

Volumetric flow rate in simulations of microfluidic devices ⁺

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Abstract. In this work, we examine the volumetric flow rate of microfluidic devices. The volumetric flow rate is a parameter which is necessary to correctly set up a simulation of a real device and to check the conformity of a simulation and a laboratory experiments [1]. Instead of defining the volumetric rate at the beginning as a simulation parameter, a parameter of external force is set. The proposed hypothesis is that for a fixed set of other parameters (topology, viscosity of the liquid, ...) the volumetric flow rate is linearly dependent on external force in typical ranges of fluid velocity used in our simulations. To confirm this linearity hypothesis and to find numerical limits of this approach, we test several values of the external force parameter. The tests are designed for three different topologies of simulation box and for various haematocrits. The topologies of the microfluidic devices are inspired by existing laboratory experiments [3 - 6]. The linear relationship between the external force and the volumetric flow rate is verified in orders of magnitudes similar to the values obtained from laboratory experiments.

1 Introduction

The research group Cell in fluid is interested in modelling of elastic objects in fluid flow in solid channels, in order to simulate the flow of the plasma and blood cells in microfluidic devices. The program is developed within open-source software ESPResSo. The model is described in [2]. The numerical simulation of such a device is faster and cheaper to establish, compare with the fabrication of the real microchip.

The long-time goal of the group is to make models of the microfluid devices capturing Circulating Tumour Cells (CTC). At today's stage of the software's development, some calibration calculations are needed, to compare the results of the numeric simulations with results of laboratory experiments. In order to do so, the input parameters influencing the movement of the cells in the blood plasma have to be the same, such as the viscosity of the liquid, the hematocrit, or the volumetric flow rate.

The volumetric flow rate is closely related with the difference of the pressures on the two ends of the device. In the numerical models, this dependence is present as well, and the difference of the pressures is somehow represented by an external fluid force parameter.

We have very often an information about a volumetric flow rate from a laboratory experiment. So we have to find a way how to introduce a known volumetric flow rate into the simulation. That's why we would like to find a way how to define it in the simulation at the beginning, as an input parameter.

It is possible to define the volumetric flow rate directly as an input parameter. The disadvantage of this approach is that we must define the course of the volumetric flow rate across the device, as well. However, it is quite difficult to obtain such an information from laboratory experiment. That's why we look for another option of defining the input flow rate, through the parameter of external fluid force.

The aim of this article is to verify, whether the volumetric flow rate in the simulations of the microfluidic devices is linearly dependent on the parameter of external force, set at the beginning as the simulation parameter. This approach demands only the knowledge of the external force parameter. If it is the case, the wished volumetric flow rate could be obtained by only one iterative simulation.

Authors in [5], p 3296, in sequence "*Using a simple syringe to control flow rate*" considered different volumetric flow rates in their experiments, and they noticed that "*The elastic PDMS device becoming deformed, and inflated, at high pressures. Consequently, the flow rate is not linearly proportional to the applied driving pressure.*" We do not ignore this feature of the microfluidic devices in our work – our goal is not to show that the volumetric flow rate can be increased infinitely in the laboratory experiments. The aim is to find a way how to define it correctly and rapidly from few calibration simulations.

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2 Methodology

To confirm the hypotheses of the linear relationship between the volumetric flow rate and the parameter of external force, we tested three different topologies. All of them are inspired by existing laboratory microfluidic devices. Each tested topology consisted of cylindrical or spherical obstacles placed in a rectangular box. The simulation parameter of the external force was set to different values, to verify the linear relationship between the volumetric flow rate and the simulation parameter. In the beginning, we tested the volumetric flow rate in the topologies without the Red Blood Cells (RBC). Afterwards, we ran the simulations without RBC, in order to include more realistic haematocrit in the laboratory experiments.

2.1 Cylindrical adhesive topology – capturing of CTCs by antibodies-coated obstacles in micro channel

First tested topology is inspired by works [3] and [4]. Both articles are concerned with a separation of the CTCs. For this purpose, they use a microfluidic device which can capture common CTCs from the patients' blood. The CTC are bigger than ordinary blood cells. While the size of ordinary blood cells ranges from approximately 4 to 18 μm , CTCs, in contrast, are larger and range from 15 to 25 μm in diameter [7]. However, those devices are not designed to isolate the cells by a size-sorting mechanism. Instead, they are using antibodies, spread over the obstacles in microchips. These antibodies capture the CTCs that flow through the channel. The dimensions of the obstacles are designed to stimulate the highest possibility of collisions between the cells and the obstacles.

In [3], the microfluidic devices contain cylindrical obstacles with radius of 50 μm placed in equilateral triangular arrangement, with a 50 μm distance between obstacles and a 50 μm shift after every 3 rows. The device is 100 μm deep. The active capture area of the laboratory device was 19 mm x 51 mm. The volumetric flow rate in the device was about 1-2 ml/h.

In [4], the experiments were made on a GEDI device. It contains cylindrical obstacles with radius of 80 μm . The gaps between the obstacles are 100 μm large, and each subsequent row is shifted by 7 μm . The device is 100 μm deep, 8 mm wide and 25 mm long. The volumetric flow rate maintained in the device was of 1 ml/h.

The chosen topology of the simulation channel is presented in the Fig. 1. The radius of the obstacles is 50 μm , the dimension of the gaps between the obstacles is 50 μm , and the obstacles are placed in an equilateral triangular array. The simulation channel is 100 μm deep, 150 μm wide and 260 μm long. The volumetric flow rate which should correspond to a 1ml/h for a 8 mm wide device, corresponds to 4,47 $\mu\text{m}^3/\mu\text{s}$.

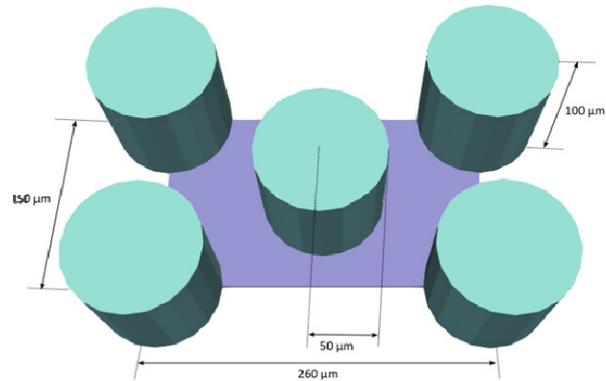


Fig. 1. Description of the simulation channel topology of cylindrical adhesive microfluidic device.

2.2 Cylindrical sorting topology – sorting of the parasites in the blood according to their size

The second tested topology is inspired by a device used in the field detection of blood parasites [5]. The aim is to sort the *T. cyclops* parasites from erythrocytes and leucocytes in human blood, using a pattern of obstacles which helps to sort the objects in blood according to their size. The microfluidic device is composed from 3 parts. First part separates the big leucocytes from the smaller erythrocytes and *T. cyclops*, which have similar dimensions. The second part of the device places the erythrocytes and the parasites close to the boundaries of the device. And the last part of the device keeps the erythrocytes close to the boundaries, while the parasites are placed in the middle stream of the device. The numerical model is inspired by the second part of the microchip.

This second part this microfluidic device consists of cylindrical obstacles, with radius of approximately 10 μm . The gap between the obstacles within one row is approximately 20 μm , and the gap between the rows is approximately 20 μm as well. Each row is shifted by one fifth of the distance between the central axes of the obstacles, which means approximately 8 μm . The depth of the device is 3,5 μm . The total width of the device is about 2 mm. The volumetric flow rates used in the experiment were in range of 2 to 4 $\mu\text{l}/\text{min}$. The device is dimensioned to be used in field, so we considered a haematocrit of a non-diluted blood, which means approximately 40%.

The chosen topology for this device is presented in the Fig.. The radius of the obstacles is 10 μm . The gap between the obstacles in the row and the distance between the rows is 20 μm as well. Each row is shifted by 8 μm . The depth of the device is 4 μm . The simulation box has dimensions 200x40x4 μm . The volumetric flow rate which correspond to 2-4 $\mu\text{l}/\text{min}$, is about 0.7-1.3 $\mu\text{m}^3/\mu\text{s}$.

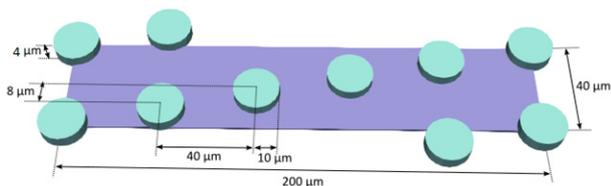


Fig. 2. Dimensions of the simulation channel of the cylindrical sorting topology.

2.2 Spherical adhesive topology – capturing of the CTCs using antibodies-coated magnetic beads

The last topology is inspired by work [6] and [8]. Both of them use magnetic beads to construct the microfluidic device. The idea, as well as in [3] and [4], is not to sort the cells according to their size, but to use antibodies which could capture the cells and keep them glued to the obstacle in the device. The advantage of the magnetic beads is that they do not need to be treated with the antibodies inside the channel during or after its construction, as it is a case for devices with cylindrical obstacles. The magnetic beads can be prepared separately, and after the treatment, they are poured into the device, and kept in the right position by an external magnetic field.

The magnetic beads used in [8] have diameter of 4.5 μm . In the bottom side of the channel, there is a magnetic pattern which defines the topology of the whole device. Once the beads are poured into the microfluidic channel, the magnetic field is activated, and the beads form chains. Those chains are starting at the specific points defined by the pattern at the bottom of the device. This pattern has a form of an equilateral triangular field. Distance between the chains is 40 μm . The depth of the device is 50 μm . It is 0.5 mm large, and 10 mm long. The volumetric flow rate in the channel in experiments is few $\mu\text{L}/\text{min}$.

The simulation box inspired by this geometry (**Fig. 3**) has dimensions of 40x70 μm . It contains five chains of 12 beads with radius of 2.25 μm . Four chains are in the corners, and one is in the middle of the box. This way, they form an array of regular equilateral triangular field. The depth of the device is 50 μm . The volumetric flow rate through this box corresponding to the laboratory experiment is fixed to be 4 $\mu\text{m}^3/\mu\text{s}$.

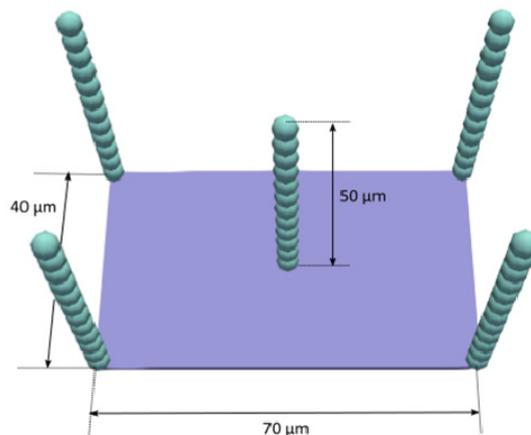


Fig. 3. Dimensions of the simulation channel of the spherical adhesive topology.

3 Measurement of the volumetric flow rate

For each topology, the volumetric flow rate was measured in three different positions. In theory, it should have the same value over any cross section of the channel. The possible differences could have origin in numerical imprecision of the simulation.

The three different positions of volumetric flow rate measurement for the cylindrical adhesive topology are shown in the **Fig. 4**. The evolution of the volumetric flow rate in the three different positions is presented in the **Fig. 5** for maximale and minimal examined value of external force parameter (with multiplier of 1 and 500).

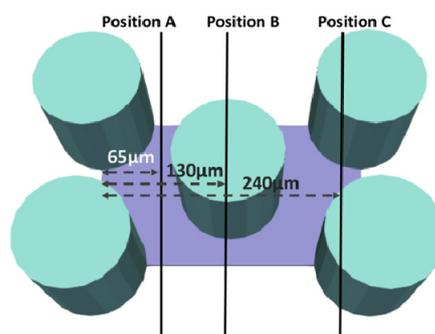


Fig. 4. Dimensions of the simulation channel of the cylindrical sorting topology.

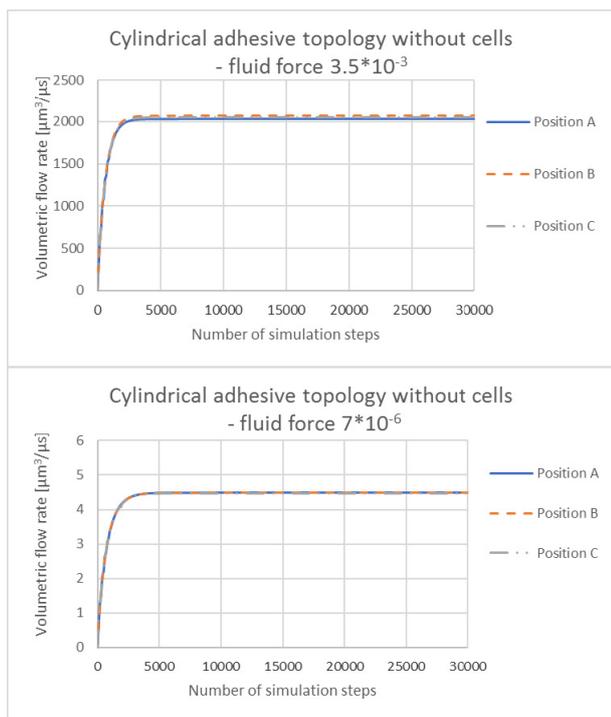


Fig. 5. Cylindrical adhesive topology without cells: comparison between the evolution of the volumetric flow rate measured in three different positions A, B and C. The presented graphs came from the simulation with the smallest and with the highest parameter of external force.

For the cylindrical sorting topology, the volumetric flow rate was measured in three different positions, as well (Fig. 6). The evolution of the volumetric flow rate in the three different positions is presented in the Fig. 7, for the maximal and the minimal examined value of external force parameter – for multipliers 1 and 100. We can observe that for the multiplier of 100, the volumetric flow rate is not the same for all positions, but its values differ by 8%. For the multiplier of 50, its values differ by 3%, and for multiplier of 10, the measured volumetric flow rate is the same for all the three positions.

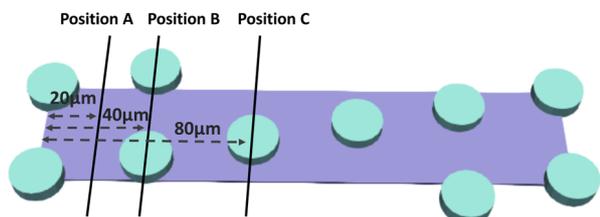


Fig. 6. Three different position of measurement of the volumetric flow rate in the channel with cylindrical sorting topology

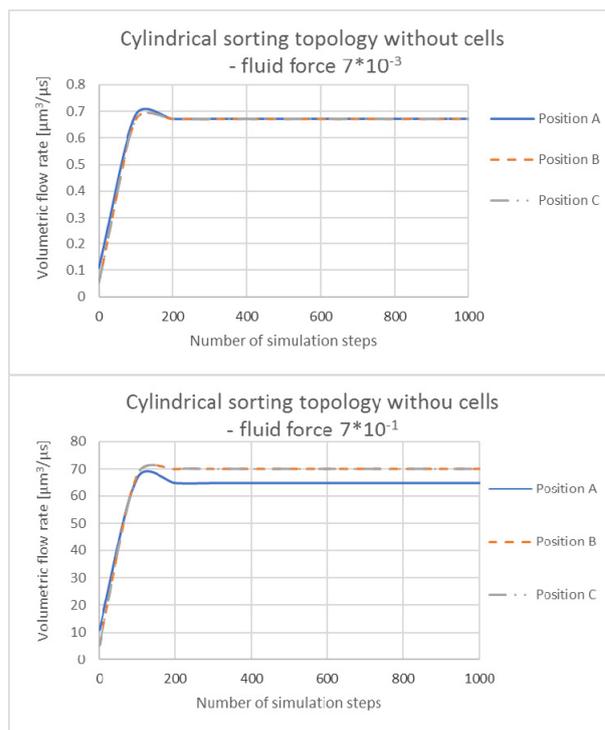


Fig. 7. Cylindrical sorting topology without cells: comparison between the evolution of the volumetric flow rate measured in three different positions A, B and C. The presented graphs came from the simulation with the smallest and with the highest parameter of external force.

In the case of spherical adhesive topology, similarly, the volumetric flow rate was measured in three different sections, presented in the Fig. 8. The volumetric flow rates measured in three different positions are presented in Fig. 9. The curves were obtained from simulations with the multiplier 1 and 500. For the multiplier 500, the difference between steady values is 2%.

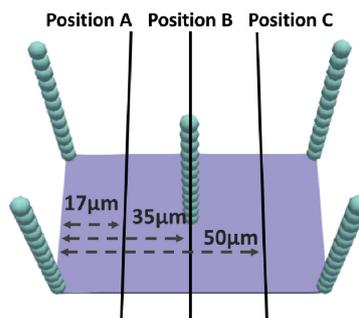


Fig. 8. Three different position of measurement of the volumetric flow rate in the channel with spherical adhesive topology

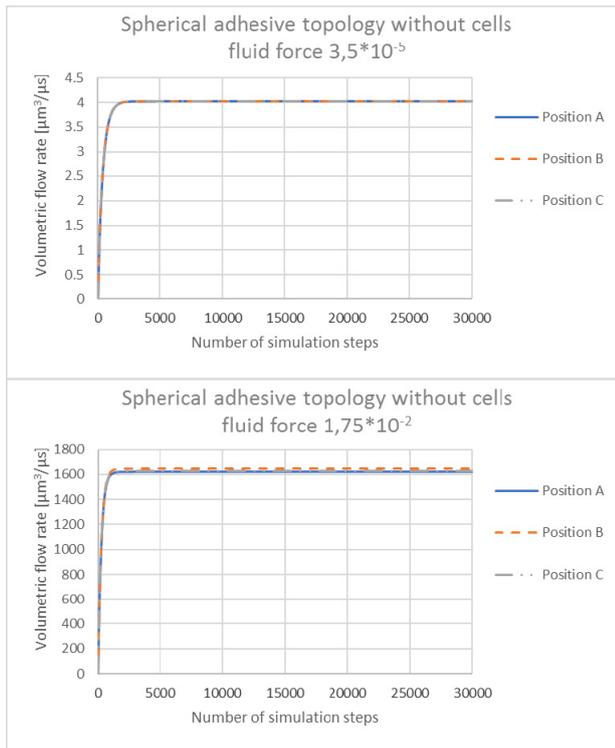


Fig. 9. Spherical adhesive topology without cells: comparison between the evolution of the volumetric flow rate measured in three different positions A, B and C. The presented graphs came from the simulation with the smallest and with the highest parameter of external force.

4 Results

4.1. Cylindrical adhesive topology

The initial value of the external force parameter for the simulation without RBCs was set to 0.000007, which results in volumetric flow rate of $4,48 \mu\text{m}^3/\mu\text{s}$. The value of external force parameter was multiplied afterwards. The results of simulations run with different multipliers are presented in **Fig. 10** and **Fig. 11**.

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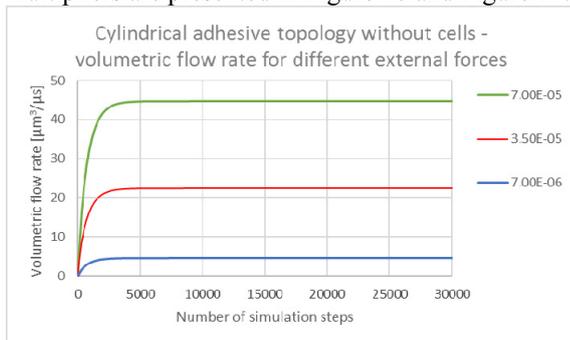


Fig. 10. Example of evolution of the volumetric flow rate during the simulation for three different values of external force for simulations without cells.

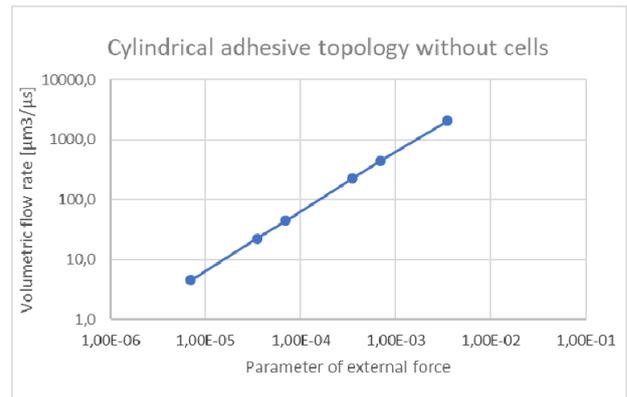


Fig. 11. Relationship between the value of external force parameter and the volumetric flow rate for simulations without cells

For the simulation with cells, we used a hematocrit of approximately 0,5 %. Three simulations were run with this level of hematocrit, with three different values of the external force parameter. The simulation with multiplier of 100 did not reached the end, it finished prematurely because of instabilities in the simulation. The obtained values of volumetric flow rates are presented in the **Fig. 12**.

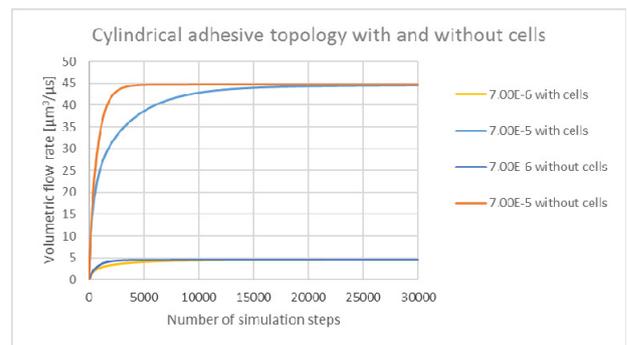


Fig. 12. Comparison between the evolution of the volumetric flow rate for simulations with and without cells.

As we can see in the **Table 1**, the volumetric flow rate is linearly dependent on the parameter of the external force for the small multipliers. The linearity is more significantly disturbed when the multiplier reaches value of 500. The percental difference between the obtained and the expected value is about 7.5%.

Another observation we can make is the similarity between values of volumetric flow rates obtained from simulations without RBCs, and the values obtained from simulations with the RBCs. The models with cells get obviously more time to get a stabilized volumetric flow rate, but the final value is similar to the one without cells. More detailed analysis of the similarity between flow rates obtained from simulations with and without cells is presented in the end of this chapter.

Table 1. Values of volumetric flow rate appropriated to different values of external force parameter.

		without cells		
external force	multiplier applied to the external force	volumetric flow rate [$\mu\text{m}^3/\mu\text{s}$]	obtained multiplier for volumetric flow rate	percentual difference between the two multipliers
7,00E-06	1	4,5	1,0	0,0%
3,50E-05	5	22,4	5,0	0,0%
7,00E-05	10	44,8	10,0	0,0%
3,50E-04	50	223,6	49,9	-0,1%
7,00E-04	100	445,7	99,6	-0,4%
3,50E-03	500	2071,4	462,7	-7,5%
		with cells		
external force	multiplier applied to the external force	volumetric flow rate [$\mu\text{m}^3/\mu\text{s}$]	obtained multiplier for volumetric flow rate	percentual difference between the two multipliers
7,00E-06	1	4,4	1,0	0,0%
1,40E-05	2	8,9	2,0	0,0%
7,00E-05	10	44,5	10,0	0,3%

4.2 Cylindrical sorting topology

The value of external force was set initially to 0,007, and the obtained volumetric flow rate is 0,67 $\mu\text{m}^3/\mu\text{s}$. The results for different values of the external force parameter settings are presented in the Fig. 13 and Fig. 14. The results from the simulations with cells with hematocrit 20 % are presented in the Fig. 15.

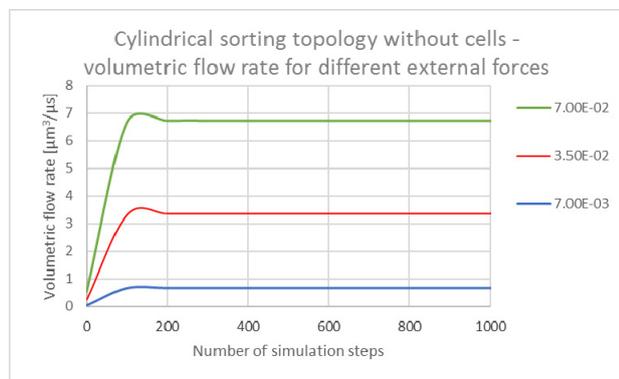


Fig. 13. Example of evolution of the volumetric flow rate during the simulation for three different values of external force for simulations without cells

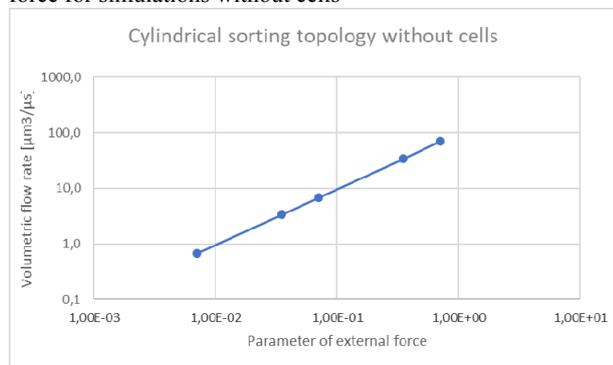


Fig. 14. Relationship between the value of external force parameter and the volumetric flow rate for simulations without cells

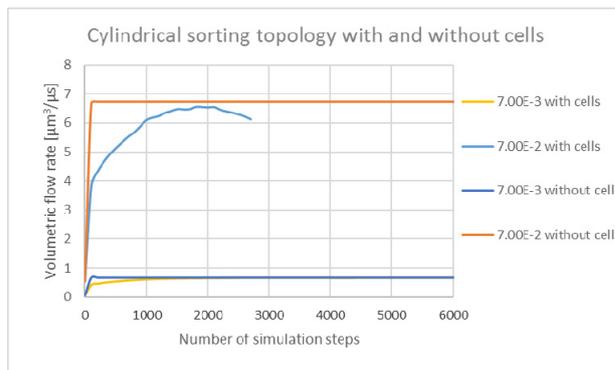


Fig. 15. Comparison between the evolution of the volumetric flow rate for simulations with and without cells.

The observation we make is the very same as for the previous topologies – the volumetric flow rate is linearly dependent on the parameter of external force for the small multipliers (Table 2). The simulations without cells and with a multiplier of 500 finished earlier than expected, due to the numerical instability. However, the expected value of the volumetric flow rate in this simulation corresponds to 1000 $\mu\text{l}/\text{min}$, a value 500 times bigger than the value of the real volumetric flow rate in the laboratory experiment. This means that the linearity is preserved for the common range of the volumetric flow rate used in laboratory experiments.

Table 2. Values of volumetric flow rate appropriated to different values of external force parameter. The linearity is slightly perturbed for the big multipliers.

		without cells		
external force	multiplier applied to the external force	volumetric flow rate [$\mu\text{m}^3/\mu\text{s}$]	obtained multiplier for volumetric flow rate	percentual difference between the two multipliers
7,00E-03	1	0,7	1,0	0,0%
3,50E-02	5	3,4	5,0	0,1%
7,00E-02	10	6,7	10,0	0,2%
3,50E-01	50	34,1	50,7	1,4%
7,00E-01	100	70,1	104,3	4,3%
		with cells		
external force	multiplier applied to the external force	volumetric flow rate [$\mu\text{m}^3/\mu\text{s}$]	obtained multiplier for volumetric flow rate	percentual difference between the two multipliers
7,00E-03	1	0,6	1	0%
1,40E-02	2	1,3	2,1	3%

4.3 Spherical adhesive topology

The initial value of the external force parameter was set to 0.000035, which results in volumetric flow rate of 4.02 $\mu\text{m}^3/\mu\text{s}$. The results obtained by modifying this parameter are presented in the Fig. 16 and Fig. 17 for without cells simulations, and in the Fig. 18 for simulations with cells. The hematocrit used for simulations with cells is 15 %.

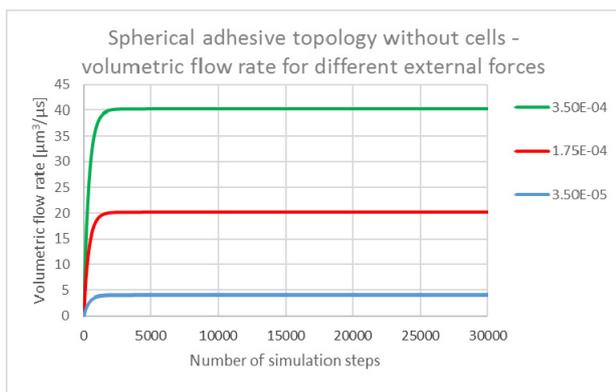


Fig. 16. Example of evolution of the volumetric flow rate during the simulation for three different values of external force for simulations without cells

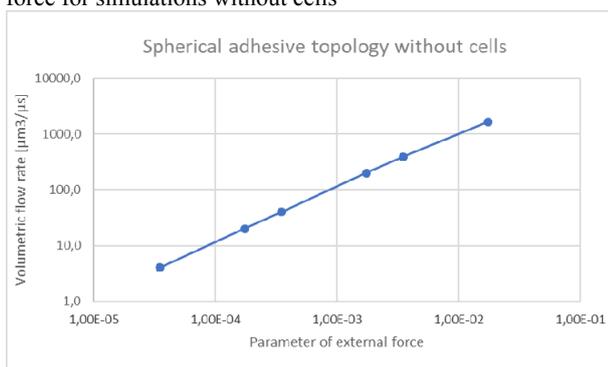


Fig. 17. Relationship between the value of external force parameter and the volumetric flow rate for simulations without cells.

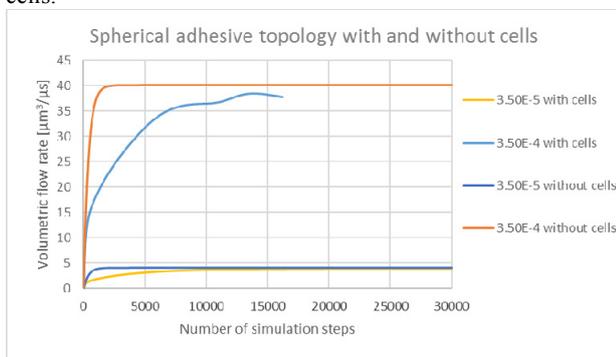


Fig. 18. Comparison between the evolution of the volumetric flow rate for simulations with and without cells.

Table 3. Values of volumetric flow rate appropriated to different values of external force parameter. The linearity is slightly perturbed for the big multipliers.

		without cells		
external force	multiplier applied to the external force	volumetric flow rate [µm³/µs]	obtained multiplier for volumetric flow rate	percentual difference between the two multipliers
3,50E-05	1	4,0	1,0	0,0%
1,75E-04	5	20,1	5,0	0,0%
3,50E-04	10	40,2	10,0	0,0%
1,75E-03	50	199,6	49,6	-0,7%
3,50E-03	100	391,8	97,4	-2,6%
1,75E-02	500	1644,6	409,1	-18,2%
		with cells		
external force	multiplier applied to the external force	volumetric flow rate [µm³/µs]	obtained multiplier for volumetric flow rate	percentual difference between the two multipliers
3,50E-05	1	4,0	1,0	0,0%
7,00E-05	2	7,83	2,0	-1,1%

The linear relationship between the volumetric flow rate and the external force parameter is partially verified (**Table 3**). The linearity is valid for multipliers smaller than 100. For the multiplier of 500, the difference between the obtained and the expected volumetric flow rate is about 18%.

The value of external force was set initially to 0,007, and the obtained

4.4 Comparison of volumetric flow rate for simulations with and without cells

The values are comparable for simulations with and without the cells, especially for simulations with a low hematocrit. The simulation with a high hematocrit shows that it influences the volumetric flow rate in a significant way, even for values close to the laboratory conditions. The comparison between the volumetric flow rates obtained from simulations with and without the cells are presented in the **Table 4**.

Table 4. Volumetric flow rate variation due to presence of the cells in the simulation.

Type of topology	External force	Volumetric flow rate [µm³/µs]		Percentual difference for volumetric flow rate with and without cells
		without cells	with cells	
Cylindrical adhesive topology	7,00E-06	4,48	4,44	-1%
	7,00E-05	44,76	44,52	-1%
Cylindrical sorting topology	7,00E-03	0,67	0,63	-6%
Spherical adhesive topology	3,50E-05	4,02	3,96	-2%

4 Conclusions

In our study, we have tested the linear dependence of the volumetric flow rate on the parameter of the external fluid force. For this purpose, we have tested the hypothesis on three different topologies inspired by existing laboratory experiments. The results show that for the range of values commonly used in the laboratory, the hypothesis is valid. For volumetric flow rates which exceed the laboratory conditions by a factor of hundreds, the linearity is not certain.

Another observation is that for high values of the external force parameter, simulations are less stable, particularly the simulations with cells.

An interesting observation can be done about the time necessary for the simulation to get a stabilized volumetric flow rate. For simulations without cells, this time is 2 to 3 times shorter than for the simulations with the same topology and with cells inside the liquid.

References

- [1] I. Cimrák, M. Gusenbauer, T. Schrefl, Computers and Mathematics with Applications, **64(3)**, 278-288 (2012)
- [2] I. Cimrák, M. Gusenbauer, I. Jančigová, Computer Physics Communications, **185**, 900-907 (2014)

- [3] S. Nagrath, L.V. Sequist et al., *Nature*, **450**, 1235–1239 (2007)
- [4] J.P. Gleghorn et al., *Lab Chip*, **10**, 27–29 (2010)
- [5] S.H. Holm, J.P. Beech, M.P. Barret, J.O. Tegenfeldt, *Anal. Methods*, **8**, 3291-3300 (2016)
- [6] D. Horák, *Biomed Mater Res Part A*, **101A**, 23–32 (2013)
- [7] S. Zheng, H. Lin, J. Liu, M. Balic, R. Datar, R. Cote, Y. Tai, *Journal of Chromatography A*, **1162**, 154–161 (2007)
- [8] A. E. Saliba et al., *PNAS*, **107**, 14524–14529 (2010)