

A strange program for LHCb

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Abstract. The LHCb experiment is expanding its physics program towards studies of rare decays of strange particles. In this talk, we reviewed the published results by LHCb in $K_S^0 \rightarrow \mu^+ \mu^-$ and $\Sigma^+ \rightarrow p \mu^+ \mu^-$ decays, as well as sensitivity studies and prospects for other strangeness decays.

1 Introduction

The Standard Model (SM) of particle physics is very successful at describing accelerator data, but it is known to be incomplete as it fails to explain gravity, dark matter, or neutrino masses. For this reason, dynamics Beyond the Standard Model (BSM) is being searched for. Direct searches at LHC are limited to few TeV by the maximum energy of the accelerator, while precise measurements of flavour decays can a priori access higher energy scales. This is particularly true for models with new sources of flavour violation, in which the CKM matrix is not the only source of flavour and CP violation. In such context, $s \rightarrow d$ quark transitions are central, because they have the strongest CKM suppression. Examples of these are kaon oscillations, rare decays such as $K \rightarrow \pi \nu \bar{\nu}$, being searched for in NA62, or $K_S^0 \rightarrow \mu^+ \mu^-$, which is being searched for at LHCb. Here we review the status and prospects for strangeness decays at the LHCb.

2 The LHCb experiment and strangeness trigger

The LHCb detector at the LHC [1] is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing b or c quarks. It has collected an integrated luminosity of 1 fb^{-1} at $\sqrt{s} = 7 \text{ TeV}$ in 2011 and 2 fb^{-1} at $\sqrt{s} = 7 \text{ TeV}$ in 2012. In 2015 it started its second Run (hereafter Run-II) of data taking, and expects to collect up to 8 fb^{-1} $\sqrt{s} = 13 \text{ TeV}$ by the end of 2018. The LHCb Upgrade, scheduled to start by 2021, is expected to collect 50 fb^{-1} at $\sqrt{s} = 14 \text{ TeV}$ by the end of 2029. Finally, an *Expression of Interest* for a Phase-II Upgrade, aiming to collect $\sim 300 \text{ fb}^{-1}$ has been submitted to the LHC Council [2].

Due to its optimization for b and c decays, the trigger efficiency for strangeness decays in its first years of data taking has been very low (1-2% for kaons with dimuon final-states, to be compared with 80-90% for dimuon B -decays). Yet, the strangeness cross-section at LHC is very large, with $\sim 10^{13}$ K_S^0 produced in the LHCb acceptance per fb^{-1} , and it allowed the publication of world best results in rare decays of strange particles. The flexibility of the software trigger (HLT) allowed increasing significantly the trigger efficiency for strangeness decays to dimuons for Run-II data [3], as shown in Table 1. Since in the LHCb upgrade there will be no hardware trigger (L0) and all the trigger steps will be software based, the trigger efficiency is expected to be as high as for B decays.

Table 1. Approximate trigger efficiencies for benchmark $s \rightarrow d \mu^+ \mu^-$ transitions at LHCb for 2012 and Run-II. Numbers obtained from [3]. The efficiency in 2011 data is lower than in 2012 data, being approximately $\sim 1\%$. The maximum efficiency that can be achieved, taking into account L0 constraints, is 30%.

Channel	$\epsilon(2012)\%$	$\epsilon(\text{Run-II})\%$
$K_S^0 \rightarrow \mu^+ \mu^-$	2.9	25
$\Sigma^+ \rightarrow p \mu^+ \mu^-$	0.83	11
$K_S^0 \rightarrow \pi^0 \mu^+ \mu^-$	2.6	24

3 Results

3.1. $K_S^0 \rightarrow \mu^+ \mu^-$

In the SM, the $K_S^0 \rightarrow \mu^+ \mu^-$ decay is long-distance dominated, with subdominant short-distance

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contributions from Z-penguin and W-box. Yet, the long-distance contribution is still very small in absolute terms, and the decay rate is very suppressed [4] [5] [6]:

$$BR(K_S^0 \rightarrow \mu^+ \mu^-)_{SM} = (5.18 \pm 1.50_{LD} \pm 0.18_{SD}) \times 10^{-12}.$$

Therefore, small BSM contributions and BSM/SM interferences can sizably modify the branching fraction. This has been proven to be the case in leptoquark models [7] [8] as well as supersymmetric models [9], despite of other experimental constraints.

The 2011 LHCb data allowed improving by a factor of thirty [10] the world best upper limit prior to LHC [11]. Recently, the analysis of the full Run-I data has been published [12], corresponding to an integrated luminosity of 3fb^{-1} and improved reconstruction, trigger and selection. The analysis uses a Boosted Decision Tree (BDT) to separate signal from combinatorial background and dedicated particle identification to separate $K_S^0 \rightarrow \mu^+ \mu^-$ decays from misidentified $K_S^0 \rightarrow \pi^+ \pi^-$ decays. The mass shift on the K_S^0 peak caused in the $\pi^\pm \rightarrow \mu^\pm$ misidentification is also crucial for the analysis, since it separates the $K_S^0 \rightarrow \pi^+ \pi^-$ peak away from the signal, and only the tail of the distribution leaks into the region in which genuine $K_S^0 \rightarrow \mu^+ \mu^-$ decays are expected. This can be seen in Figure 1 (top), where the invariant mass distribution of the dimuon candidates is shown, for one of the most sensitive BDT bins. A simultaneous maximum likelihood fit of the dimuon invariant mass spectrum of all the BDT bins is performed, yielding no significant signal. A Bayesian upper limit (see Figure 1 bottom) is derived by integrating the likelihood times a flat prior on the BR, resulting in an upper limit

$$BR(K_S^0 \rightarrow \mu^+ \mu^-)_{EXP} < 8(10) \times 10^{-10}$$

The extrapolated sensitivity as a function of the integrated luminosity times trigger efficiency is shown in Figure 2. The trigger efficiency for Run-II is estimated to be 25%, while for the LHCb upgrade is expected to be close to 100%. It can be seen that the LHCb upgrade can explore a good part of the 10^{-11} region, and a Phase-II upgrade could get very close to the SM prediction.

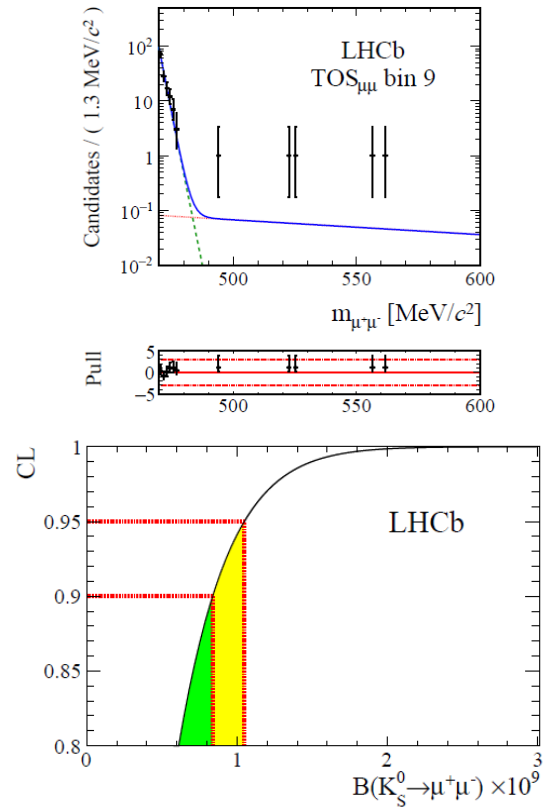


Figure 1. *Top:* Invariant mass distribution of the dimuon candidates in one of the most sensitive BDT bins. The green dashed line shows the contribution from misidentified $K_S^0 \rightarrow \pi^+ \pi^-$ decays. The red dotted line shows the contribution from combinatorial background. No significant signal (which would have shown as a nearly-Gaussian peak at $\sim 500 \text{ MeV}/c^2$) has been found. *Bottom:* cumulative function of the posterior probability of the branching fraction. The intersections with (upper limit) 90% and 95% credible intervals are shown with dashed red lines.

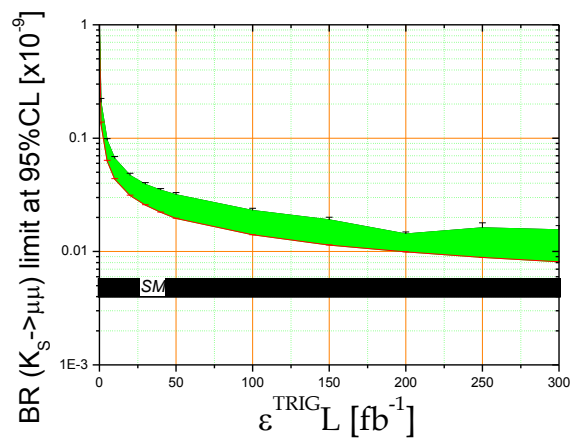


Figure 2. Expected LHCb sensitivity to $BR(K_S^0 \rightarrow \mu^+ \mu^-)$ as a function of the integrated luminosity (times trigger efficiency). The trigger efficiency for Run-II is estimated to be $\sim 25\%$, while for LHCb upgrade is expected to be close to

the 100%. The black band shows the SM prediction with its uncertainty. The width of the green band shows an estimated uncertainty on the sensitivity prediction, including model systematics and statistical subtraction of the $K_L^0 \rightarrow \mu^+\mu^-$ contribution.

3.2 $\Sigma^+ \rightarrow p\mu^+\mu^-$

The $\Sigma^+ \rightarrow p\mu^+\mu^-$ decay is a flavour-changing neutral-current process, allowed only at loop level in the SM. The process is dominated by long-distance contributions for a predicted branching fraction of [13]

$$1.6 \times 10^{-8} < BR(\Sigma^+ \rightarrow p\mu^+\mu^-)_{SM} < 9.0 \times 10^{-8},$$

while the short-distance SM contributions are suppressed at a branching fraction of about 10^{-12} . Evidence for this decay was seen by the HyperCP experiment [14] with a measured branching fraction

$$BR(\Sigma^+ \rightarrow p\mu^+\mu^-)_{EXP}^{HyperCP} = (8.6_{-5.4}^{+6.6} \pm 5.5) \times 10^{-8},$$

which is compatible with the SM prediction. Remarkably, the three observed decays have almost the same dimuon invariant mass of $m_{\chi^0} = 214.3 \pm 0.5$ MeV/ c^2 , which lies close to the kinematic limit. This result is sometimes referred to as “the Hyper-CP anomaly”, since it suggests a potential BSM particle at a mass of ~ 214 MeV/ c^2 .

The LHCb collaboration has recently published [15] a preliminary search for the $\Sigma^+ \rightarrow p\mu^+\mu^-$ decay, showing a strong evidence for this decay with a significance of four standard deviations (see Figure 3, top). The dimuon invariant mass distribution does not show evidence for a low mass resonance, disproving the HyperCP-anomaly (see Figure 3, bottom). This measurement was based on Run-I data with no trigger path dedicated to this channel. In the time between the FCCP workshop and the write-up of the proceedings, the search has been published into a paper [16], providing in addition a branching fraction measurement

$$BR(\Sigma^+ \rightarrow p\mu^+\mu^-)_{EXP}^{LHCb} = (2.1_{-1.2}^{+1.6}) \times 10^{-8}.$$

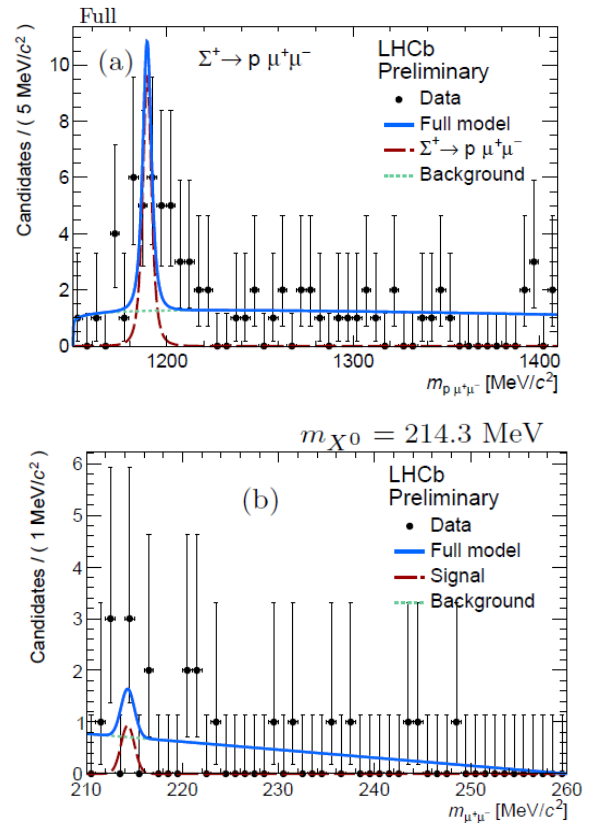


Figure 3. *Top:* Invariant mass distribution of the $\Sigma^+ \rightarrow p\mu^+\mu^-$ candidates. The green dotted line represents the background, and the red dashed line the $\Sigma^+ \rightarrow p\mu^+\mu^-$ signal, with a significance of four standard deviations. *Bottom:* dimuon mass distribution of the candidates. The green dotted line represents an exponential continuum, and the red dashed line a potential $X^0 \rightarrow \mu^+\mu^-$ signal, which is found to be not significant.

4 Sensitivity studies

4.1 $K_S^0 \rightarrow \pi^0\mu^+\mu^-$

The $K_S^0 \rightarrow \pi^0\mu^+\mu^-$ decay is another $s \rightarrow d\mu^+\mu^-$ transition with a very low branching fraction, predicted at the 10^{-9} level in the SM. The closely related $K_L^0 \rightarrow \pi^0\mu^+\mu^-$ decay has been shown to have sensitivity to BSM dynamics, such as Extra Dimensions [17] or leptoquarks [8]. Its sensitivity is limited by the uncertainty on the SM contribution, which is dominated by the experimental precision on the K_S^0 mode, the NA48 measurement of $BR(K_S^0 \rightarrow \pi^0\mu^+\mu^-)$ [18]:

$$BR(K_S^0 \rightarrow \pi^0\mu^+\mu^-)_{EXP} = (2.9_{-1.2}^{+1.5} \pm 0.2) \times 10^{-9}$$

Hence, improvements on the experimental measurement of $BR(K_S^0 \rightarrow \pi^0\mu^+\mu^-)$ translate into improved BSM search potential of $K_L^0 \rightarrow \pi^0\mu^+\mu^-$. The LHCb experiment has evaluated its sensitivity to $BR(K_S^0 \rightarrow \pi^0\mu^+\mu^-)$ [19] with two different reconstruction strategies: explicit requirement of the

$\pi^0 \rightarrow \gamma\gamma$ decay (FULL), and reconstruction based on the dimuon information only (PARTIAL). In both cases the LHCb upgrade is expected to overtake the precision of NA48, as shown in Figure 4. The sensitivity as a function of the dimuon mass has not been evaluated yet.

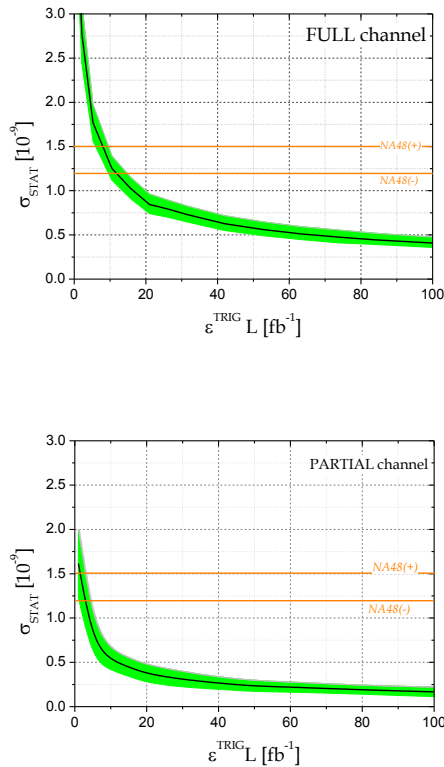


Figure 4. Expected LHCb sensitivity to $BR(K_S^0 \rightarrow \pi^0 \mu^+ \mu^-)$ as a function of the integrated luminosity (times trigger efficiency). The trigger efficiency for Run-II is estimated to

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be $\sim 24\%$, while for LHCb upgrade is expected to be close to 100%.

4.2 $K_S^0 \rightarrow \pi^+ \pi^- e^+ e^-$

Although most of the efforts have been put in dimuon modes, sensitivities to final states with electrons, such as $K_S^0 \rightarrow \pi^+ \pi^- e^+ e^-$, [20] are also under study, showing that an LHCb strangeness program with electrons is also possible.

4.3 Other analyses

It must be noted that the LHCb capabilities are not restricted to the four channels discussed above. Indeed, other measurements such as K^+ mass, $K_S^0 \rightarrow \gamma \mu^+ \mu^-$, kaons from Φ decays, Lepton Flavour Violation, Lepton Universality Violation, and semileptonic decays of hyperons could be feasible and are under study.

4 Conclusions

Results and prospects of the LHCb experiment in rare decays of strange particles have been discussed. These include updated limits in $BR(K_S^0 \rightarrow \mu^+ \mu^-)$ as well as the disproof of the Hyper-CP anomaly in $\Sigma^+ \rightarrow p \mu^+ \mu^-$ decays. Future prospects have also been discussed, showing the potential to reach $BR(K_S^0 \rightarrow \mu^+ \mu^-)$ at nearly the SM level, as well as to overtake NA48 in the determination of $BR(K_S^0 \rightarrow \pi^0 \mu^+ \mu^-)$.

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