Multichannel all-optical switch based on a thin slab of resonant two-level emitters

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Abstract. We discuss the possibility of using a thin layer of inhomogeneously broadened resonant emitters as a multichannel all-optical switch. Switching time from the lower stable branch of the system’s bistable characteristics to the upper one and vice versa, which determines the speed of operation of a bistable device, is studied.

Bistability of an optical system is an extraordinary nonlinear effect, the essence of which is that the system has two stable states of transmission at the same value of the external field [1]. This phenomenon is of ongoing interest for optical technologies, in particular, for creation of fully optical switches, optical transistors, and optical memory.

In this work, we model the optical response of a thin resonant slab (TRS) of thickness L with inhomogeneously broadened two-level emitters of density N₀. As is well known, for such systems, the local field correction (LFC) to the field acting on an emitter is of importance [2,3]. The basis of our model consists of equations for the density matrix elements ρ_{ab} (a,b = 1,2) of a two-level emitter and the field E inside the slab, which in the rotating wave approximation reads [2,3]

\[ \dot{R} = (i\Delta - \Gamma_2)R - \Omega Z, \]
\[ \dot{Z} = \frac{1}{2}(\Omega^* R + \Omega R^*) - \Gamma_1 (1 + Z), \]
\[ \Omega = \Omega_0 + (i\Delta_L - \Gamma_\phi)\langle R \rangle, \]
\[ \langle R \rangle = \int R(x,t,\Delta)G(\Delta)d\Delta. \]

Here, R is the amplitude of the off-diagonal density matrix element ρ_{01} = -(i/2)Re^{i\omega t}; Z = ρ_{22} - ρ_{11}; Γ₁ and Γ₂ are the relaxation constants of the population and coherence, respectively; Ω₀ = dΕ₀/ℏ and Ω = dΕ/ℏ are the external field Ε₀ and the field E inside the slab in frequency units; d is the transition dipole moment of the emitter; ℏ is the reduced Plank constant; Δ = ω - ω₂₁ is the detuning away from resonance of an isolated emitter, ω₂₁; Γ₂ = 2πd²N₀kL/ℏ is the superradiant constant [2,3] and Δ_L = 4πd²N₀/(3ℏ) is the LFC [2,3]; G(Δ) is the inhomogeneous contour, in what follows, a Lorentzian of width Γ₂*.

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Fig. 1. (a) – Optical hysteresis of the field intensity \( y = |\Omega|^2/(\Gamma_1 \Gamma_2) \) inside the slab, obtained under adiabatically sweeping the external field intensity \( x = \Omega_0^2/(\Gamma_1 \Gamma_2) \) up and down (shown by arrows). The parameters of calculations are: \( \Delta L = 20 \Gamma_2 \), \( kL = 0.1 \), \( \Delta_0 = -5 \Gamma_2 \), \( \Gamma_1 = 2 \Gamma_2 \), \( \Gamma_2^* = 0.5 \Gamma_2 \). (b) and (c) – Time evolution of the integrated population difference \( \langle Z \rangle \) after an abrupt switching of the external field intensity \( x \) from zero to a value that exceeds the critical intensity \( x_{cr} = (\Omega_0^{cr})^2/(\Gamma_1 \Gamma_2) = 56.096 \). The parameters of calculations are the same as in (a), except \( \Gamma_2^* = 2 \Gamma_2 \). (d) – Intensity dependence of the switching time \( \tau \).

Examples of numerical calculations of the optical response of the TRS are presented in Fig. 1. Plot (a) shows the optical hysteresis of the system: a sudden switching from one stable state to the other one and back under the adiabatically sweeping the driving field \( \Omega_0 \) up and down (shown by arrows), occurring at different magnitudes of \( \Omega_0 \). So the system is bistable. Plots (b) and (c) display the dynamics of switching at the high-intensity switching point \( x_{cr} = (\Omega_0^{cr})^2/(\Gamma_1 \Gamma_2) = 56.096 \) (for the parameters specified in the figure caption). Plot (d) shows the intensity dependence of the switching time \( \tau \). Conspicuous is the fact that the dynamics of switching reveals a critical slowing down around the high-intensity switching point \( x_{cr} \). Switching down always occurs during a time interval on the order of \( \Gamma_1^{-1} \).

Fig. 2. Schematics of a multichannel all-optical switch (see details in the text).

Bistable properties of a thin resonant slab allow one to view this system as an all-optical switch. Apparently, different regions of the TRS are switched independently of each other. Then, this system may operate as a multichannel all-optical switch that is of importance from the viewpoint of optical information processing and optical computing. As an implementation of a device of such a type, an array of dielectric optical waveguides coupled with a TRS could be proposed, as illustrated in Fig. 2.

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References