

Heavy light tetraquarks from Lattice QCD

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Abstract. We present preliminary results from a lattice calculation of tetraquark states in the charm and bottom sector of the type $ud\bar{b}\bar{b}$, $us\bar{b}\bar{b}$, $ud\bar{c}\bar{c}$ and $sc\bar{b}\bar{b}$. These calculations are performed on $N_f = 2 + 1 + 1$ MILC ensembles with lattice spacing of $a = 0.12$ fm and $a = 0.06$ fm. A relativistic action with overlap fermions is employed for the light and charm quarks while a non-relativistic action with non-perturbatively improved coefficients is used in the bottom sector. Preliminary results provide a clear indication of presence of energy levels below the relevant thresholds of different tetraquark states in the double bottom sector while a scattering state is observed in the charm sector.

1 Introduction

The discovery of the resonances $Z_b(10607)$ & $Z_b(10650)$ by BELLE [1] in 2012 has shown the existence of multi-quark exotic states in the bottom sector. This has generated a lot of interest in exploring its hadronic structure with the leading candidate being that of a tetraquark state. A tetraquark state, first employed by Jaffe [2] in the context of light scalar mesons and later for exotic spectroscopy [3], is colour neutral state formed as a bound system of diquarks and antidiquarks. The tetraquark structure has been recently employed to identify favourable flavour, spin channels in the bottom sector. However, the tools employed in such searches are typically sum rule type calculations. A first principles approach of lattice QCD is more desirable for such searches and in the past year, two lattice studies [4–6] have identified a promising channel with the flavor content $ud\bar{b}\bar{b}$. The calculation in Ref [4] computed a potential of two heavy static antiquarks in presence of two light quarks in lattice QCD. This was then used to solve a coupled non-relativistic Schrödinger equation to find a binding energy of $\Delta E = 90^{+43}_{-36}$ MeV in the $ud\bar{b}\bar{b}$ channel with $I(J^P) = 0(1^+)$. In a later extension to this work in Ref [6], a resonance prediction was made by searching for the poles in the S and T matrices decaying in two B mesons. The work in Ref [5] also confirmed a presence of a deeply bound level in the flavour channels of $ud\bar{b}\bar{b}$ and $us\bar{b}\bar{b}$ with the binding of 189(10) MeV and 98(7) MeV respectively. In the work presented at the conference, we explore the tetraquarks of type $ud\bar{b}\bar{b}$, as in Ref [5], confirming a presence of a level well below the threshold state and explore other flavour channels.

2 Operator setup

We consider two types of operators here, namely a tetraquark operator with two quarks and two antiquarks with the quantum numbers of the interested state and a two meson operator corresponding

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with equivalent quantum numbers of that of the tetraquark state. The construction of the tetraquark operator employs a product of a diquark and antidiquark as suggested by Jaffe [3]. The diquarks and antidiquarks are constructed with the so called ‘‘good diquark’’ [3] prescription. We would like to construct a tetraquark operator with $I(J^P) = 0(1^+)$ in the diquark-antidiquark picture with two antibottom and two light quarks in the following configuration:

$$(uq) \rightarrow (\bar{\mathbf{3}}_c, 0, F_A), \quad (\bar{b}\bar{b}) \rightarrow (\mathbf{3}_c, 1, F_S),$$

$$D(x) = u_\alpha^a(x) (C\gamma_5)_{\alpha\beta} q_\beta^b(x) \bar{b}_\kappa^a(x) (C\gamma_i)_{\kappa\rho} \bar{b}_\rho^b(x). \quad (1)$$

where the braces indicate the Color, Spin and Flavor (C,S,F) degrees of freedom. For the case of light quarks F_A indicates antisymmetric flavour combination which in this case will be in $\mathbf{3}_f$. For the double antibottom quarks, the flavour symmetry is manifestly symmetric F_S . In the light diquark, the flavour $q \in (d, s, c)$ allows for studying different flavours of tetraquark states. The tetraquark operator shown on line two, indicated by $D(x)$ (keeping consistent notation with Ref [5]), is constructed by taking a dot product of the aforementioned diquark and antidiquark in color space. A two meson operator with the quantum numbers as $I(J^P) = 0(1^+)$ can be constructed with different flavours $q \in (d, s, c)$ as :

$$M^d(x) = B^+(x)B^{0*}(x) - B^0(x)B^{+*}(x) \rightarrow (q = d)$$

$$M^s(x) = B^+(x)B_s^*(x) - B_s^0(x)B^{+*}(x) \rightarrow (q = s)$$

$$M^c(x) = B^+(x)B_c^*(x) - B_c^+(x)B^{+*}(x) \rightarrow (q = c) \quad (2)$$

With these operators, correlation functions were computed on the lattices which will be described in the next section.

3 Lattice setup

The calculations presented at the conference are performed on MILC ensembles employing HISQ gauge action and $N_f = 2 + 1 + 1$ flavours. The ensemble parameters are shown in Table 1. For both the ensembles the charm and strange quark masses are tuned to their physical values, while the ratio m_s/m_l is fixed to 5. The details can be found in Ref [7].

Table 1. Ensemble parameters used in this work

V	β	a(fm)	m_π (MeV)	$m_\pi L$	N_{confs}
$24^3 \times 64$	6.00	0.1207(14)	305	4.54	236
$48^3 \times 144$	6.72	0.0582(5)	319	4.51	70

For the propagator computation, we have used wall sources ¹ and the configurations were gauge fixed in Coulomb gauge and the links were then HYP smeared. In the valence sector, we employ an overlap action where the details can be found in Refs [8–10]. The use of overlap action eliminates $O(a)$ lattice artifacts. In addition with the use of a multimass algorithm, a range of input bare masses can be accommodated. The strange quark mass is tuned by equating the fictitious pseudoscalar $\bar{s}s$ to 685 MeV [10]. The charm quark mass was tuned by setting its spin averaged kinetic mass $(m_{\eta_c} + 3m_{J/\psi})/4$ to its physical value [8] and the bare values of the $am_c = 0.529, 0.290$ were used for the $24^3 \times 64$ and $48^3 \times 144$ lattices respectively.

¹Propagators were also computed on point sources although the results are not shown here.

The bottom sector employs a NRQCD action as shown in Ref [11]. In this set up, all terms upto $1/M_0^2$ and leading term in $1/M_0^3$ are included in the action where $M_0 = am_b$ corresponds to the bare bottom quark mass. The interaction part of NRQCD Hamiltonian includes $O(a)$ improved derivatives and also six improvement coefficients $c_1..c_6$. The details of the action can be found in Ref [9]. We use the non-perturbative determination of these coefficients as done by the HPQCD collaboration [12] for the coarser lattice. For the finer lattices, they are set to their tree level values. The bottom quark mass was tuned by setting its lattice spin averaged mass of (1S) bottomonium:

$$\overline{M}_{\text{kin}}(1S) = \frac{3}{4}aM_{\text{kin}}(\Upsilon) + \frac{1}{4}aM_{\text{kin}}(\eta_b), \quad M_{\text{kin}} = \frac{a^2 p^2 - (a\Delta E)^2}{2a\Delta E} \quad (3)$$

to its experimental value. The kinetic mass is computed as shown in the right equation above. The scale is set independently using the Ω baryon is found to be consistent with the MILC determinations using the r_1 parameter.

4 Results

In obtaining the ground states of the correlation functions, we employ the variational method [13, 14]. With the tetraquark operator $D(x)$ and the two meson operator $M(x)$, we compute a correlator matrix:

$$C_{ij}(t) = \sum_x \langle 0 | O_i(x, t) O_j^\dagger(0, 0) | 0 \rangle \quad O_i(x, t) \in \{D(x, t), M(x, t)\} \quad (4)$$

The correlator matrix is a 2×2 matrix and we solve a generalised eigenvalue problem to obtain the principle (ground state) correlation function as:

$$C_{ij}(t + \Delta t)v_j(t) = \lambda(t)C_{ij}(t)v_j(t) \quad m_{\text{eff}}(t) = -\frac{\log \lambda(t)}{\Delta t} \quad (5)$$

where $\lambda(t), v_j(t)$ are the eigenvalue and eigenvectors from the solution of the GEVP. It is convenient to construct effective masses as shown in right equation above.

4.1 Results for $ud\bar{b}\bar{b}$

The results for the $ud\bar{b}\bar{b}$ tetraquarks are shown in Fig 1. In the case of the $ud\bar{b}\bar{b}$, the threshold states are those of two free mesons namely B and B^* . The plot in the left panel shows the effective masses of the two free meson states obtained from the correlator $C = C_B C_{B^*}$ (shown in green). The data in blue corresponds to the lowest level of the GEVP solution of the 2×2 correlator matrix constructed from the correlation functions of the operators mentioned in section 2. The solution of the GEVP yields two levels, the excited level is found to be noisy and is not shown here. With the use of wall sources, the ground states are seen to approach a plateau from below. The data in blue provides a clear indication of a level below the effective mass of the threshold correlator with the binding indicated on the plot. The results are also seen to be noisy for $t > 25$. The calculation has been extended to smaller pion masses and the preliminary results are shown in the right panel. At the pion mass close to the SU(3) point, a comparison can be with the results at the finer lattice spacing (right panel, data in blue) where the binding energies results are seen to be consistent. As the pion mass is lowered, the binding is seen to get deeper, albeit with higher uncertainties. This trend is found to be consistent with observations made in Ref [5].

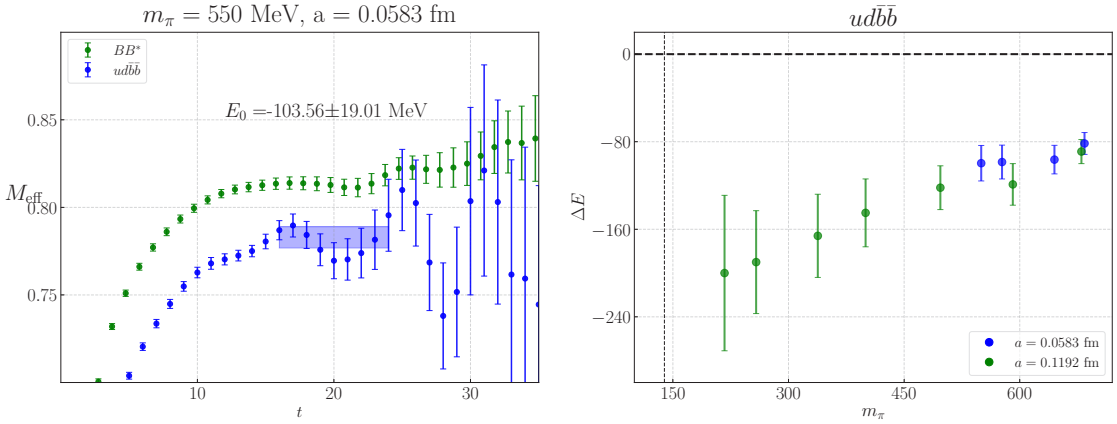


Figure 1. Left panel: Preliminary results for $ud\bar{b}\bar{b}$ tetraquark state. See text for the description of effective masses. Right panel : Binding energies for $a = 0.0583$ fm (data in blue) and $a = 0.1207$ fm (data in green)

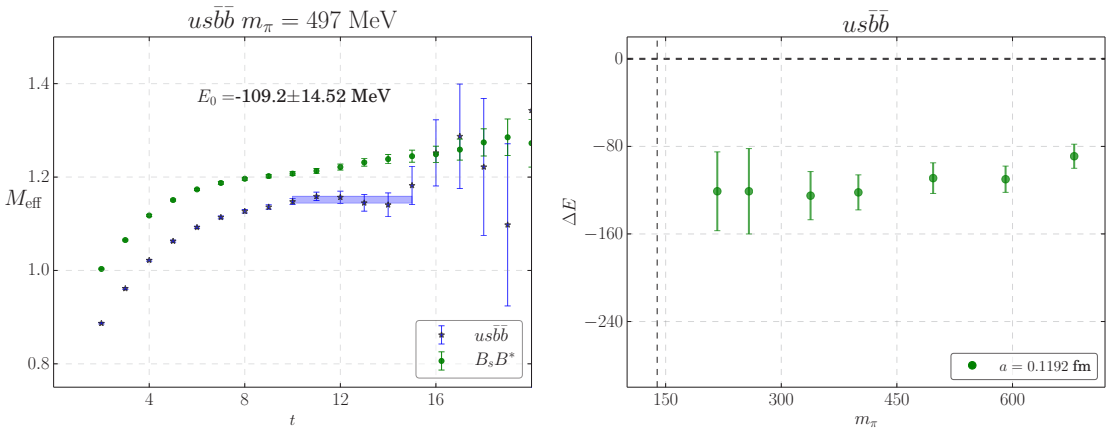


Figure 2. Preliminary results for $us\bar{b}\bar{b}$. Left panel : Effective energies for the threshold state and the lowest level of the GEVP solution. Right panel : Summary of binding energies at lower pion masses at $a = 0.1207$ fm.

4.2 Results for $us\bar{b}\bar{b}$

The results for $us\bar{b}\bar{b}$ tetraquarks are shown in Fig 2. The threshold here is that of B_s meson and B^* meson. As before, the left panel indicates effective masses of the product of the correlators of B_s and B^* (shown in green) and the lowest level of the GEVP solution (data in blue). A clear indication of a level below the threshold state is seen for $m_\pi = 497$ MeV with the binding indicated on the plot. The results shown here are computed on $24^3 \times 64$ lattice with $a = 0.1207$ fm. The right panel in fig 2 show the results of the pion mass dependence of the binding energies where the slope of the binding energies w.r.t pion masses is not as pronounced in comparison with the of $ud\bar{b}\bar{b}$ which

possibly indicates a shallower binding at the physical pion. These results however are preliminary and will be improved upon in a later publication.

4.3 Results for $sc\bar{b}\bar{b}$

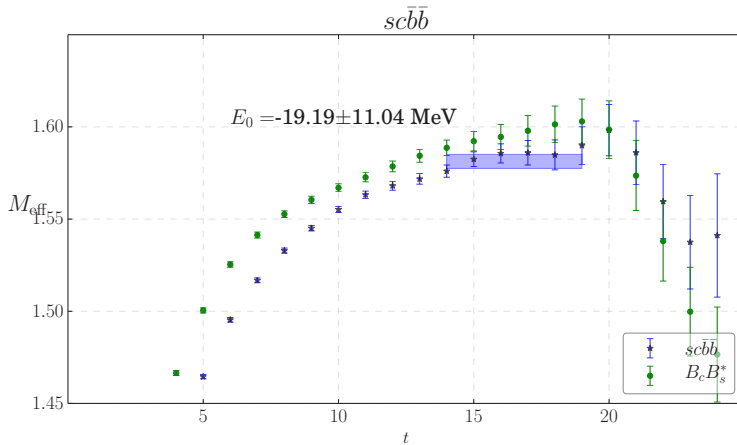


Figure 3. Preliminary results for $sc\bar{b}\bar{b}$ with all flavours at physical quark mass.

We have also included the tetraquark state $sc\bar{b}\bar{b}$ in our calculation with the results shown in the Fig 3. As before the data in green indicates the effective mass of the threshold correlator which in this case is two free mesons namely B_c meson and B_s^* meson. We also note that this calculation was done with quark masses for all flavours at their physical value. The main systematic in case will be the lattice spacing dependence of the binding energy which is currently ongoing. Due to the shallow result of the binding energy, study of finite volume effects in this case may also be important.

4.4 Results for $ud\bar{c}\bar{c}$

The charm analogue of the doubly bottom tetraquark state is the $ud\bar{c}\bar{c}$ state. The two meson thresholds in this case are the D and D^* mesons. The results of this calculation are shown in Fig 4. As before the data in green are the effective masses of the threshold correlator and the data in blue are those of the $ud\bar{c}\bar{c}$. The results in this case are seen lie well above the threshold clear indicating no binding. This trend in this case is in contrast to what is seen with the $ud\bar{b}\bar{b}$. The results presented here are at lattice spacing $a = 0.0583 \text{ fm}$ and at heavier pion masses. The extension to lower pion masses and another lattice spacing is currently underway.

5 Conclusions and outlook

In the work shown here, we have explored heavy light tetraquarks in the bottom and charm sector. The results on $a = 0.1207 \text{ fm}$ is the progress since the conference. In the $ud\bar{b}\bar{b}$ sector, we find a very clear indication of deeply bound levels as the pion mass is lowered. The results shown here are

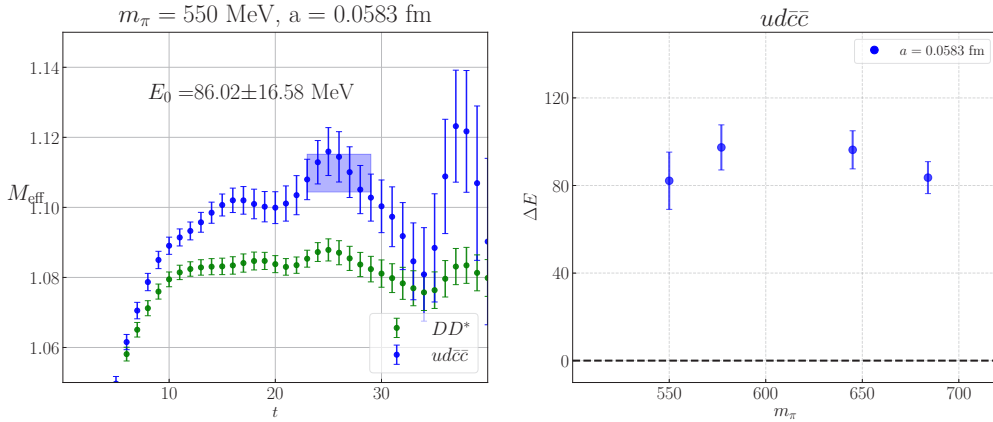


Figure 4. Preliminary results for $ud\bar{c}\bar{c}$. The color notation is the same as in previous plots.

preliminary and with added statistics are subject to change. The results at lattice spacings $a = 0.0583$ fm and $a = 0.1207$ fm are seen to be consistent indicating no significant lattice spacing dependence. The work is being currently extended to improve statistics and include results lower pion masses for atleast one more lattice spacing. The results on $us\bar{b}\bar{b}$ are also seen to provide a clear indication of levels below the threshold state at various pion masses. The trend to the physical point is not as pronounced as that of $ud\bar{b}\bar{b}$ indicating that the binding may be shallower. This has been noted in Ref [5]. The effective mass of $sc\bar{b}\bar{b}$ state is seen to be close to the threshold at the physical values of the quark masses. The results on $ud\bar{c}\bar{c}$ are seen to lie well above the threshold states.

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