

An Investigation on the Effects of Different Stratifications on Negatively Buoyant Jets

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Abstract. Negatively buoyant jets develop when fluids are released upwards into a lighter fluid or, vice versa, downwards into a heavier fluid. There are many engineering applications, such as the discharge, via submerged outfalls, of brine from desalination plants into the sea. Some concerns are raised about the potential negative environmental impacts of this discharge. The increase in salinity is the major cause for environmental impact, as it is very harmful to many marine species. The diffusers for brine discharge are typically inclined upwards, to increase the path before the brine reaches the sea bottom, as it tends to fall downwards driven by negative buoyancy. The negatively buoyant jet that develops conserves axisymmetry only when released vertically, so that it is not possible to use the well-known equations for axisymmetric jets. The main target of this paper is to investigate on a laboratory model the effects of different stratifications on the features of negatively buoyant jets. This has been done via a LIF (Light Induced Fluorescence) technique, testing various release angles on the horizontal and densimetric Froude numbers. Except for the initial stage, a different widening rate for the upper boundary and the lower boundary has been highlighted.

1 Introduction

Negatively buoyant jets (NBJs) develop when fluids are released upwards into a lighter fluid or, vice versa, downwards into a heavier fluid.

NBJs are implied in many fields, from the passive containment cooling system of nuclear power plants (Wang et al., 2017 [1]) to the volcanic eruption columns (Kaminski et al. 2005 [2]), see Ferrari and Querzoli, 2010 [3], for a list. In particular, the discharge of brine from desalination plants is one of the most investigated (see, for instance, Cipollina et al., 2004 [4]; Kikkert et al., 2007 [5], Besalduch et al., 2013 [6], Ferrari and Querzoli, 2015 [7], Crowe et al., 2016 [8], Nikiforakis et al., 2017 [9]), because of the concerns linked to the potential negative environmental impacts of the brine discharge (Campbell et al., 2005 [10]): this is mainly due to the increase in salinity which brine discharges impose on the receiving body, harmful to a lot of marine species, in particular to the *Posidonia Oceanica* (see Latorre, 2005 [11]; Sánchez-Lizaso et al., 2008 [12]).

The typical configuration for marine brine discharge is a submerged outfall equipped with upward-inclined diffusers: in this way, the path followed by the brine is increased, as brine is driven downwards by negative buoyancy and the path is delimited by the sea bottom (e.g. Roberts and Toms, 1988 [13], and Jirka, 2008 [14]). Even if, as above showed, there have been several investigations about NBJs, the high complexity of the physical phenomenon is not yet fully understood. As a matter of fact, NBJs conserve axisymmetry only as long

as they are vertically released, so it is not possible to use the well-known equations for axisymmetric jets. As this lack in axisymmetry is mainly due to the different conditions of stratifications in the upper and lower boundaries of NBJs, the main target of this paper is to investigate the different effects of these stratifications on the features of NBJs.

2 Dimensional Analysis

When a NBJ is released from an upward-inclined diffuser, the flow is driven from two sources: one of momentum and one of negative buoyancy. As a consequence (see Figure 1), if the initial part is driven mostly by the momentum (more “simple-jet-like”), as far as the fluid moves away from the outlet, the buoyancy prevails and the axis bends down (more “plume-like”). As stated by List, 1979 [15], when the receiving fluid is at rest, the only parameters characterizing the boundary conditions are the initial flux of volume Q , the initial flux of momentum M and the initial flux of buoyancy per unit mass B : performing a dimensional analysis, two length scales l_Q and l_M can be obtained. The former, i.e. l_Q , gives the magnitude of the distance from the orifice where the jet becomes independent from the initial conditions:

$$l_Q = \frac{Q}{\sqrt{M}} \quad (1)$$

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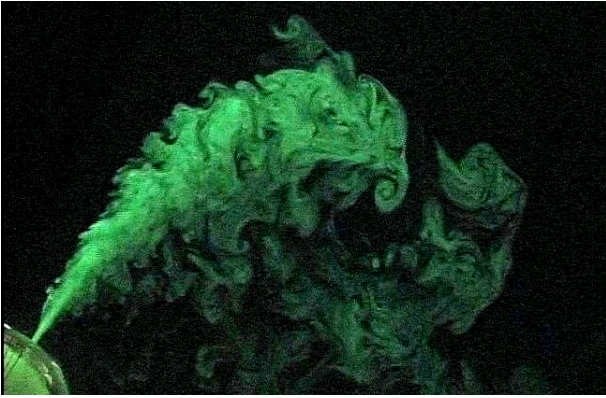


Fig 1. Instantaneous concentration field of a NBJ (densimetric Froude number Fr : 14.8; release angle to the horizontal θ : 55°).

In the case of a circular outlet, such as the one employed in the experiments presented here, l_Q becomes:

$$l_Q = \sqrt{\pi/4}D \cong D \quad (2)$$

The latter, i.e. l_M , gives the magnitude of the distance from the orifice where the plume-like behaviour prevails:

$$l_M = \frac{M^{3/4}}{\sqrt{B}} \quad (3)$$

From the ratio of l_Q to l_M , a non-dimensional number, the Richardson number Ri , measuring the relative importance of buoyancy and momentum, can be obtained:

$$Ri = \frac{l_Q}{l_M} \quad (4)$$

The most important non-dimensional parameter for NBJ is usually considered the densimetric Froude number Fr , i.e. the ratio of inertia to buoyancy:

$$Fr = \frac{U}{\sqrt{g \frac{\rho_{DISC} - \rho_{REC}}{\rho_{REC}} D}} \quad (5)$$

where U is the mean initial jet velocity, ρ_{DISC} the discharged fluid density, ρ_{REC} the receiving fluid density and D the outlet diameter. Fr and Ri are inversely proportional:

$$Fr = \sqrt[4]{\frac{\pi}{4}} Ri^{-1} \quad (6)$$

In addition to Fr , the non-dimensional parameters influencing NBJs are the Reynolds number

$$Re = \frac{UD}{\nu} \quad (7)$$

where ν is the kinematic viscosity of the brine, and the angle to the horizontal θ . As θ measures the misalignment between M and B , a NBJ will be axisymmetric only if θ is $\pm 90^\circ$.

3 Experimental set-up and data elaboration

Experiments were carried out in the Hydraulics Laboratory of the DICAAR, University of Cagliari. In particular, a 30 cm wide, 21 m long flume was employed. The flume was filled with tap water to simulate the sea at rest and a solution of water, sodium-sulphate (a low corrosive salt employed to increase the density of the solution) and fluoresceine (a fluorescent salt) was used to simulate the brine. The solution flew (through a pipe) from a constant head tank (supplied by a closed hydraulic circuit) to a cylindrical vessel, simulating a portion of a pipe laid down the sea bottom, with a sharp-edged orifice on its lateral wall. This was done to simulate a widely employed configuration for the marine outfall design (see, for instance, Wright et al., 1982 [16]; Avanzini, 2006 [17]).

To allow measurements via a Digital Image Analysis (DIA) technique (see below), the NBJ middle vertical section was lighted by a 4 mm thick light sheet and a video-camera, orthogonal to the light sheet, filmed the experiments at a frequency of 25 Hz and a resolution of 720x576 pixels. The flow was investigated by means of LIF (Light Induced Fluorescence), a DIA technique: in Fig. 1, the reader can see an example of an instantaneous image, where a light NBJ is above a dark background. With the employed concentrations, the concentration field is proportional to the light intensity, so a quantitative measure of the concentration is allowed. Digital Image Analysis (DIA) techniques have many advantages, if compared to the traditional techniques employing physical point probes: DIA are non-intrusive and quasi-continuous in space, as every pixel on the camera sensor works as a single probe (see Ferrari, 2017 [18] for a review). The LIF codes were designed in-house by the Hydraulics Group of the University of Cagliari, which has a considerable expertise in the development of DIA techniques to measure concentrations (Ferrari and Querzoli, 2010 [3], Ferrari et al., 2016 [19]), velocity (Besalduch et al., 2013 [6] and 2014 [20], Ferrari et al, 2017 [21]) and accelerations (Ferrari et al, 2007 [22], Ferrari and Rossi, 2008 [23], Ferrari et al, 2008 [24], Lardeau et al., 2008 [25], Rossi et al, 2009 [26]) in the environmental (such as in the present paper), biomedical (Badas et al., 2015 [27], Espa et al, 2012 [28]) or civil and industrial fields (Badas et al., 2017 [29]).

Experiments we were performed with $\theta = 45^\circ \div 90^\circ$ and $Fr = 4.9 \div 40.2$: by this way, the whole field of practical interest was investigated. Moreover, Q was

kept constant and high enough to have a Re larger than its critical value for the apparatus.

Under the hypothesis of ergodicity, the statistics of mean and variance of the concentration field were obtained by time averaging. On Fig. 2 and 3, the mean and the variance of the concentration field are shown. The mean concentration is normalized by the initial concentration, the variance by the squared maximum concentration. The jet axis was assumed to be coincident with the locus of the concentration maxima on jet cross-sections. Detailed information on these parameters, as well as on the experimental set-up and data elaboration, can be found on Ferrari and Querzoli, 2010 [3] and on Ferrari et al., 2016 [19].

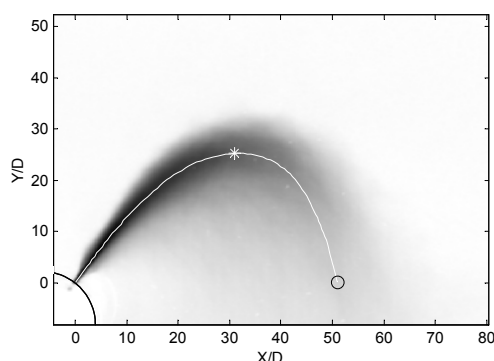


Fig. 2. Map of the non-dimensional mean concentration for the same NBJ of Fig. 1; the white line is the jet axis, the white asterisk is the point of maximum height and the black circle is the impact point.

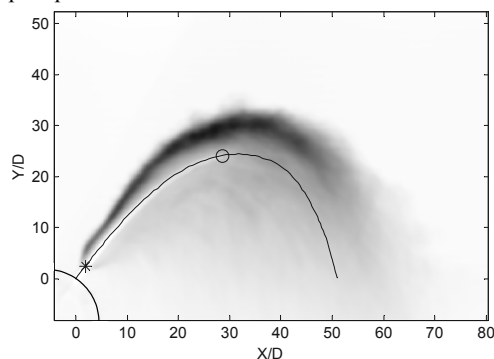


Fig. 3. Map of the variance of the concentration, non-dimensionalised by the square of the maximum concentration, for the same NBJ of Fig. 1; the black line is the jet axis, the black asterisk is the point of onset of the Kelvin-Helmholtz instability and the black circle is the point of transition to turbulence of the whole NBJ.

4 Results

This paragraph is organised as follows: first of all, a discussion about the different stratifications, stable on the upper region and unstable in the lower region of the NBJ, will be given; then, the concentration profiles, orthogonal to the NBJ axis, at various distances from the outlet will be shown, in order to discuss their shape, with particular regard to the Gaussian shape hypothesis often employed in the study of these phenomena; lastly, the focus will be on the widening features of the NBJ upper and lower region.

4.1. The different stratifications on the NBJ boundaries

From Figures 1, 2 and 3, we can see that NBJs released upwards are characterised by an initial region with a width similar to the outlet diameter D : afterwards, the onset of Kelvin-Helmholtz (KH) instabilities (in the region pointed out by a black asterisk on Figure 3) causes a sudden widening, as the consequent billows drag external fluid into the NBJ, resulting in mixing and dilution. This dilution, as demonstrated by the reduction in the grey intensity along the axis on Figure 2, grows along the trajectory followed by the jet.

On Figure 2, we can see that the mean concentration field of a NBJ is symmetrical in the first part of the jet, where the flow is driven essentially by M ; then, as B becomes relevant, the span-wise sections become neither symmetrical nor self-similar (see chapter 4.2). As a consequence, above the axis there is a narrow region characterized by high upwards concentration gradients, whilst below the axis there is a wide region of smooth variations. This is due to the fact that, on the upper boundary, the local stable stratification (M and B have different directions) allows a full development of compact KH billows, followed by their breaking; on the contrary, the local unstable stratification on the lower boundary (M and B have the same direction) tends to transform the growing waves into downward plumes.

On Figure 3, the variance field shows the lowest values in the core region of the NBJ, close to the axis, values that increase towards the upper and lower boundaries. The stable stratification at the upper boundary of the NBJ has the consequence that the maximum variance values are located at the upper, external region, where the KH waves develop upwards but rapidly break, because of the downwards B contrasting their growth. On the opposite, the unstable stratification at the lower boundary of the NBJ has the consequence that the downwards-developing, plume-like structures are driven away from the NBJ core by the buoyancy, with more distributed but less intense concentration fluctuations.

4.2 Shape and Gaussianity of the concentration cross-sectional profiles

As previously stated, when the receiving fluid is at rest, a NBJ conserves axisymmetry only if it is released vertically (M and B have the same direction) but, if the discharge is done with a different θ , the NBJ rapidly loses its symmetry. This is clearly visible in Figure 4 and 5, where cross-sectional profiles of C/C_0 (the non-dimensional mean concentration C , normalized by the initial mean concentration C_0), are shown; R/D is the non-dimensional coordinate orthogonal to the jet axis (positive values of R/D identify the lower portion of the jet, while negative ones identify the upper portion). The legend shows the non-dimensional distance S/D , measured from the outlet along the axis, where the profiles are measured. The different shape between the upper and lower region increases with increasing S/D .

The measured points on the concentration profiles (circles, asterisks, crosses and triangles) are shown together with the best-fit Gaussian curves (in a least-squares sense), are shown on Figure 6; the NBJ upper region is on the left-hand side (negative R/D), the lower one on the right-hand side (positive R/D).

This graph highlights another difference between the upper and lower regions: the concentration profiles have a Gaussian shape, both in the lower and upper part, in the initial region of the NBJ, close to the outlet (solid lines in Figure 6), even if they are not identical. As S/D grows, the upper region conserves a Gaussian shape (dashed and dotted lines on the left), whilst the lower region loses its Gaussian shape (as shown by the dashed and dotted lines on the right that do not correctly fit the experimental data). When the NBJ becomes to an end (last section, $S/D = 50$), both the profiles have a Gaussian shape (dash-dot lines), but again with different coefficients.

We can consequently state that an inclined NBJ cannot be considered as a Gaussian phenomenon.

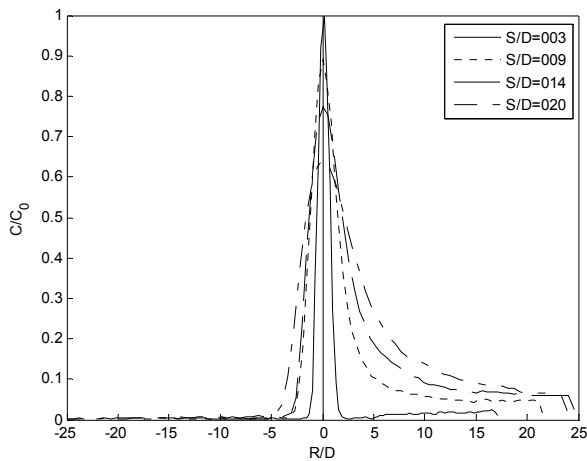


Fig. 4. cross-sectional non-dimensional concentration profiles for a NBJ with $\theta = 45^\circ$, $Fr = 8.0$; positive values of R/D identify the lower portion of the NBJ, negative ones identify the upper portion; S/D is the non-dimensional distance from the outlet measured along the axis; $S/D = 3 \div 20$.

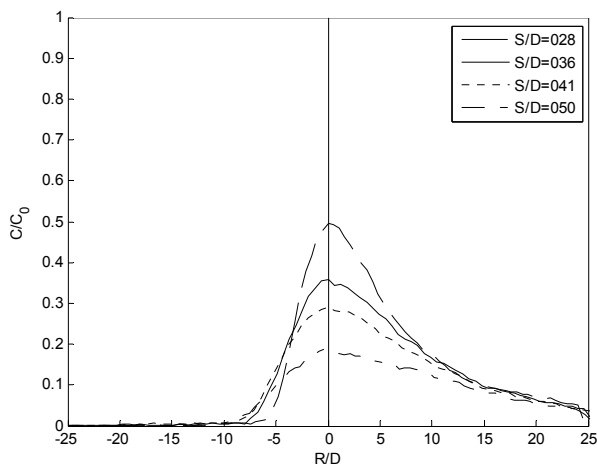


Fig. 5. cross-sectional non-dimensional concentration profiles for the same NBJ of Fig. 4; $S/D = 28 \div 50$.

4.3 NBJ widening features

To study the widening features of inclined NBJs, Figures 7, 8, 9 and 10 show the experimental data for the widening of the upper boundary (red asterisks) and lower boundary (blue circles) of NBJs, taking into account all the θ in the range $45^\circ \div 70^\circ$ and Fr of 30.8 (Fig.7), 23.8 (8), 14.8 (9) and 8.0 (10). The data for the highest θ have been not taken into account because at those θ s NBJs experiences the so called re-entrainment. Re-entrainment is the mixing of the NBJ fluid with itself instead of with external fluid; under these conditions, there is not a sharp separation between the uprising and descending branches of the NBJs, as shown by Ferrari and Querzoli, 2010 [3], and consequently it is not possible to define the lower boundaries of the NBJs. The widening is computed as the R/D where the concentration decreases to a fraction $1/e$ of the maximum concentration on that particular profile.

The NBJs with the highest Fr (30.8 on Figure 7) experience the largest values of S/D , because they are the lightest and, as a consequence, they cover the longest distance before falling down to the sea bottom; of course, exactly the opposite happens for the heaviest NBJs ($Fr = 8.0$ on Figure 10), that cover the shortest path. For S/D lower than the length scale l_M (in the range between 29, for $Fr = 30.8$ and 7.6, for $Fr = 8.0$), the widening is similar for the upper and lower regions, as the NBJ is still behaving in a simple-jet-like matter, that with a similar linear growth for both the upper and lower region. When S/D becomes much larger than l_M , the lower region, because of the unstable stratification, begins to widen with a faster rate than the upper region, whose widening is slowed down by the stable stratification (see paragraph 4.1).

To study the law of widening rate of the upper and lower region of these NBJs, the upper region data have been fitted, on a least-square sense, with the linear law of eq. 8, the lower boundary with both the eq. 8 and the 5/3 power law of eq. 9:

$$\frac{R}{D} = c \frac{S}{D} + \frac{1}{e} \quad (8)$$

$$\frac{R}{D} = \kappa \left(\frac{S}{D} \right)^{5/3} + \frac{1}{e} \quad (9)$$

Obviously, the constant $1/e$ has been added because, as a consequence of the chosen definition for the widening, the NBJ width at the outlet has this value for all the Fr and θ .

Looking at Figures from 7 to 10, equation 8 seems to correctly fit the data relative to the upper region. Focusing on the lower region, both equations 8 and 9 are able to catch some of the data features, with some lacks in precision for the lowest and highest S/D . More in details, equation 8 strongly overestimates the initial widening of the lower region and underestimates the final growth.

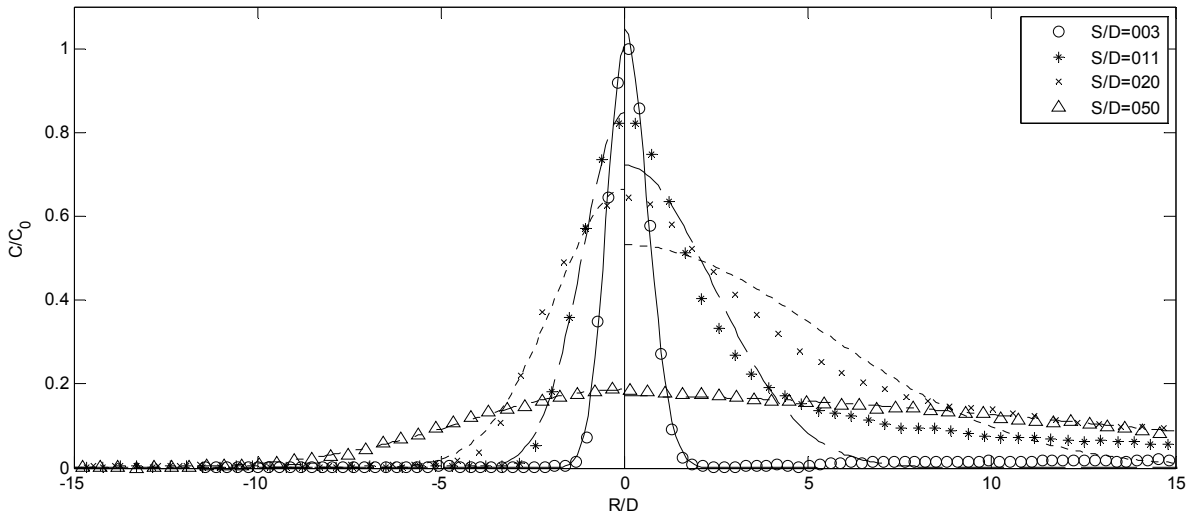


Fig. 6. Interpolation with best-fit (least mean squares) Gaussian curves of cross-sectional non-dimensional concentration profiles: on the left the upper region of the NBJ, on the right the lower region of the NBJ (same NBJ as in Figures 4 and 5).

Equation 9, instead, slightly underestimates the initial widening (and it better follows the trend of the data for low values of S/D) but it is able to better catch the trend for increasing S/D , particularly for the lowest Fr s. As a consequence, we can state that equation 9 seems to better fit the points with the highest S/D , which could mean that, if the nearest region to the origin is not taken into account, the lower region tends to a faster than linear evolution or, in other words, doesn't behave like a simple jet of a plume (both linear).

It seems, thereby, that the upper boundary of the inclined NBJs tends to develop following a linear law, whilst the lower boundary tends to evolve in a more complex way, nor jet- neither plume-like, in particular when S/D increases and becomes much larger than l_M . On Table 1, the coefficient of eq. 8 and 9 for the curves of Figures 7, 8, 9 and 10 have been reported; the subscript U refers to the upper region, the subscript L to the lower region. The values of the coefficient c_U (representing the rate of growth of the upper part of the inclined NBJs) decrease with decreasing Fr , suggesting that, as the NBJs become heavier, the stable stratification on the upper boundary prevents the Kelvin-Helmholtz billows to develop vertically and, consequently, the NBJs to widen more. On the opposite, the coefficients κ increase with decreasing Fr , suggesting that, as the NBJs become heavier, the unstable stratification on the lower region helps the descending plumes to detach and, as a consequence, the NBJs to widen more. Given a fixed S/D , as Fr decreases (i.e., B increases), the lower region of an inclined NBJ tends to widen more whilst, vice-versa, the upper boundary tends to widen less.

Finally, considering all the showed results, we can state that an incline NBJ cannot be considered as an axisymmetric or Gaussian phenomenon, as some features of simple jets, plumes and more complex ones coexist: the upper region tends to evolve in a simple-jet-like way, following a linear law widening rate; the lower boundary in a more complex manner, following a power law widening rate. The investigation on inclined NBJs can be considered, consequently, a complicated one, because we are not allowed to employ the well-known

integral equations usually adopted to forecast the behaviour of axisymmetric jets or plumes.

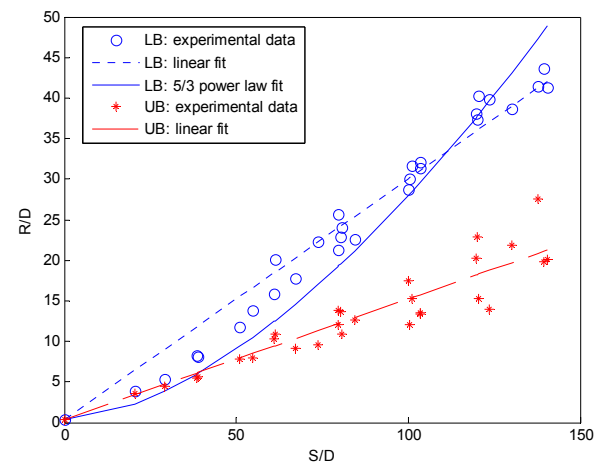


Fig. 7. Experimental data for the widening of the upper boundary (UB) and lower boundary (LB) of NBJs with $Fr = 30.8$ and $\theta = 45^\circ\text{--}70^\circ$.

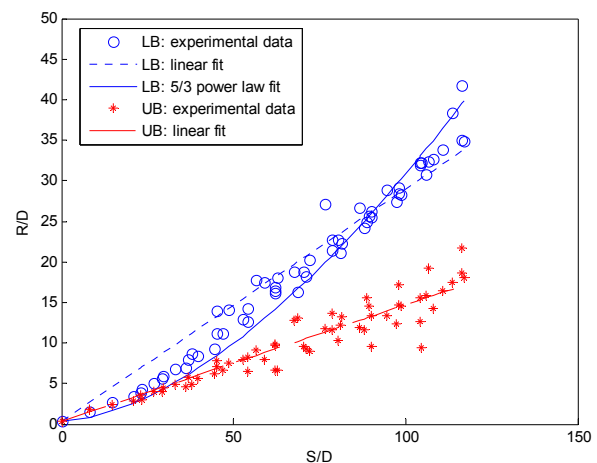


Fig. 8. Experimental data for the widening of the upper boundary (UB) and lower boundary (LB) of NBJs with $Fr = 23.8$ and $\theta = 45^\circ\text{--}70^\circ$.

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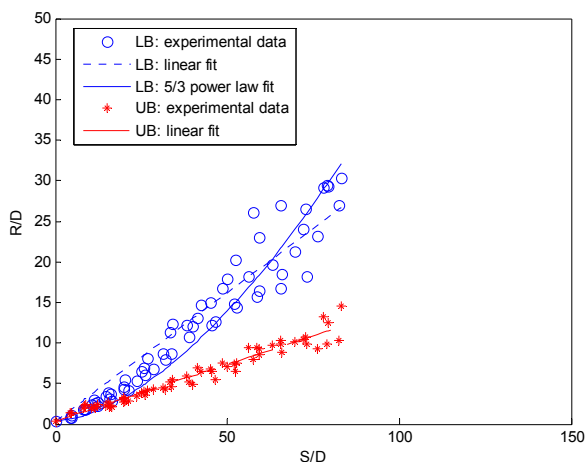


Fig. 9. Experimental data for the widening of the upper boundary (UB) and lower boundary (LB) of NBJs with $Fr = 14.8$ and $\theta = 45^\circ\text{--}70^\circ$.

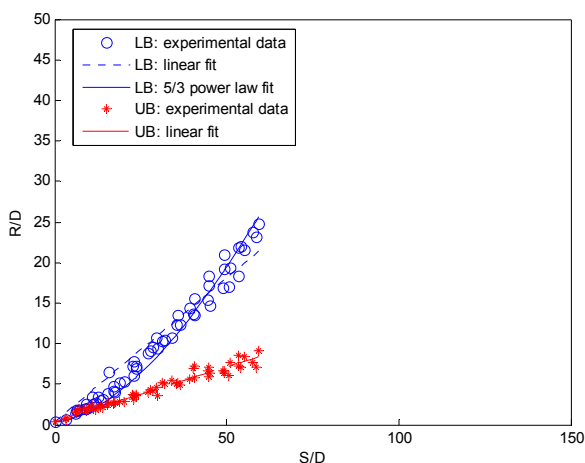


Fig. 10. Experimental data for the widening of the upper boundary (UB) and lower boundary (LB) of NBJs with $Fr = 8.0$ and $\theta = 45^\circ\text{--}70^\circ$.

Table 1. coefficients of the eq. 6 and 7 for the data of Figure 5; U = upper boundary, L = lower boundary.

Fr	c_L	κ	c_U
30.8	0.297	0.013	0.149
23.8	0.286	0.014	0.143
14.8	0.317	0.020	0.140
8.0	0.355	0.028	0.136

5 Conclusions

In order to have a better insight into the effects of stable and unstable stratification on NBJs released with a certain angle to the horizontal θ , an experimental investigation have been carried out in the Hydraulic Laboratory of the DICAAR, University of Cagliari. A non-intrusive Digital Image Analysis technique, namely LIF (Light Induced Fluorescence), was in-house developed and employed. The experimental parameters, in particular θ and the densimetric Froude number Fr , were chosen in order to investigate the whole field of practical interest for the marine brine discharge.

The results show that the different conditions of stratifications, in particular a stable stratification on the upper region and an unstable stratification on the lower region, force the lower region of the NBJ to widen with a faster rate than the upper region, leading to a strongly asymmetric and non-Gaussian phenomenon. As a consequence, a great care should then be employed when applying the integral equations for axisymmetric jets or plume to inclined NBJs. Moreover, the experimental data have shown that the upper region has a linear widening rate, whilst the lower region tends to widen with a faster 5/3 power law rate.

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