

# Breakup of ${}^8\text{B}$ on ${}^{58}\text{Ni}$ at energies around the Coulomb barrier and the astrophysical $S_{17}(0)$ factor revisited

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**Abstract.** Calculations of breakup and direct proton transfer for the  ${}^8\text{B}+{}^{58}\text{Ni}$  system at energies around the Coulomb barrier ( $E_{\text{B,lab}}=22.95$  MeV) were performed by the continuum-discretized coupled channels (CDCC) method and the coupled-reaction-channels (CRC) method, respectively. For the  ${}^7\text{Be}+{}^{58}\text{Ni}$  interaction, we used a semimicroscopic optical model potential (OMP) that combines microscopic calculations of the mean-field double folding potential and a phenomenological construction of the dynamical polarization potential (DPP). The  ${}^7\text{Be}$  angular distribution at  $E_{\text{lab}}=25.75$  MeV from the  ${}^8\text{B}$  breakup on  ${}^{58}\text{Ni}$  was calculated and the spectroscopic factor for  ${}^8\text{B} \rightarrow {}^7\text{Be}+p$  vertex,  $S_{\text{expt}} = 1.10 \pm 0.05$ , was deduced. The astrophysical  $S_{17}(0)$  factor was calculated equal to  $20.7 \pm 1.1$  eV·b, being in good agreement with the previously reported values.

## 1 Introduction

The study of nuclear reactions with the proton-halo exotic  ${}^8\text{B}$  nuclei is of great interest for nuclear astrophysics in view of the problem of stellar nucleosynthesis and the production of high-energy neutrinos in the Sun. In particular, the breakup of  ${}^8\text{B}$  in the field of heavy targets can provide information about an inverse process, the proton radiative capture by  ${}^7\text{Be}$ , which occurs in the Sun at energies about 20 keV. The  ${}^8\text{B}+{}^{58}\text{Ni}$  system has been extensively studied both experimentally and theoretically by different research groups around the world. We studied the breakup of  ${}^8\text{B}$  in the field of  ${}^{58}\text{Ni}$  with the realistic  ${}^7\text{Be}$  core-target potential calculated in the semi-microscopic OMP. The analysis of the breakup, transfer and elastic-scattering cross sections allowed us to obtain the experimental spectroscopic factor  $S_{\text{expt}}$  and extract the astrophysical  $S_{17}(0)$  factor by using the ANC method. A comparison was made with calculations performed by using the Woods-Saxon potentials previously reported [1].

## 2 Elastic scattering calculations

For the  ${}^7\text{Be}+{}^{58}\text{Ni}$  interaction an optical potential of the form  $U = V_F + V_P + iW + V_C$  was used, where  $V_F$  is a double folded potential,  $iW = i(W_V f(x_V) + W_D f(x_D))$ ,  $V_C$  represent the absorption and Coulomb potential, respectively and  $V_P$  is the DPP, implemented by S. A. Goncharov [2]:

$$V_P = \alpha(E)W_V(E)f(x_V) + \beta(E) \cdot 4W_D(E)\frac{df(x_D)}{dx_D}, \quad (1)$$

where the Woods-Saxon form factor  $f(x_{V,D})$  was used. To calculate  $V_F$ , we used for the projectile an empirical density model that was constructed on the basis of the global parametrization, described in [3]; for the target, we considered an appropriately normalized empirical charge density in a three-parameter modified Fermi form with values taken from [4]. The radial part of the potential were calculated by using the CDM3Y6-Paris nucleon-nucleon effective interaction [5]. To find the OMP parameters for the  $p+{}^7\text{Be}$  and  $p+{}^{58}\text{Ni}$  interactions, the systematics proposed in [6] was used.

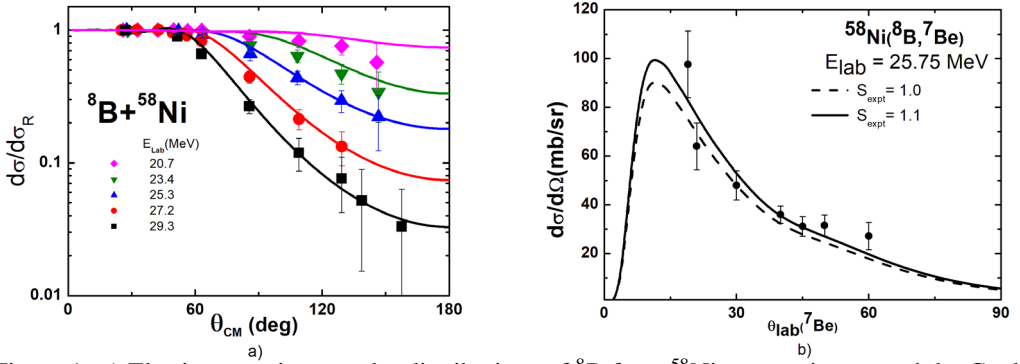


Figure 1: a) Elastic scattering angular distributions of  ${}^8\text{B}$  from  ${}^{58}\text{Ni}$  at energies around the Coulomb barrier calculated by the CDCC in comparison with the experimental data [8]. b) The differential breakup cross section of the  ${}^8\text{B}+{}^{58}\text{Ni}$  reaction at  $E_{\text{lab}}=25.75$  MeV. CDCC calculations performed with two values of the spectroscopic factors are shown in comparison with the experimental data of [9].

The elastic scattering angular distributions were calculated for the  ${}^8\text{B}+{}^{58}\text{Ni}$  system at laboratory energies  $E_{\text{lab}} = 20.7, 23.4, 25.3, 27.2$  and  $29.3$  MeV using the FRESKO code [7]. Figure 1 (a) shows the calculations in comparison with the data reported in [8]. Our results agree well with the data, particularly at energies above the Coulomb barrier. At energies below the barrier, the calculations slightly differ from the experimental data. Table 1 shows the DPP parameters used to fit the  ${}^7\text{Be}+{}^{58}\text{Ni}$  system data of [8]. For the energies 20.7 and 23.4 MeV, values of  $r_{V,D} = 1.2$  and  $a_{V,D} = 0.4$  fm were used, while for the rest of the energies, the values  $r_{V,D} = 1.28$  and  $a_{V,D} = 0.45$  fm showed the best fit to the data. The values  $W_V = 90$  and  $W_D = 5$  MeV were kept constant for all energies.

Table 1: DPP parameters used for the  ${}^7\text{Be}+{}^{58}\text{Ni}$  system.

$E_{\text{lab}}$ [MeV]	$\alpha$	$\beta$	$E_{\text{lab}}$ [MeV]	$\alpha$	$\beta$	$E_{\text{lab}}$ [MeV]	$\alpha$	$\beta$
20.7	-1.25	15.0	25.3	-0.28	0	27.2	-0.08	0
23.4	0	0	25.75	-0.23	0	29.3	0	0

### 3 Breakup and transfer analysis

We assume a cluster structure of  ${}^8\text{B} = p \oplus {}^7\text{Be}$ . The valence proton has an orbital angular momentum  $l$ , thus having total angular momentum relative to the core  $\mathbf{J} = \mathbf{l} + \mathbf{s}$ . In the case of breakup above

Table 2: Breakup and reaction cross sections for the  ${}^8\text{B}+{}^{58}\text{Ni}$  system.

$E_{\text{lab}}$ (MeV)	$\sigma_{\text{bu th}}$ (mb)	$\sigma_{\text{R th}}$ (mb)	$\sigma_{\text{R exp}}^a$ (mb)
23.4	194.56	382.26	365±50
25.3	198.92	606.92	515±50
27.2	204.80	797.80	827±45
29.3	209.24	978.24	1007±40

<sup>a</sup> Experimental data taken from ref. [8]

the Coulomb barrier, the excited states of proton in the continuum were represented by 167 bins with orbital angular momenta  $l=0-4$  up to energies of 6 MeV. Figure 1 (b) shows the CDCC calculations of the  ${}^8\text{B}$  breakup differential cross section for the  ${}^8\text{B}+{}^{58}\text{Ni}$  reaction with spectroscopic factors  $S_{\text{expt}} = 1.0$  and  $S_{\text{expt}} = 1.1$ , respectively, corresponding to the  ${}^8\text{B} \rightarrow {}^7\text{Be}+p$  vertex. The results are compared with the data of [9].

We calculated the direct proton transfer in the  ${}^{58}\text{Ni}({}^8\text{B}, {}^7\text{Be}){}^{59}\text{Cu}$  reaction, which can contribute to the  ${}^{58}\text{Ni}({}^8\text{B}, {}^7\text{Be})$  reaction cross section. Excited states of  ${}^{59}\text{Cu}$  up to  $E_x = 3.580$  MeV were taken into consideration. The calculation showed that proton stripping provides less than 3% of the total  ${}^7\text{Be}$  emission cross sections.

Table 2 shows the breakup and reaction cross sections calculated for the  ${}^8\text{B}+{}^{58}\text{Ni}$  system at energies around the Coulomb barrier in comparison with the data taken from Ref. [8]. The reaction cross sections were obtained by fitting the elastic scattering angular distributions using CDCC calculations. An accepted value of the spectroscopic factor for the  ${}^8\text{B} \rightarrow {}^7\text{Be}+p$  vertex,  $S_{\text{expt}}=1.10 \pm 0.05$ , allowed us to estimate the ANC,  $C^2 = 0.54 \pm 0.03 \text{ fm}^{-1}$ , and the astrophysical  $S_{17}(0)$  factor to be equal to  $20.7 \pm 1.1 \text{ eV}\cdot\text{b}$ , which are in good accordance with the previously published results [1].

## 4 Conclusions

We have performed CDCC calculations of the elastic scattering, breakup, direct proton transfer, and reaction cross sections for the  ${}^8\text{B}+{}^{58}\text{Ni}$  system at energies around the Coulomb barrier. All cross sections were calculated by using the  ${}^7\text{Be}+{}^{58}\text{Ni}$  semi-microscopical optical model potential containing the folding and DPP parts. The direct proton transfer contribution to the reaction cross section is about 3%. The astrophysical  $S_{17}(0)$  factor equal to  $20.7 \pm 1.1 \text{ eV}\cdot\text{b}$  was calculated using the ANC method, being in good agreement with previously reported values.

## References

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