

# Validation of the MCNP6 electron-photon transport algorithm: multiple-scattering of 13- and 20-MeV electrons in thin foils

David A. Dixon<sup>1</sup> and H. Grady Hughes<sup>1</sup>

<sup>1</sup>*Los Alamos National Laboratory, Los Alamos, NM 87545*

This paper presents a validation test comparing angular distributions from an electron multiple-scattering experiment with those generated using the MCNP6 Monte Carlo code system. In this experiment, a 13- and 20-MeV electron pencil beam is deflected by thin foils with atomic numbers from 4 to 79. To determine the angular distribution, the fluence is measured down range of the scattering foil at various radii orthogonal to the beam line. The characteristic angle (the angle for which the max of the distribution is reduced by  $1/e$ ) is then determined from the angular distribution and compared with experiment. Multiple scattering foils tested herein include beryllium, carbon, aluminum, copper, and gold. For the default electron-photon transport settings, the calculated characteristic angle was statistically distinguishable from measurement and generally broader than the measured distributions. The average relative difference ranged from 5.8% to 12.2% over all of the foils, source energies, and physics settings tested. This validation illuminated a deficiency in the computation of the underlying angular distributions that is well understood. As a result, code enhancements were made to stabilize the angular distributions in the presence of very small substeps. However, the enhancement only marginally improved results indicating that additional algorithmic details should be studied.

**KEYWORDS:** MCNP6, Monte Carlo, electron-photon transport, multiple-scattering, validation

## I. Introduction

It is of interest to establish a comprehensive electron-photon validation suite for the MCNP<sup>®</sup> version 6.2 code system.<sup>(1)</sup> Though previous efforts to validate the electron-photon transport in earlier versions of MCNP6 were made,<sup>(2-4)</sup> these validations correspond to unsupported versions of the code and are somewhat limited in scope. New validation efforts are underway<sup>(5)</sup> that test the current version of the code containing improvements to the electron-photon transport capabilities.<sup>(6,7)</sup>

The purpose of this paper is to benchmark MCNP6 against the measurements of Ross *et al.*<sup>(8)</sup> Ross *et al* measured angular distributions corresponding to electron pencil beams of 13- and 20-MeV normally incident on single-element (with exception of a titanium alloy), thin foils with atomic numbers from 4 to 79. Electron fluence was measured roughly 120 cm from the beam exit window at various radii orthogonal to the beam line. While a number of similar benchmarks exist, uncertainties in the accelerator descriptions associated with these similar experiments somewhat limit their utility. Whereas, Ross *et al* determined angular distributions where the beam characteristics are well known, with no free parameters, making this benchmark ideal for validation of electron transport codes. Though this benchmark is valuable to codes intended for medical physics applications, it is also valuable for general validation of multiple-scattering models. As such, this experimental benchmark will be added to the MCNP6 electron-photon validation suite.

As the emphasis of this paper is the validation test rather than MCNP6 electron-photon transport algorithm, we briefly describe the relevant physics and identify additional references that provide greater detail. The MCNP6 electron-photon trans-

port algorithm is considered a class I condensed history method. In MCNP6, electrons traverse a substep distance where both angular deflection and energy-loss are sampled (i.e. the primary processes impacting the angular distributions studied herein). In MCNP6, angular deflection resulting from multiple elastic scattering of electrons is given by the Goudsmit-Saunderson distribution,<sup>(9,10)</sup> and energy-loss resulting from multiple inelastic scattering is given by the Landau distribution.<sup>(11)</sup> Additionally, secondary electrons and bremsstrahlung photons are sampled along the substep, and the necessary atomic relaxation processes are considered. A substep is determined by dividing an energy step, or the range corresponding to some  $\Delta E_k$ , into  $M$  smaller steps, where  $M$  is an increasing function of the target atomic number. Calculation of the energy steps, substeps, and the necessary electron transport data are completed during a preprocessing stage for a fixed grid of  $k$  energies. While electron and photon transport is coupled, the impact of photons on the angular distribution is nominal, and therefore, not discussed herein. More detail on the electron-photon transport algorithm is given by Hughes.<sup>(12)</sup>

The accuracy of the aforementioned processes depends on both algorithmic limitations and the underlying differential cross sections. In essence, this validation tests both potential sources of error. However, assume for a moment that the underlying differential cross sections are sufficiently accurate. Under these conditions, algorithmic error would account for any disagreement between simulation and experiment. Specific examples of algorithmic error relevant to this study include the stability of the underlying angular distributions, the boundary crossing approximation, the solution dependence on the

step-size, and the treatment of angular deflection attributed to impact-ionization collisions for both the primary and the knock-on electron (again, details on each of these can be found in the literature<sup>(12,13)</sup>).

## II. Method

### 1. Experiment

The characteristics of the experiment are detailed in Ross *et al.*<sup>(8)</sup> Below we provide sufficient detail to understand the simulation setup and the validation results. For the purposes of this paper, we consider a beam of electrons emerging from an accelerator exit window (see Fig. 1). The beam is effectively a pencil beam with some spread in space, angle, and energy. However, only the spread in space is considered. The spatial distribution of the experimental beam in transverse directions did not have cylindrical symmetry, but could be approximated by a gaussian with a full-width at half-max of approximately 1-mm.<sup>(8)</sup> Immediately upstream of the exit window is the scattering foil that is followed by the remainder of the apparatus including a transmission monitor chamber, a helium bag, and then the detector.

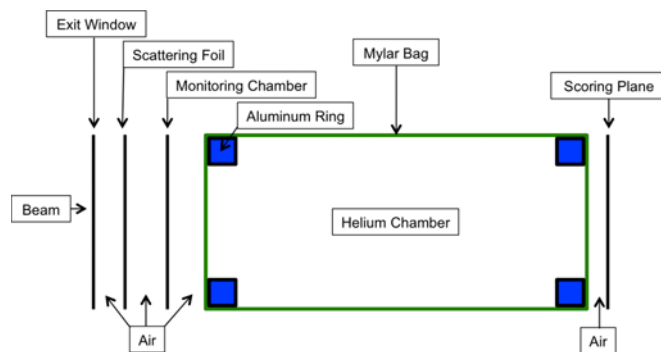


Figure 1: Schematic of experimental apparatus (not to scale).

To sufficiently resolve the scattered electrons a significant distance between the foil and the detector was required. What would be an air column was replaced by a helium column to reduce the scattering power of the region between the foil and the detector. For this reason, a Mylar bag was constructed to contain the helium.

The scattered electron distributions were measured using both an ionization chamber and diode detectors. The point of measurement was assumed to be the center of the detector. The detectors were mounted on a linear translator positioned orthogonal to the beam line. The detectors were then moved along the translator to determine the fluence at a given radial distance from the beam line that is then converted to angle.

### 2. Simulation

In this work, the experimental apparatus was approximated by several cylinders. The position and thickness of the cylinders were taken from the description of the key experimental components in Ross *et al.*<sup>(8)</sup> The electron source was placed at the origin and directed along the positive x-axis. The beam was assumed to be Gaussian in the transverse directions with a

FWHM=1-mm. One billion source particles were followed in each simulation. The scoring plane was taken to be the plane orthogonal to the beam located at a point furthest from the source. The F2 surface flux tally was applied to this plane, which was subdivided into rings with increasing radii at 0.1 cm spacings from 0.1 cm out to 21 cm (this is accomplished using the *FS* or tally segment card).

The default photon physics was used to transport photons; however, it is of interest to note that the new EPRDATA library<sup>(7)</sup> (i.e. the ".12p" extension) was used for photon data.

For electrons, three physics settings were studied including the default physics, and two other settings where the impact of modifying *ESTEP* and *EFAC* are studied in isolation. *ESTEP* controls the number of substeps per energy step and presumably improves accuracy in problems with optically thin regions where electrons stream through these regions without colliding. Under such conditions, truncation of the electron substep is likely and results in an approximation of the angular deflection distribution. Therefore, one can increase *ESTEP* to ensure at least a few of collisions occur before reaching a boundary. *ESTEP* is applied to specific materials, so one can modify the substep size only for the material of interest that is contained in the thin cells. One should always keep in mind that reducing the substep size increases runtime. The following two tables present the *ESTEP* settings necessary to ensure that roughly 10 condensed collisions occur in each foil. These settings correspond to results labeled as *ESTEP*. Default values of *ESTEP* are for the materials used herein are: Be = 2, C = 3, Al = 5, Cu = 7, Ta = 12, and Au = 13.

Table 1: *ESTEP* parameter analysis for 13-MeV simulations. *ESTEP* was determined such that there were approximately 10 substeps per foil thickness (that is, electrons will suffer roughly 10 condensed collisions when traversing the foil).

Material	Thickness (mg/cm <sup>2</sup> )	Step size (g/cm <sup>2</sup> )	<i>ESTEP</i>
Be	926	$6.29 \times 10^{-1}$	8
Al	70.1	$5.33 \times 10^{-1}$	80
	140		40
Cu	43.0	$4.95 \times 10^{-1}$	119
	86.4		63
	129.6		42
	174.5		35
Ta	44.3	$4.13 \times 10^{-1}$	24
Au	31.2	$3.96 \times 10^{-1}$	130
	54.8		78
	93.7		52

In the context of this work, *EFAC* has a similar effect, but *EFAC* is a global parameter in the sense that it is used to adjust the coarseness, or fineness, of the energy grid that is material independent. So, rather than selectively reducing the substep size in thin cells, the substep size is reduced in all cells. By default *EFAC* = 0.917. Increasing it to 0.99 (maximum permissible value) the substep size decreases, and runtime increases.

**Table 2: ESTEP parameter analysis for 20-MeV simulations. ESTEP was determined such that there were approximately 10 substeps per foil thickness (that is, electrons will suffer roughly 10 condensed collisions when traversing the foil).**

Material	Thickness (mg/cm <sup>2</sup> )	Step size (g/cm <sup>2</sup> )	ESTEP
Be	926	9.078×10 <sup>-1</sup>	10
C	546	7.83×10 <sup>-1</sup>	15
Al	70.1	7.22×10 <sup>-1</sup>	105
	140		55
	274		30
Cu	43.0	6.20×10 <sup>-1</sup>	147
	86.4		77
	129.6		49
	174.5		42
Ta	44.3	4.76×10 <sup>-1</sup>	24
Au	54.8	4.53×10 <sup>-1</sup>	91
	164.2		39

### 3. Results Processing

Characteristic angles are determined by fitting a Gaussian to the measured and calculated angular distributions,  $R(\theta)$ , over a partial range. The partial range corresponds to all values where the distributions are greater than approximately  $R_{max}/e$  or  $1/e$  if the distribution is normalized such that the amplitude is unity. The fit is then used to obtain the characteristic angle by setting the Gaussian equal to  $1/e$  and solving for  $\theta = \theta_{1/e}$ . That is,

$$R(\theta_{1/e}) = \frac{R_{max}}{e} = a \exp(-\theta_{1/e}^2/b), \quad (1)$$

where  $a$  and  $b$  are the Gaussian fit parameters obtained in the first step. Therefore,

$$\theta_{1/e} = \sqrt{-b \ln\left(\frac{R_{max}}{ea}\right)}, \quad (2)$$

or if  $R_{max} = a = 1$ ,  $\theta_{1/e} = \sqrt{b}$ .

Ross *et al* provide the experimental results where the independent variable is given in either the raw form, or shifted. The shifted results are used herein, as they correspond to a symmetric distribution necessary for a consistent comparison.

### III. Results

The characteristic angles of the measured and calculated angular distributions are presented in Tables 4 and 5. In these tables there are several characteristic angles reported as described below:

- Ross *et al.*: Characteristic angles reported by Ross *et al.*<sup>(8)</sup> using the fitting technique described therein.
- Dixon and Hughes: Characteristic angles determined by applying the fitting technique described herein to the data from Ross.<sup>(8)</sup>
- Calc.: Characteristic angles determined by applying the fitting technique described herein to the angular distributions calculated using MCNP6.

Furthermore, there are three different calculated results corresponding to different physics settings: the default physics; the default physics with modified ESTEP parameters, and the default physics modified EFAC parameter of 0.99.

The average difference (i.e. an arithmetic average of the rel. diff. over each of the foils) in the measured and calculated characteristic angle for 13-MeV and 20-MeV electrons for the physics settings tested are presented in Table 3. According to Faddegon<sup>(14)</sup> *et al.*, for 13-MeV and 20-MeV electrons the average difference in the measured and calculated characteristic angle for EGSnrc was 1.3% and 1.1% respectively; for PENELOPE the average difference was 1.1% and 1.1% respectively; and for Geant4 the average difference was 0.7% and 0.9% respectively.

**Table 3: Average of the relative differences between measured and calculated characteristic angles for 13-MeV and 20-MeV electrons for three different physics configurations including: default, ESTEP, and EFAC.**

Physics Setting	Average Relative Difference	
	13-MeV	20-MeV
Default	6.4%	5.8%
ESTEP	10.1%	12.2%
EFAC	7.3%	9.5%

The average values presented above were determined from the relative difference results in Tables 4 and 5. The values in Tables 4 and 5 compare characteristic angles corresponding to MCNP6 computed angular distributions (i.e. Calc.) with the Dixon and Hughes column such that a consistent fitting technique is used between the compared results. The characteristic angle reported by Ross<sup>(8)</sup> was included to demonstrate that the fitting technique used by Dixon and Hughes gives characteristic angles that are in agreement with those reported by Ross.<sup>(8)</sup>

In general, the characteristic angles are overestimated as seen in Tables 4 and 5 and in Figs. 2 and 3. In Figs. 2 and 3, the angular distribution for 13- and 20-MeV electrons on several foils with the default physics are presented.

### IV. Conclusions

The multiple-scattering validation test of the MCNP6 Monte Carlo code system reveals deficiencies in the treatment of angular deflection in MCNP6 for the default electron-photon transport settings. Moreover, refinement of the electron transport physics via ESTEP and EFAC only exacerbate the disagreement between the measured and calculated results. The cause of the disagreement remains an open question, but it is widely known that the class I condensed algorithm for electron transport contained in MCNP6 has inherent limitations. Of relevance to this work, these limitations include both the boundary crossing approximation and stepsize artifacts. Boundary crossing approximations are particularly sensitive when transporting electrons between high and low density regions. The Fano cavity test should be considered to quantify the impact of this effect. Stepwise artifacts result from the application of angular

**Table 4: Comparison of measured and calculated characteristic angles for 13-MeV electrons on Be, Al, Cu, Ta, and Au. Three different physics configurations are studied: default, ESTEP, and EFAC=0.99.**

Material	Thickness (mg/cm <sup>2</sup> )	Ross <i>et al.</i> (deg.)	Dixon & Hughes (deg.)	Default		ESTEP		EFAC	
				Calc. (deg.)	Rel. Diff. (%)	Calc. (deg.)	Rel. Diff. (%)	Calc. (deg.)	Rel. Diff. (%)
Be	926	8.143	8.089	8.363	3.382	8.846	9.353	9.081	12.26
Al	70.1	4.003	3.981	4.308	8.223	4.558	14.48	4.289	7.745
	140	5.268	5.226	5.602	7.196	5.798	10.93	5.652	8.150
Cu	43.0	4.219	4.167	4.541	8.968	4.745	13.86	4.450	6.786
	86.4	5.630	5.562	5.970	7.310	6.161	10.74	5.956	7.056
	129.6	6.861	6.803	7.115	4.593	7.323	7.645	7.210	5.989
	174.5	7.956	7.911	8.521	7.704	8.681	9.722	8.619	8.938
Ta	44.3	5.558	5.503	5.812	5.608	5.787	5.158	5.799	5.385
Au	31.2	4.878	4.798	5.061	5.473	5.346	11.43	5.164	7.626
	54.8	6.329	6.260	6.568	4.911	6.769	8.124	6.650	6.216
	93.7	8.243	8.231	8.812	7.052	9.013	9.495	8.986	9.163

**Table 5: Comparison of measured and calculated characteristic angles for 20-MeV electrons on Be, C, Al, Cu, Ta, and Au. Three different electron physics configurations are studied: default, ESTEP, and EFAC=0.99.**

Material	Thickness (mg/cm <sup>2</sup> )	Ross <i>et al.</i> (deg.)	Dixon & Hughes (deg.)	Default		ESTEP		EFAC	
				Calc. (deg.)	Rel. Diff. (%)	Calc. (deg.)	Rel. Diff. (%)	Calc. (deg.)	Rel. Diff. (%)
Be	926	5.238	5.214	5.434	4.234	5.694	9.208	5.839	11.99
C	546	5.132	5.108	5.198	1.764	5.456	6.825	5.561	8.873
Al	70.1	2.653	2.634	2.873	9.094	3.148	19.51	2.877	9.221
	140	3.484	3.463	3.657	5.611	4.018	16.02	3.823	10.39
	274	4.777	4.750	4.978	4.798	5.269	10.94	5.268	10.91
Cu	43.0	2.790	2.768	3.025	9.284	3.279	18.49	2.999	8.375
	86.4	3.714	3.685	3.891	5.605	4.245	15.21	4.041	9.666
	129.6	4.493	4.454	4.758	6.843	4.979	11.81	4.900	10.02
	174.5	5.198	5.147	5.429	5.496	5.724	11.22	5.688	10.51
Ta	206.3	7.913	7.809	8.207	5.099	8.306	6.364	8.793	12.60
Au	54.8	4.127	4.111	4.382	6.590	4.652	13.17	4.563	10.99
	164.2	7.278	7.258	7.593	4.605	7.775	7.113	8.025	10.57

deflection data that is not fully converged. That is, MCNP6 currently allows for no more than 240 series solution terms when computing the angular deflection distributions. In many cases, additional terms are required to ensure that the solution has converged. The code was modified to determine if allowing more terms would improve the angular distributions. We found that allowing for several thousand terms in the computation of the underlying angular deflection distributions, the agreement in the characteristic angle improved by roughly 1% for each of the different physics settings.

## Acknowledgments

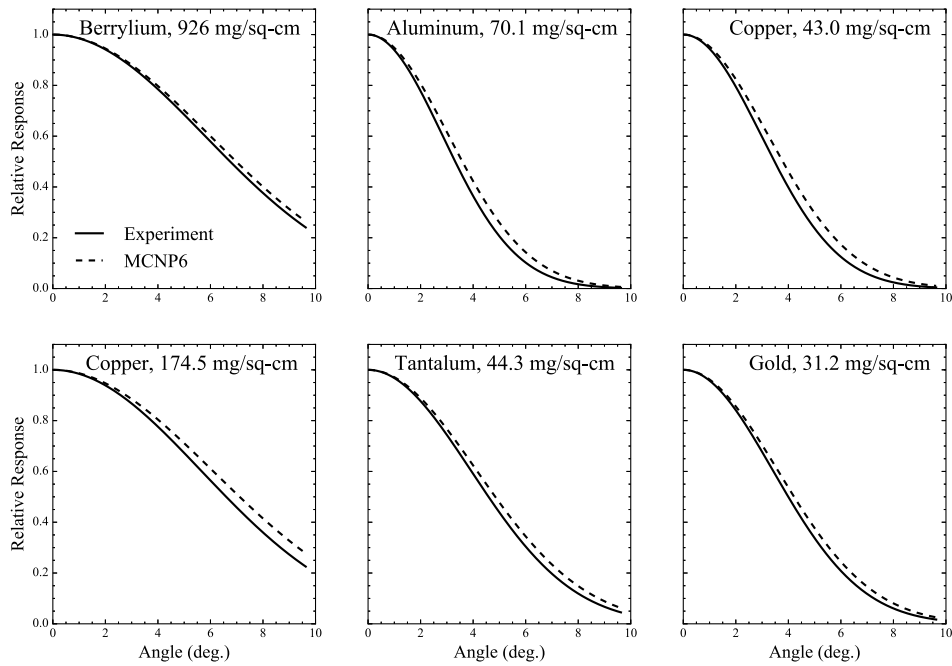
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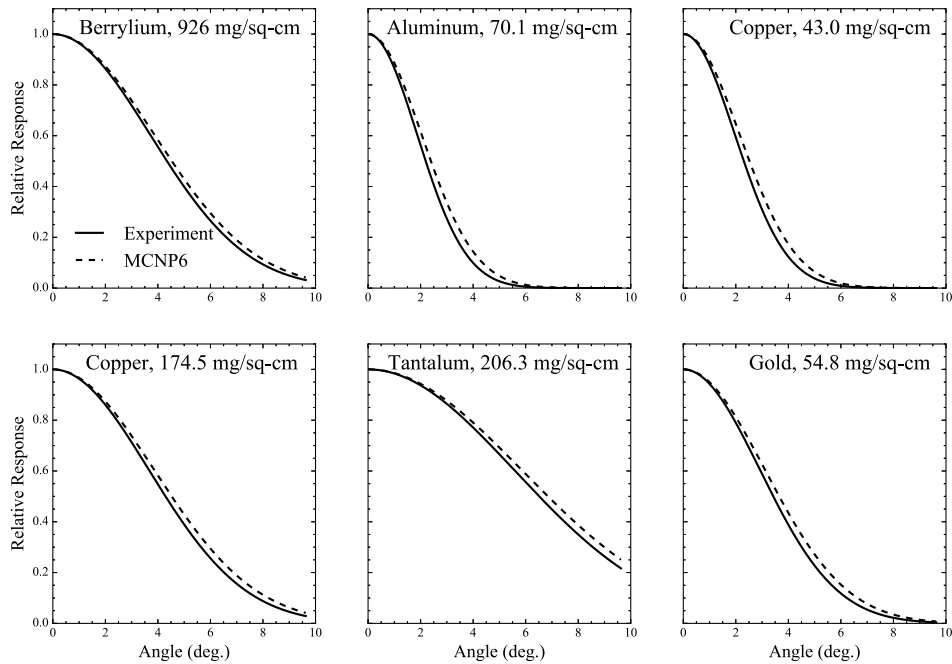
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**Figure 2: Comparison of Gaussian fits to experiment (solid curve) and calculated (dashed curve) angular distributions for 13-MeV electrons on various foils.**



**Figure 3: Comparison of Gaussian fits to experiment (solid curve) and calculated (dashed curve) angular distributions for 20-MeV electrons on various foils.**