{\ell} \beta \bar{c}_k \beta c_k)$, etc., gets more points from recent discoveries of the multiquark hadrons [26].

3 Excited leptons

The excited lepton states defined in the introduction can be particularly important if their masses are smaller than the leptoquark and leptoquark masses, which can be natural due to the absence of the color charge. The contact interactions among the SM fermions f and the excited fermions f^* can be generically written as [9]

$$\begin{aligned} \mathcal{L}_{CI} = & \frac{g_*^2}{2\Lambda^2} \sum_{\alpha, \beta=L,R} \left[\eta_{\alpha\beta} (\bar{f}_\alpha \gamma^\mu f_\alpha) (\bar{f}_\beta \gamma_\mu f_\beta) \right. \\ & + \eta'_{\alpha\beta} (\bar{f}_\alpha \gamma^\mu f_\alpha) (\bar{f}_\beta^* \gamma_\mu f_\beta^*) + \tilde{\eta}'_{\alpha\beta} (\bar{f}_\alpha^* \gamma^\mu f_\alpha^*) (\bar{f}_\beta \gamma_\mu f_\beta) \\ & \left. + \eta''_{\alpha\beta} (\bar{f}_\alpha \gamma^\mu f_\alpha) (\bar{f}_\beta \gamma_\mu f_\beta) + \text{H.c.} + \dots \right], \quad (1) \end{aligned}$$

where Λ is the contact interaction scale, $g_*^2 = 4\pi$, and the new parameter values are usually taken of $|\eta^j| \leq 1$.

Assuming nearly maximal couplings of $|\eta^j| \simeq 1$ and the excited fermion masses of $M_{f^*} \simeq \Lambda$, the present lower bounds for $\Lambda/\sqrt{|\eta^j|}$ ratios are of the order of few TeV [9]. However, if Eq. (1) expresses a “residue” effective interactions between the composites (with respect to the fundamental interactions among their subcomponents) then $|\eta^j|$ couplings can be small, and even the case of $M_{f^*} \lesssim \Lambda < 1$ TeV is not excluded for $|\eta^j| \ll 1$ and $\Lambda/\sqrt{|\eta^j|} \gg M_{f^*}$.

A particular type of excited leptons that at low energies interact with the SM fermions dominantly through the contact terms we refer as *leptomesons* (LM).⁴ The relevant contact terms (with η'' couplings) can be realized, e.g., through the leptoquark exchange.

³For supersymmetric models with “prepreons” see, e.g., Ref. [19].

⁴Notice that the same term “leptomeson” was used in the literature for the bound states of colored excitations of e^+ and e^- [27].

4 Baryogenesis

Possible BG and the dark matter generation by a scalar 4-haplon state was considered in Ref. [12]. In this proceedings we discuss if fermionic LM states can provide a successful BG [28]. Similarly to the sterile neutrino ν_R case, depending on the LM properties, deviation from thermal equilibrium can occur at either production or freeze-out and decay (compare to the BG via ν_R oscillations [29] and the usual leptogenesis [30], respectively). In both scenarios one should replace the Yukawa interactions of ν_R by the contact interactions of LMs, which may result in promising effects.

4.1 BG from LM oscillations

Once created in the early universe neutral long-lived LMs oscillate and interact with ordinary matter. These processes do not violate the total lepton number L^{tot} (for Dirac LMs). However the oscillations violate CP and therefore do not conserve individual lepton numbers L_i for LMs. Hence the initial state with all zero lepton numbers evolves into a state with $L^{\text{tot}} = L_0 + \sum_i L_i = 0$ but $L_i \neq 0$, where L_0 is the lepton number of other particles.

At temperatures $T \ll \Lambda$ the LMs communicate their lepton asymmetry to neutrinos ν_ℓ and charged leptons ℓ through the effective interactions, e.g., B -conserving (and L -conserving for Dirac LMs) vector couplings

$$\begin{aligned} & \sum_{\psi_\ell, f, f'} \sum_{\alpha, \beta=L,R} \left[\frac{\tilde{\epsilon}_{ff'\psi_\ell}^{\alpha\beta}}{\Lambda^2} (\bar{f}_\alpha \gamma^\mu f'_\alpha) (\bar{\psi}_{\ell\beta} \gamma_\mu N_{\ell\beta}) \right. \\ & \left. + \frac{\tilde{\epsilon}_{ff'\psi_\ell}^{\alpha\beta}}{\Lambda^2} (\bar{\psi}_{\ell\alpha} \gamma^\mu f'_\alpha) (\bar{f}_\beta \gamma_\mu N_{\ell\alpha}) \right] + \text{H.c.}, \quad (2) \end{aligned}$$

where $\psi_\ell = \ell, \nu_\ell$ ($\ell = e, \mu, \tau$), the constants $\tilde{\epsilon} = 4\pi\eta''$ can be real, f and f' denote either quarks or leptons such that $Q_{f_\alpha} + Q_{f'_\alpha} + Q_{\psi_\ell} = 0$, and N_ℓ is the neutral LM flavor state related to the mass eigenstates N_i as $N_{\ell\alpha} = \sum_{i=1}^n U_{\ell i}^\alpha N_i$.

Suppose that LMs of at least one type N_i remain in thermal equilibrium till the electroweak symmetry breaking time t_{EW} at which sphalerons become ineffective, and those of at least one other type N_j come out-of-equilibrium by t_{EW} . Hence the lepton number of former (later) affects (has no effect on) the BG. In result, the final baryon asymmetry after t_{EW} is nonzero. At the time $t \gg t_{EW}$ all LMs decay into the leptons and the quarks (hadrons). For this reason they do not contribute to the dark matter in the universe, and do not destroy the Big Bang nucleosynthesis.

The system of n types of singlet LMs of a given momentum $k(t) \propto T(t)$ that interact with the primordial plasma can be described by the $n \times n$ density matrix $\rho(t)$. In a simplified picture it satisfies the kinetic equation [29]

$$i \frac{d\rho}{dt} = [\hat{H}, \rho] - \frac{i}{2} [\Gamma, \rho] + \frac{i}{2} [\Gamma^\rho, 1 - \rho], \quad (3)$$

where Γ (Γ^ρ) is the destruction (production) rate, and the effective Hamiltonian can be written as

$$\hat{H} = V(t) + U \frac{\hat{M}^2}{2k(t)} U^\dagger, \quad (4)$$

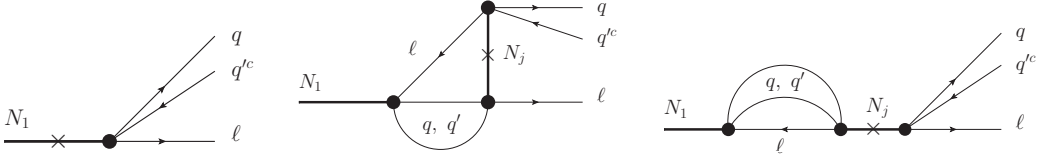


Figure 1. Feynman diagrams for the discussed contributions to the CP asymmetry, where the line direction shows either L or B flow, “ \times ” represents a Majorana mass insertion, and the black bulb represents a subprocess (e.g., a leptoquark exchange).

where $\hat{M}^2 = \text{diag}(M_1^2, \dots, M_n^2)$ with the masses M_i of N_i , and V is a real potential.

One of the main features of the discussed BG from LMs is that the 4-particle interaction cross section that contributes to the destruction rate is proportional to the total energy of the process s instead of the inverse proportionality that takes place in the BG from ν_R oscillations. Indeed, this cross section can be written as

$$\sigma(a + b \rightarrow c + d) = \frac{C}{4\pi} |\epsilon|^2 \frac{s}{\Lambda^4} \propto s, \quad (5)$$

where a, b, c and d denote the four interacting particles (f, f', ψ_ℓ and N_ℓ), and $C = O(1)$ is the constant that includes the color factor in the case of the interaction with quarks. In result, the interaction rate that equilibrates LMs,

$$\Gamma \propto |\epsilon|^2 \frac{T^5}{\Lambda^4}, \quad (6)$$

is suppressed by $(T/\Lambda)^4$ with respect to the Higgs mediated interaction rate in usual BG via ν_R oscillations.

The conditions that LMs of type N_i remain in thermal equilibrium till t_{EW} , while N_j do not, can be written as

$$\Gamma_i(T_{\text{EW}}) > H(T_{\text{EW}}), \quad \Gamma_j(T_{\text{EW}}) < H(T_{\text{EW}}), \quad (7)$$

where $H(T)$ is the Hubble expansion rate. Due to the suppression factor of $(T_{\text{EW}}/\Lambda)^4$ the successful BG can be realized with the relatively large couplings $|\epsilon|$ with respect to the sterile neutrino Yukawas of $Y \sim 10^{-7}$ in the BG via ν_R oscillations [29]. In particular, for $\Lambda \gtrsim 10$ and 30 TeV we have $|\epsilon| \gtrsim 10^{-4}$ and 10^{-3} , respectively. Hence the considered BG scenario can be relevant for the LHC and next colliders without unnatural hierarchy of the couplings.

4.2 BG from LM decays

Suppose that the neutral LMs are Majorana particles ($N = N^c$). Consider their out-of-equilibrium, CP - and L -violating decays in the early universe. The relevant interactions can be written as

$$\begin{aligned} & \frac{\epsilon_{ff'}^{\alpha R}}{\Lambda^2} (\bar{f}_\alpha \gamma^\mu f'_\alpha) (\bar{\psi}_{lR} \gamma_\mu N_{lR}) + \frac{\epsilon_{ff'}^S}{\Lambda^2} (\bar{f}_R f'_L) (\bar{\psi}_{lL} N_{lR}) \\ & + \frac{\epsilon_{ff'}^T}{\Lambda^2} (\bar{f} \sigma^{\mu\nu} f') (\bar{\psi}_{lL} \sigma_{\mu\nu} N_{lR}) + \text{H.c.} \end{aligned} \quad (8)$$

To be more specific in the following we consider the term

$$\frac{\lambda_{li}}{\Lambda^2} (\bar{q}_\alpha \gamma^\mu q'_\alpha) (\bar{\ell}_R \gamma_\mu N_{iR}), \quad (9)$$

where $\lambda_{li} = \epsilon_{qq'}^{\alpha R} U_{li}^R$ is a complex parameter. Consider the interference of tree and one-loop diagrams⁵ shown in Fig. 1. The final CP asymmetry that is produced in decays of the lightest LMs N_1

$$\varepsilon_1 = \frac{1}{\Gamma_1} \sum_\ell [\Gamma(N_1 \rightarrow \ell_R q_\alpha q'_\alpha) - \Gamma(N_1 \rightarrow \ell_R^c q_\alpha q'_\alpha)], \quad (10)$$

can be non-zero if $\text{Im}[(\lambda^\dagger \lambda)_{1j}^2] \neq 0$. Using the width [33],

$$\begin{aligned} \Gamma_1 &= \sum_\ell [\Gamma(N_1 \rightarrow \ell_R q_\alpha q'_\alpha) + \Gamma(N_1 \rightarrow \ell_R^c q_\alpha q'_\alpha)] \\ &\approx \frac{1}{128\pi^3} (\lambda^\dagger \lambda)_{11} \frac{M_1^5}{\Lambda^4}, \end{aligned} \quad (11)$$

the condition for the decay parameter $K \equiv \Gamma_1/H(M_1) > 3$ (strong washout regime) translates into the limit of

$$(\lambda^\dagger \lambda)_{11} \gtrsim 4 \times 10^{-7} \times \left(\frac{\Lambda}{10 \text{ TeV}} \right)^4 \times \left(\frac{1 \text{ TeV}}{M_1} \right)^3. \quad (12)$$

The final baryon asymmetry can be written as

$$\frac{n_B - n_{\bar{B}}}{s} = \left(-\frac{28}{79} \right) \times \frac{n_L - n_{\bar{L}}}{s} = \left(-\frac{28}{79} \right) \times \frac{\varepsilon_1 \kappa}{g_*}, \quad (13)$$

where n_B (n_L) is the baryon (lepton) number density, s is the entropy density, $-28/79$ is the sphaleron lepton-to-baryon factor, and $\kappa \leq 1$ is the washout coefficient that can be determined by solving the set of Boltzmann equations. The observed baryon asymmetry of the universe of [9]

$$\eta_B = \frac{n_B - n_{\bar{B}}}{n_\gamma} = 7.04 \times \frac{n_B - n_{\bar{B}}}{s} \simeq 6 \times 10^{-10}, \quad (14)$$

where n_γ is the photon number density, can be produced, e.g., for $K \sim 100$ with the degeneracy factor of

$$\mu \equiv \frac{M_2 - M_1}{M_1} \lesssim 10^{-6} \left(\frac{M_1}{1 \text{ TeV}} \right), \quad (15)$$

which enters the resonant CP asymmetry of

$$\varepsilon_i \sim \frac{\text{Im}\{[(\lambda^\dagger \lambda)_{ij}]^2\}}{(\lambda^\dagger \lambda)_{ii} (\lambda^\dagger \lambda)_{jj}} \frac{\Gamma_j}{M_j} \frac{M_i M_j}{M_i^2 - M_j^2} \sim \mu^{-1} \frac{\Gamma_1}{M_1}. \quad (16)$$

Notice that the discussed effective LM-quark-quark-lepton vertex can be realized, e.g., through the exchange

⁵Same two-loop self-energy graph as in Fig. 1 (right) was discussed in the resonant BG scenarios of Refs. [31, 32], where the baryon asymmetry is directly produced in the three-body decays of the sterile neutrinos. Although these mechanisms involve B -violating interactions of $qqq\nu_R$ type, they do not lead to fast proton decay due to the large values of ν_R mass and the B -violating interaction scale of $O(1)$ TeV.

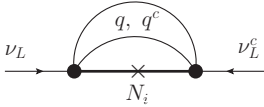


Figure 2. The discussed contribution to the neutrino masses (line direction shows L flow).

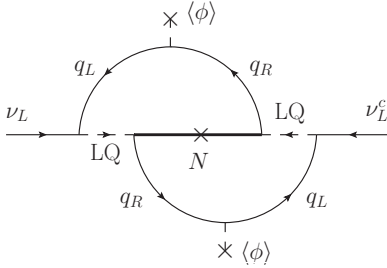


Figure 3. Leptoquark (LQ) contribution to the neutrino masses.

of the $SU(2)_L$ singlet scalar leptoquark S_{0R} with $Y = 1/3$. Relevant interaction terms can be written as

$$-\mathcal{L}_{\text{int}} = (g_{ij} \bar{d}_R^c N_{iR} + f_j \bar{u}_R^c \ell_R) S_{0R}^j + \text{H.c.} \quad (17)$$

Then the above expressions are valid with the replacements of $\lambda \rightarrow g f^*$ and $\Lambda \rightarrow M_{S_{0R}}$. The relevant to the successful BG values of the new couplings of $|g| \sim |f| \sim 0.01 - 0.1$, can be interesting for the collider searches.

4.3 Neutrino masses

For Majorana LMs the effective terms of

$$\frac{\epsilon_{ff\nu\ell}^S}{\Lambda^2} \bar{f}_R f_L \bar{\nu}_{iL} N_{iR} + \frac{\epsilon_{ff\nu\ell}^T}{\Lambda^2} \bar{f} \sigma^{\mu\nu} f \bar{\nu}_{iL} \sigma_{\mu\nu} N_{iR} + \text{H.c.} \quad (18)$$

can generate the two-loop contributions to the observable light neutrino masses m_{ν_ℓ} that is illustrated for $f = q$ in Fig. 2, and for a particular model with the leptoquarks in Fig. 3. A simple estimate of this contribution is

$$m_{\nu_\ell} \sim \sum_i \frac{(\epsilon U_{\ell i}^R)^2}{(16\pi^2)^2} \frac{M_i^3 m_q^2}{\Lambda^4}, \quad (19)$$

where m_q is the quark mass. Hence the present bound of $m_\nu \lesssim 2 \text{ eV}$ [34] can be easily satisfied.

5 Conclusion

The two new testable baryogenesis scenarios in the models with excited leptons are introduced, which do not contradict to the observable neutrino masses. First, *the BG from LM oscillations* may take place for relatively light and long-lived LMs, which do not all decay before t_{EW} . Second, *the BG from LM decays* can be realized if all LMs decay before t_{EW} . A particular models based on the former (later) BG proposal require detailed study of the neutrino

potential (of the Boltzmann equations) to be verified in the future experiments. In both scenarios the baryon number is violated only by the sphaleron processes that does not affect the proton stability. Due to the relatively low temperatures of the discussed BG mechanisms, an analog of the gravitino problem [35, 36] does not exist there.

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