Production spectra of zero-degree neutral particles measured by the LHCf experiment

DOI: 10.1051/epjconf/201716407028

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Abstract. The LHC-forward (LHCf) experiment, situated at the LHC accelerator, has measured neutral particles production in a very forward region (pseudo-rapidity $\eta > 8.4$) in proton-proton and proton-lead collisions. The main purpose of the LHCf experiment is to test hadronic interaction models used in ground based cosmic rays experiments to simulate cosmic rays induced air-showers in the Earth's atmosphere.

The experiment is composed of two independent detectors located at 140m from the AT-LAS interaction point (IP1) on opposite sides; each detector is composed of two sampling and position sensitive calorimeters.

In this paper, latest published physics results from p-p and p-Pb collisions (at $\sqrt{s} = 7$, 2.76 TeV and $\sqrt{s_{NN}} = 5.02$ TeV, respectively) compared with Monte Carlo predictions of DPMJET, EPOS, PYTHIA, QGSJET and SIBYLL event generators will be presented. In particular, the inclusive energy spectra of neutrons in p-p collisions and the transverse and longitudinal momentum spectra of neutral pions for different pseudo-rapidity ranges in p-p and p-Pb collisions will be shown; then, test of Feynman scaling hypothesis using neutral pion spectra will be discussed. Preliminary results of photon inclusive energy spectra in p-p collisions at $\sqrt{s} = 13$ TeV will be also presented.

1 Introduction

The LHC-forward experiment (LHCf) has measured neutral particles production in a very forward region in proton-proton and proton-lead collisions at the Large Hadron Collider. The main purpose of LHCf is to improve hadronic interaction models of Monte Carlo (MC) simulations used in cosmic rays indirect measurements. Highest energy cosmic rays can only be detected from secondary particles which are produced in the interaction of the primary particle with nuclei of the atmosphere, the so-called air showers. Studying the development of air showers, it is possible to reconstruct the type and kinematic parameters of primary particles. To reproduce the development of air showers, MC simulations with accurate hadronic interaction models are needed. Since the energy flow of secondary particles is concentrated in the forward direction, measurements of particle production in the high pseudo-rapidity region (i.e. small angles) are very important. Furthermore, in the forward region soft QCD interactions (non perturbative regime) dominates and MC simulations of air showers are based on phenomenological models, so inputs from experimental data are crucial for the tuning of that

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models. LHC accelerator gives the possibility to study a wide range of collision energies, from 0.9 TeV to 13 TeV in the center of mass frame, which corresponds to an energy range in the laboratory frame from 10^{14} eV to 0.9×10^{17} eV. This energy range covers the "knee" region of cosmic rays spectrum, which occurs around 10^{15} eV.

2 The detector

LHCf is composed of two independent detectors, called *Arm1* and *Arm2*. Arm1 is located 140 meters away from ATLAS interaction point (IP1) in the IP8 direction, while Arm2 is placed 140 meters away from IP1 in the opposite direction (toward IP2). Detectors are placed inside Target Neutral Absorber (TAN), where the beam pipe turns into two separates tubes. Since charged particles are deviated by the D1 dipole magnet (which bends colliding beams into the two separate beam pipes), only neutral particles, mainly photons and neutrons, reach the detector.

Each detector was made of two sampling and imaging calorimeters (called *towers* hereafter). Each tower is composed of 16 tungsten layers and 16 scintillator layers to measure the energy deposit and it also contains 4 position sensitive layers. During 0.9 TeV, 2.76 TeV and 7 TeV operations at the LHC, plastic scintillators (EJ-260) were used. Arm1 detector used scintillating fiber (SciFi) to measure position, while Arm2 used silicon microstrip detectors. For 13 TeV operation both detectors were upgraded: all the plastic scintillators were replaced by Gd_2SiO_5 (GSO) scintillators because of their radiation hardness; also the Arm1 SciFi were replaced by GSO bars. In Arm2 the signal of silicon detectors was reduced using a new bonding scheme of the microstrips to avoid saturation of readout electronics due to the higher energy deposit expected at $\sqrt{s} = 13$ TeV.

Transverse cross sections of towers are $20 \times 20 \text{ mm}^2$ and $40 \times 40 \text{ mm}^2$ for Arm1 and $25 \times 25 \text{ mm}^2$ and $32 \times 32 \text{ mm}^2$ for Arm2. Longitudinal dimension of towers is of 44 radiation lengths, which correspond to 1.6 nuclear interaction lengths. Energy resolution is better than 5 % for photons and of about 40 % for neutrons. Position resolution for photons is $200 \, \mu \text{m}$ and $40 \, \mu \text{m}$ for Arm1 and Arm2 respectively, while position resolution for neutrons is of about 1 mm. Smaller tower of each detector is placed on the beam center and covers the pseudo-rapidity range $\eta > 9.6$, while larger tower covers the pseudo-rapidity range $\theta = 8.4 < \eta < 9.4$. More detailed descriptions of detector performance are reported elsewhere [1–3].

3 Experimental results

LHCf collected data in three different low-luminosity dedicated runs in 2010, 2013 and 2015. Protonproton collisions at $\sqrt{s} = 900$ GeV (2010), $\sqrt{s} = 2.76$ TeV (2013), $\sqrt{s} = 7$ TeV (2010) and $\sqrt{s} = 13$ TeV (2015) were produced; proton-lead collisions at $\sqrt{s_{NN}} = 5.02$ TeV were produced in 2013.

Results for single photon energy spectra in p-p collisions at $\sqrt{s} = 900$ GeV, 7 TeV and for π^0 transverse momentum spectra in p-p and p-Pb collisions at $\sqrt{s} = 7$ TeV and $\sqrt{s_{NN}} = 5.02$ TeV, respectively, have already been published [4–7]. Latest published results are presented in this paper: neutron energy spectra in p-p collisions at $\sqrt{s} = 7$ TeV [8] are shown in Sect. 3.1 while results from π^0 transverse (P_T) and longitudinal (P_Z) momentum analysis in p-p at $\sqrt{s} = 2.76$, 7 TeV and p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV [9] are shown in Sect. 3.2. Finally, also preliminary result from photon analysis at $\sqrt{s} = 13$ TeV will be presented in Sect. 3.3.

3.1 Neutron energy spectra in p-p collisions

In this section results from neutron analysis in p-p collisions at $\sqrt{s} = 7$ TeV are presented [8]. Data were taken during LHC Fill #1104 on May 15, 2010. Integrated luminosities were 0.68 nb⁻¹ and 0.53

nb⁻¹ for Arm1 and Arm2 respectively, with a typical luminosity of 6.4×10^{28} cm⁻²s⁻¹, derived from LHCf front counters [10].

Due to the limited dimensions of calorimeter, hadronic showers are not fully contained in the tower and the study of hadronic neutral particles is less precise and more complicated than photons. Hadronic particles (mostly neutrons) are selected using information of longitudinal distribution of energy deposit in the calorimeter. Since the energy resolution for neutrons is $\sim 40\%$, the true energy spectra are smeared by the detector response. Smeared energy spectra of neutrons measured by Arm1 and compared with MC simulation predictions of several models are shown in figure 1 in three pseudorapidity ranges.

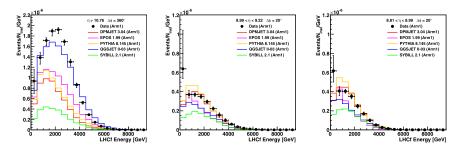


Figure 1. Energy distribution of neutrons in p-p collisions at 7 TeV before unfolding in three pseudo-rapidity regions: coloured histograms show results from Monte Carlo simulations with different hadronic interaction models, black points show experimental data of Arm1. Particles above maximum available energy (3500 GeV) are clearly due to the smearing effect of the detector.

Unfolding technique is therefore needed to estimate the true energy spectra: a multidimensional-spectra unfolding Bayesian method [11] relying on energy and P_T measured variables was used. Results after unfolding are shown in figure 2. It is evident from the results that in the very forward pseudo-rapidity region ($\eta > 10.76$) there is a large amount of high energy neutrons, only predicted by the QGSJETII-03 [12] model. In other regions data lie between models predictions, but there is not an unique model that represents data.

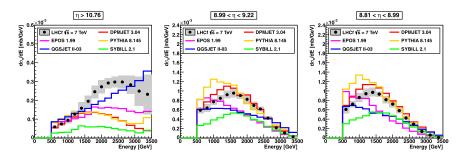


Figure 2. Energy distribution of neutrons in p-p collisions at 7 TeV after unfolding in three pseudo-rapidity regions: coloured histograms show results from Monte Carlo simulations with different hadronic interaction models, black points show combined experimental data of Arm1 and Arm2. A large amount of very-forward ($\eta > 10.76$) high energy neutrons is predicted by QGSJETII-03 and confirmed by data.

3.2 Neutral pion P_T and P_Z spectra in p-p and p-Pb collisions

Published results [9] for π^0 transverse and longitudinal momentum spectra in proton-proton collisions at $\sqrt{s} = 7$ TeV will be shown in Sect. 3.2.1. Results from comparison with proton-lead collisions data collected at $\sqrt{s_{NN}} = 5.02$ TeV will be presented in Sect. 3.2.2. A test for the Feynman scaling rule will be shown in Sect. 3.2.3.

In previous published results [6, 7], π^0 s were reconstructed from energy and incident position of their decay photons in events with one incident photon in each tower ("Type I" event), while events with two incident photons in the same tower ("Type II" event) were rejected. In this analysis Type II events were included for the first time. The inclusion of Type II events allowed to extend the coverage in the $P_T - P_Z$ phase space as shown in figure 3.

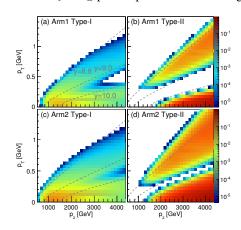


Figure 3. Acceptance map of π^0 s detected by LHCf in $P_T - P_Z$ phase space for Arm1 (top figures) and Arm2 (bottom figures). Type II π^0 s (right figures) extended the phase space coverage up to $P_T \sim 1$ GeV, while with only Type I events the maximum detectable P_T was ~ 0.5 GeV.

3.2.1 P_T and P_Z distributions at $\sqrt{s} = 7$ TeV

The data in p-p collision at $\sqrt{s} = 7$ TeV were obtained during LHC Fills #1104, #1107, #1112 and #1117 from May 15 to 22, 2010. The integrated luminosity was 2.67 nb⁻¹ for Arm1 and 2.10 nb⁻¹ for Arm2

Measured P_T spectra are compared with predictions from different MC interaction models in several regions of the rapidity y, as shown in figure 4. QGSJET II-04 model [12] shows a good overall agreement with LHCf data, while EPOS-LHC [13] has a good agreement at low $P_T(<0.5~{\rm GeV})$ and shows a slightly harder distribution for $P_T>0.5~{\rm GeV}$. Both models are based on the parton-based Gribov-Regge approach [14, 15] and they are tuned using ALICE, ATLAS, CMS and TOTEM data. The prediction of SIBYLL 2.1 [16] agrees well with the data for 8.8 < y < 9.2 and $P_T<0.4~{\rm GeV}$ while its absolute yield is about half that of the data for y>9.2. The predictions of DPMJET 3.06 [17] and PYTHIA 8.185 [18, 19] are compatible with data for 9.0 < y < 9.8 and $P_T<0.2~{\rm GeV}$, while for $P_T>0.2~{\rm GeV}$ they become significantly harder than both LHCf data and the predictions of other models.

 P_Z spectra in several P_T intervals are presented and compared with models predictions in figure 5. There is again an overall agreement between LHCf data and the QGSJET prediction, especially for 0.0 $< P_T < 0.2$ GeV. The EPOS prediction is compatible with data for $P_Z < 2$ TeV, while showing a hard slope for $P_Z > 2$ TeV in all P_T regions. The predictions by DPMJET and PYTHIA agree with data for $P_T < 0.2$ GeV and $P_Z < 1.6$ TeV, while showing a harder distribution for the higher P_Z regions. SIBYLL predicts a smaller production of π^0 s for $P_T < 0.2$ GeV and shows a similar behaviour with DPMJET and PYTHIA with increasing P_T .

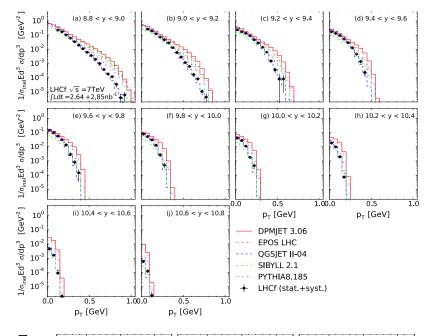


Figure 4. Transverse momentum distribution of neutral pions in several rapidity regions at $\sqrt{s} = 7$ TeV: coloured histograms show results from Monte Carlo simulations with different hadronic interaction models, black points show experimental data.

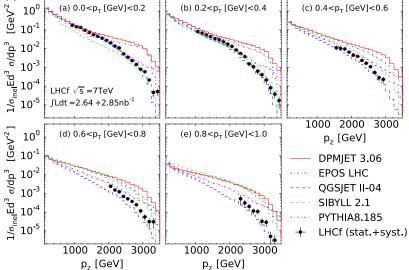


Figure 5. Longitudinal momentum distribution of neutral pions in several P_T regions at $\sqrt{s} = 7$ TeV: coloured histograms show results from Monte Carlo simulations with different hadronic interaction models, black points show experimental data.

3.2.2 Nuclear modification factor

During the p-Pb run in 2013, only Arm2 detector was installed because of two different reasons: to minimize interference with ATLAS ZDC operation and because MC simulations predicted too high multiplicity of neutrons in Pb-remnant side. As a consequence, Arm2 were installed on p-remnant side for almost all the time. Data recorded on p-remnant side in LHC Fills #3478 and #3481 on January 21-22, 2013 were used in the analysis, with an integrated luminosities of 0.63 nb⁻¹.

Proton-lead collisions are very important to understand nuclear effects on forward particles production, because the first collision of an air-shower happens between the primary cosmic ray (a proton

or a nucleus) and a nucleus of the atmosphere. Comparing spectra in p-p and p-Pb collisions it is possible to estimate nuclear effects; in particular it is possible to measure the nuclear modification factor (R_{pPb}) which is defined as

$$R_{pPb} \equiv \frac{\sigma_{inel}^{pp}}{\langle N_{coll} \rangle \sigma_{inel}^{pPb}} \frac{E d^3 \sigma^{pPb} / dp^3}{E d^3 \sigma^{pp} / dp^3}$$

where $Ed^3\sigma^{pp}/dp^3$ and $Ed^3\sigma^{pPb}/dp^3$ are the inclusive cross sections of π^0 production in p-p and p-Pb collisions at 5.02 TeV respectively, and $\langle N_{coll} \rangle$ is the mean number of binary nucleon-nucleon collisions. $\langle N_{coll} \rangle$ is estimated as 6.9 from MC simulations. Since there are no data in p-p collisions at $\sqrt{s} = 5.02$ TeV, the distribution $Ed^3\sigma^{pp}/dp^3$ is obtained interpolating results at 7 TeV and 2.76 TeV in p-p collisions. Nuclear modification factor as a function of P_T in different rapidity regions is shown in figure 6. MC models predict a strong suppression $(R_{pPb} \sim 0.1)$ over all range of P_T and it is consistent with the LHCf data within the uncertainties.

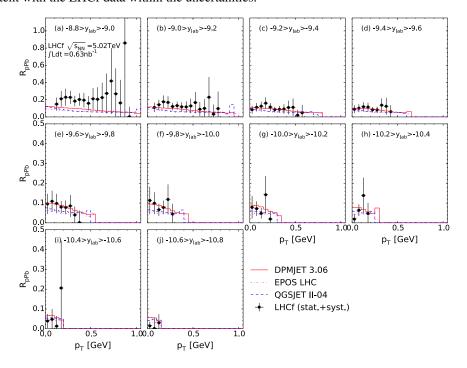


Figure 6. Nuclear modification factor for neutral pions in p-Pb collisions at 5.02 TeV: coloured histograms show results from Monte Carlo simulations with different hadronic interaction models, experimental data are shown as black points.

3.2.3 Test of Feynman scaling

The Feynman scaling rule [20] states that the production cross sections of secondary particles as a function of the Feynman-x variable $(X_F \equiv 2P_Z/\sqrt{s})$ were independent of the incident energy in the forward region. If the Feynman scaling holds, the differential cross section $(X_F/\sigma_{inel})(d\sigma/dX_F)$ should be independent of the center of mass energy for $X_F \gtrsim 0.2$.

This scaling rule was tested with the LHCf data at $\sqrt{s} = 2.76$ TeV and $\sqrt{s} = 7$ TeV. The test was possible only in the P_T range 0.0 GeV $< P_T < 0.4$ GeV because there are no data at 2.76 TeV outside this range. The X_F distributions of data at 7 TeV and 2.76 TeV are shown in figure 7 for two different P_T ranges. The comparison of the two data sample indicates that the Feynman scaling holds at the $\sim 20\%$ level.

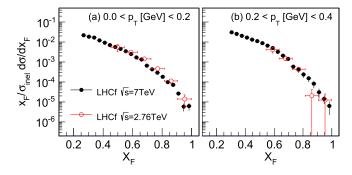


Figure 7. X_F distributions of LHCf data at $\sqrt{s} = 7$ TeV (black points) and $\sqrt{s} = 2.76$ TeV (red points) in two P_T ranges. The distributions are consistent at the $\sim 20\%$ level.

3.3 Photon inclusive energy spectra in p-p collisions at \sqrt{s} = 13 TeV

Proton-proton collisions at $\sqrt{s} = 13$ TeV were produced for the first time in 2015 at LHC. LHCf had a dedicated low-luminosity run from 9th to 13th of June 2015, with an instantaneously luminosity of $0.5 \div 1.6 \times 10^{29} \text{cm}^{-2} \text{s}^{-1}$. The preliminary analysis of photon spectrum was done on the data obtained during the LHC Fill #3855. The integrated luminosity was 0.23 nb^{-1} for both Arm1 and Arm2. The preliminary inclusive energy spectra of photon produced in p-p collisions at $\sqrt{s} = 13 \text{ TeV}$ are presented in figure 8 for two pseudo-rapidity ranges together with the predictions of QGSJET, EPOS, SIBYLL, DPMJET and PYTHIA hadronic interaction models. The LHCf data lie between MC predictions but there is not an unique model with a good agreement in the whole energy range.

4 Conclusions

LHCf experiment performed measurements on forward production of neutral particles in proton-proton and proton-lead collisions at LHC accelerator. These measurements are necessary to calibrate hadronic interaction models used in cosmic rays physics to understand the development of atmospheric showers. LHCf performed measurements at $\sqrt{s} = 0.9$, 2.76, 7 and 13 TeV in p-p collisions and at $\sqrt{s_{NN}} = 5.02$ TeV in p-Pb collisions.

In this paper latest published results about π^0 analysis (p-p at 2.76 and 7 TeV, p-Pb at 5.02 TeV) and neutron analysis (p-p at 7 TeV) were presented. Preliminary results from neutron analysis indicate a large amount of high-energy hadrons in the very forward region, predicted only by QGSJETII-03 model. π^0 P_T and P_Z measured spectra show a better agreement with QGSJET II-04 (in the whole range) and EPOS-LHC (for $P_T < 0.5$ GeV and $P_Z < 2$ TeV). π^0 analysis in p-Pb indicates a strong suppression of π^0 production compared to p-p case; nuclear modification factor from data is consistent with MC models predictions within errors. Feynman scaling hypotesis was tested comparing π^0 spectra at 7 and 2.76 TeV; the results indicate that the rule holds at the $\sim 20\%$ level. Preliminary photon energy spectra indicate that there is not an unique model representing the data, but the measured data lie between MC predictions.

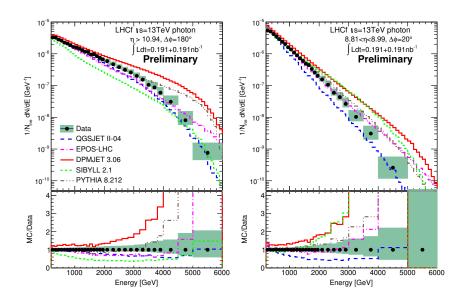


Figure 8. Preliminary inclusive photon energy spectrum for pseudorapidity $\eta > 10.94$ (left figure) and $8.81 < \eta < 8.99$ (right picture). Data are represented by black point while MC prediction from several models are represented by coloured histograms. Green shaded area represents statistical+systematic errors of data.

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