Determination of the applicability limits of the ideal gas model for the calculation of moist air properties

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Abstract. The submitted paper deals with the finding of such moist air states in which the components of moist air and hence the humid air itself can be described by the ideal gas model while maintaining a predefined accuracy. Both components of moist air (dry air and water vapor) can be described by a model of ideal gas at sufficiently low pressures and sufficiently high temperatures. In the paper, we are looking for such combinations of pressures and temperatures for both components, where the relative deviation in the density calculation using the ideal gas model does not exceed the desired value. In addition, on the basis of the mixture theory, such moist air conditions (characterized by pressure, temperature and specific humidity) are searched, on which the accuracy of the calculation meets the required conditions. Subsequently, diagrams are constructed that can be used to help identify the interface between a moist air area that can be described by a simple ideal gas model, and areas where it is necessary to use a more accurate model for one of the components.

1 Introduction

Moist air is a mixture of dry air and water. To simplify the physical description of moist air that occurs in nature or participates in technical processes, only these components can be considered. The presence of other substances in the air (such as various gaseous pollutants, liquid components of different technologies, solids dust) will not be considered.

If all the humidity in the moist air is in gaseous state, the moist air is a homogeneous mixture. The moist air as a heterogeneous mixture contains water partly in the gas state and partly in the liquid state or partly in the gas state and partly in the solid state or partly in the gas state, partly in the liquid state and partly in the solid state.

Humidity is present in relatively small quantities in atmospheric air but is of considerable importance both in the environment and in technological processes too.

In the case of homogeneous moist air, the properties of its components are very often described by the ideal gas model. Therefore, the aim of the present paper is to determine the applicability range of this model when the ideal gas is defined [1] as a theoretical model of compressible fluid, which is characterized by the relationship between the quantities of pressure \( p \) [\( \text{Pa} \)], temperature \( T \) [\( \text{K} \)] and density \( \rho \) [\( \text{kg} \cdot \text{m}^{-3} \)]

\[
p = \rho \cdot r \cdot T ,
\]

where \( r \) [\( \text{J} \cdot \text{kg}^{-1} \cdot \text{K} \)] is the specific gas constant, and by constant values of specific heat capacities at constant pressure \( c_r \) [\( \text{J} \cdot \text{kg}^{-1} \cdot \text{K} \)] = const. and at constant volume \( c_v \) [\( \text{J} \cdot \text{kg}^{-1} \cdot \text{K} \)] = const.

2 Components of moist air

Dry air

The basic physical and thermodynamic properties of dry air as an ideal gas are, according to NIST (National Institute for Standards and Technology, USA) [2]:

- middle molar mass: \( M = 28.958 \text{ kg} \cdot \text{kmol}^{-1} \)
- specific gas constant: \( r = 287.12 \text{ J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1} \)
- specific heat capacity at constant pressure: \( c_r = 1005.9 \text{ J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1} \)
- specific heat capacity at constant volume: \( c_v = 718.8 \text{ J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1} \)
- Poisson constant: \( \kappa = 1.399 \)
The theoretical model of ideal gas is very advantageous for its simplicity, and is usually used to calculate the dry air properties near the atmospheric conditions. For conditions significantly different from standard atmospheric, the ideal gas model exhibits large deviations from the actual values. Figure 1 shows the actual temperature-dependence of density for selected constant pressure values (solid lines) and for a comparison of the same dependences obtained with the ideal gas model (dashed lines).

It can be seen that the ideal gas model is useful for calculating the density of dry air at low pressures and high temperatures. With increasing pressure and decreasing temperature, the differences between actual density and density calculated using the ideal gas model increase.

Figure 2 shows relative dry air density deviations calculated using the ideal gas model for the same pressures as were used in Figure 1. The area of possible use of the ideal gas model depends on the required accuracy: if we would consider the accuracy of 10% to be sufficient, it is possible to use the ideal gas model to calculate the dry air density up to pressures of 20 MPa (for temperatures higher than −50 °C) (for pressures of 50 MPa and 100 MPa, the ideal gas model in the entire temperature range of −150 °C to 1000 °C gives deviations greater than 10%). However, if we require an accuracy of up to 1%, it would be possible to use the ideal gas model to calculate the dry air density only within a certain range of pressures and temperatures. For example, for a pressure of 0.1 MPa, the ideal gas model is usable for temperatures above about −140 °C, for a pressure of 1 MPa it is usable for temperatures above about −25 °C, for a pressure of 5 MPa in a temperature range of 25 °C to 105 °C, for a pressure of 10 MPa only in a temperature range of 25 °C to 75 °C.

**Water** is a relatively complex substance from a thermodynamic point of view. It is the most widely used and best available energy transfer medium. Therefore, extensive research has focused on it. For the work of experts, documents are gradually approved by the International Association for Properties of Water and Steam (IAPWS) as recommendations for determination of thermodynamic and thermophysical properties of water and water vapour. Important documents are:

- 1995 (revision 2014) – Formulation of the thermodynamic properties of a common water substance for general and scientific use (in short IAPWS-95),

The state equations of water and water vapour according to these two documents are substantially distant from the state equation of the ideal gas.

The basic physical properties of water and water vapour are according to [3]:

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**Fig. 1.** Dependence of dry air density on pressures of 0.1 MPa (black), 1 MPa (yellow), 5 MPa (blue), 10 MPa (purple), 20 MPa (green), 50 MPa (brown) and 100 MPa (red) according to [2] (solid lines) and dependence of dry air density obtained with the ideal gas model (dashed lines).

**Fig. 2.** The course of relative deviations of dry air densities as an ideal gas in dependence on temperature for pressures of 0.1 MPa (black), 1 MPa (yellow), 5 MPa (blue), 10 MPa (purple), 20 MPa (green), 50 MPa (brown) and 100 MPa (red).
molar mass: \[ M = 18.015 \, \text{kg/kmol} \]

specific gas constant: \[ r = 461.526 \, \text{J/kg} \cdot \text{K} \]

triple point:
- temperature of triple point: \[ T_c = 0.01 \, ^\circ\text{C} = 273.16 \, \text{K} \]
- pressure of triple point: \[ p_c = 611.637 \, \text{Pa} \]

critical point:
- critical temperature: \[ T_m = 373.946 \, ^\circ\text{C} = 647.096 \, \text{K} \]
- critical pressure: \[ p_m = 22.064 \, \text{MPa} \]

The thermodynamic behaviour of water occurring in homogeneous moist air (i.e. water vapour) can be described under normal atmospheric conditions by the ideal gas model.

The basic thermodynamic properties of water vapour as an ideal gas:
- specific heat capacity at constant pressure:
  \[ c_v = 1.898 \, \text{J/kg} \cdot \text{K} \]
- specific heat capacity at constant volume:
  \[ c_p = 1.398 \, \text{J/kg} \cdot \text{K} \]
- Poisson constant:
  \[ \kappa = 1.33 \]

For thermodynamics of humid air, it is important to determine the thermodynamic properties of water vapour within the limits of saturation.

On the liquid-steam saturation limit from the triple point to the critical point, the scientific formulation (IAPWS 95) \([4]\) determines the dependence of the saturated water vapour pressure on the temperature function

\[ p^* = p_m \cdot \exp \left[ \frac{1}{1 - \tau} \left( a_0 + a_1 \cdot \tau + a_2 \cdot \tau^2 + \ldots \right) \right], \quad (2) \]

where:

\[ \tau = 1 - \frac{T}{T_m} \quad (3) \]

is the temperature function, \[ \tau [K] \] is the temperature, the constants in equation (2) are given in Table 1.

| \( a_0 \) | \( -7.859 \) | \( 517 \) | \( 83 \) |
| \( a_1 \) | \( 1.844 \) | \( 682 \) | \( 59 \) |
| \( a_2 \) | \( -11.786 \) | \( 649 \) | \( 7 \) |

On the saturation limit between the solid and gaseous phases \( H_2O \), i.e. for sublimation and desublimation, the pressure in relation to the temperature from \( 50 \, \text{K} \) to the triple point is solved according to equation \([5]\)

\[ p^w = p_m \cdot \exp \left( \frac{1}{1 - \tau} \sum a_i \cdot \tau^i \right), \quad (4) \]

where:

\[ \theta = \frac{\tau}{T_m}, \quad (5) \]

is the temperature function, \[ \tau [K] \] is the temperature, the constants in equation (4) are given in Table 2.

| \( a_0 \) | \( -2.124 \) | \( 400 \) | \( 6 \) |
| \( a_1 \) | \( 27.320 \) | \( 381 \) | \( 9 \) |
| \( a_2 \) | \( -6.105 \) | \( 981 \) | \( 3 \) |

Similarly to dry air, it is advantageous to describe the properties of water steam by the ideal gas model. The scope of application of the ideal gas model is limited by the beginning of condensation, it is by achieving a saturation state (a condition where moist air becomes a heterogeneous mixture, it is when water begins excreted from moist air in the form of a liquid or solid phase), and next by the range of pressures and temperatures where the steam is a homogeneous mixture, but the ideal gas model has large deviations from actual data. Figure 3 shows the actual dependence of density of water vapor at the temperature for selected constant pressure values (full line) and, for comparison, the same dependence obtained using the ideal gas model (dashed lines).

Similarly to dry air, at lower pressures, the ideal gas model produces results comparable to real values, even up to the saturation curve. At higher pressures, the densities calculated using the ideal gas model deviate from the actual properties near the saturation curve, and for very high pressures, the ideal gas model does not produce the right results even at high temperatures.

Figure 4 shows relative water vapor density deviations calculated using the ideal gas model for the pressures used in Figure 3. As with dry air, the area of possible use of the ideal gas model depends on the required accuracy: if we would consider the accuracy of \( 10\% \) to be sufficient, it is possible to use the ideal gas model to calculate the water vapor density up to pressure of \( 2 \, \text{MPa} \) and it up to the saturation curve, or even up to pressure of \( 10 \, \text{MPa} \) for temperatures over \( 500 \, ^\circ\text{C} \). However, if we require an accuracy of up to \( 1\% \), it would be possible to use the ideal gas model to calculate the water vapor density only up to pressure about \( 50 \, \text{kPa} \) (an accuracy of up to \( 1\% \) would be met up to the saturation vaporous state), or to pressure about \( 0.1 \, \text{MPa} \), but only for temperatures higher than about \( 135 \, ^\circ\text{C} \), or to pressure about \( 0.2 \, \text{MPa} \), but only for temperatures higher than about \( 205 \, ^\circ\text{C} \); at a pressure of \( 1 \, \text{MPa} \) the temperature would have to be above \( 450 \, ^\circ\text{C} \).
3 Moist air

Moist air as a thermodynamic system is a mixture of dry air and water substances. The water contained in the moist air is always in the form of steam, but in addition (in the case of a heterogeneous mixture), part of the water in moist air may also be in a liquid or solid state.

Temperature of moist air \( t_{MA} \) [°C] is equal to the temperature of its components

\[ t_{MA} = t_{MA} = t_v = t . \]  

(6)

Pressure of moist air \( p_{MA} \) [Pa] is equal according to Dalton’s law to the sum of the partial pressures of dry air and water vapor

\[ p_{MA} = p_{MA} + p_v . \]  

(7)

Mass of moist air \( m_{MA} \) [kg] is equal

\[ m_{MA} = m_{MA} + m_{v,MA} . \]  

(8)

Density of moist air \( \rho_{MA} \) [kg m\(^{-3}\)] is the ratio of the mass of moist air to its volume, it is provided that the water in the humid air is only as vapor

\[ \rho_{MA} = \rho_{MA} + \rho_v . \]  

(9)

Pressure of saturated water vapor \( p_v^* \) [Pa] is a significant variable for moist air, which determines the water content limit in the air for the existence of a homogeneous mixture:

- if \( p_v < p_v^* \), the moist air is unsaturated and the mixture homogeneous;
- if \( p_v = p_v^* \), the moist air is saturated and the mixture homogeneous;
- if \( p_v > p_v^* \), the moist air is supersaturated and the mixture is heterogeneous.

Pressure of saturated water vapor is a function of temperature

\[ p_v^* = p_v^*(t) , \]  

(10)

(see equation (2), (4)).

Dew point temperature \( t_{DP} \) [°C] is the moist air temperature at which the moist air reaches the saturation state under isobaric cooling

\[ p_v(t_{DP}) = p_v^*(t_{DP}) , \]  

(11)

it is the water vapor contained therein is in the state of saturated vapor.

Fig. 3. The dependence of the density of water vapor at the temperature for pressures \( 611.657 \) Pa (light blue), \( 50 \) kPa (yellow), \( 0.1 \) MPa (purple), \( 0.2 \) MPa (blue), \( 0.5 \) MPa (green), \( 1 \) MPa (red) and \( 2 \) MPa (brown) according to [3] (solid lines) and the dependence of water vapor density obtained with the ideal gas model (dashed lines).

Fig. 4. The course of relative deviations of water vapor densities as an ideal gas in dependence on temperature for pressures \( 611.657 \) Pa (light blue), \( 50 \) kPa (yellow), \( 0.1 \) MPa (purple), \( 0.2 \) MPa (blue), \( 1 \) MPa (red), \( 2 \) MPa (brown), \( 5 \) MPa (green) and \( 10 \) MPa (violet).

Absolute humidity \( a \) [kg m\(^{-3}\)] is the mass amount of steam in unit air volume. It’s the vapor density \( \rho_v \) at vapor pressure \( p_v \) and temperature \( t \). Absolute humidity is defined only for a homogeneous mixture and for given temperature can be within a range \( 0: \alpha = \rho_v^*(t) \).

Relative humidity \( \varphi \) [-] is the ratio of absolute humidity to absolute humidity of saturated humid air at a given temperature

\[ \varphi = \frac{\alpha}{\alpha^*} . \]  

(12)
Relative humidity is defined only for a homogeneous mixture and may be within the range \((0; 1)\). Assuming water vapor contained in moist air can be described by the ideal gas model

\[
\varphi = \frac{p_r}{p^*_r} \quad (13)
\]

Specific air humidity \(x \left[ \text{kg} \cdot \text{kg}^{-1} \right]\) is the ratio of the mass of water to the dry air mass

\[
x = \frac{m_w}{m_{_{\text{ar}}}}. \quad (14)
\]

In the case of a homogeneous mixture in which water is contained in the air only in the vapor form

\[
x = \frac{m_r}{m_{_{\text{ar}}}}. \quad (15)
\]

Specific gas constant of moist air \(r_{_{\text{sw}}} \left[ \text{J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1} \right]\) is equal

\[
r_{_{\text{sw}}} = \frac{1}{1 + x} r_{_{\text{sw}}} + \frac{x}{1 + x} r_r. \quad (16)
\]

and specific heat capacity at constant pressure of moist air \(c_{_{r_{_{\text{sw}}}}} \left[ \text{J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1} \right]\) is equal

\[
c_{_{r_{_{\text{sw}}}}} = \frac{1}{1 + x} c_{_{r_{_{\text{sw}}}}} + \frac{x}{1 + x} c_r. \quad (17)
\]

So that thermodynamic constant of moist air as and ideal gas depend on specific humidity \(x\) ; at \(x = \text{const}\) they are constant.

If the moist air forms a homogeneous mixture and its components (dry air and steam) can be described by the ideal gas model, it is possible to derive the relationship between specific and relative humidity

\[
x = 0.622 \cdot \frac{p_r}{p_{_{\text{sw}}} - p_r}, \quad (18)
\]

resp.

\[
x = 0.622 \cdot \frac{\varphi p^*_r}{p_{_{\text{sw}}} - \varphi p^*_r}. \quad (19)
\]

For the specific humidity of the saturated homogeneous mixture \((\varphi = 1)\)

\[
x^* = 0.622 \cdot \frac{p^*_{r_{(t)}}}{p_{_{\text{sw}}} - p^*_{r_{(t)}}} = x^*_{(t)}. \quad (20)
\]

Our task is to find a boundary in moist air (combination of pressure, temperature and humidity), where there is possible to use the ideal gas model to calculate the properties of components of the moist air with sufficient accuracy.

One limit is the boundary on which ends the existence of water in the form of steam only. The paper [6] deals by searching for this boundary: For a given value of specific humidity of moist air, the dependence of saturated pressure of moist air on the temperature was solved from equations (20) and (2), respectively at the case of sublimation limit may be (4) In the equilibrium \(p-T\) diagram, a curve (with a parameter of the specific humidity \(x\)) was constructed, on which water would be excreted from moist air (the curves of the saturated moist air for different water contents, it is for different values of specific humidity \(x\)), see Figure 5.

![Fig. 5. Curves of saturated moist air for different values of specific humidity \(x\) [6].](image-url)
required accuracy; between the two curves is the area where the moist air is still homogeneous, but for calculating the vapor properties it is necessary to use one of the more precise models; to the left of the saturated moist air curves, the humid air is already heterogeneous.

\[ P_{ma} = P_{ma} \left( 1 + \frac{x}{0.622} \right), \]  

(21)

from which we can calculate the maximum moist air pressure for the selected specific humidity \( x \) in which the density of the dry air contained in the moist air at a given temperature and specific humidity can be calculated using the ideal gas model to meet the required accuracy.

In Figure 6, for the selected values of specific humidity \( x \) is plotted line (dotted curve) of maximum pressure of moist air at which it is possible to calculate the partial dry air density using the ideal gas model, while meeting the required accuracy. It can be seen from the Figure 6 that the upper limit of pressure of moist air depends very little on the specific humidity \( x \) and the curves almost coincide with the dry air partial pressure curve (for the maximum humidity used by us: \( x = 0.1 \), the total pressure is equal to 1.16 times the dry air partial pressure; for humidity \( x = 0.01 \) is equal to 1.016 times and for humidity \( x = 0.001 \) it's just 1.0016 times). Thus, in the diagram in Figure 6, the dependence of the maximum dry air partial pressure on the temperature to fulfill the required accuracy is plotted. This curve is only a little different from the curves describing the total pressure of moist air for the selected values of specific humidity \( x \). In the other figures, therefore, only the dependence of dry air partial pressure is plotted for expression of boundary of usability of model of ideal gas for dry air.

Thus, if the state of moist air is in the \( p-T \) diagram shown by a point below the maximum dry air partial pressure curve, the partial density of the dry air (contained in the moist air) may be calculated with sufficient precision using the ideal gas model.

The area where the densities of the both components of moist air can be calculated using the equation state of the ideal gas lies in the \( p-T \) diagrams "bottom right": it is bounded by a saturated moist air curve for the corresponding specific humidity \( x \), by a curve showing the maximum value of the air pressure (at given temperature and specific humidity \( x \)), when it is possible to model water vapor with the ideal gas model, and by a curve showing the maximum pressure value (at
given temperature) when it is possible to model dry air with the ideal gas model.

![Diagram](https://example.com/diagram.png)

**Fig. 8.** Area of usability of the ideal gas model for the calculation of properties of components of moist air at the required accuracy of the calculation to 5% for specific humidity of 0.1, 0.01 and 0.001.

Figure 6 shows an area in which the ideal gas model achieves an accuracy of 1%, in Figure 7 an area with accuracy of 2%, and in Figure 8 an area with accuracy of 5%.

In all of the figures, an isentropy network (a set of brown curves) is plotted for better idea:

1) If the reversible adiabatic compression occurs in homogeneous moist air, the moist air will "move away" from the saturation curve, it means: it is not possible to get from homogeneous moist air to heterogeneous moist air during isentropic compression, on the other hand, in a reversible adiabatic expansion, it is possible to reach a state of saturated moist air.

2) If the reversible adiabatic compression occurs in homogeneous moist air, the accuracy of the description of the water vapor properties by the ideal gas model will be increased ("we move away" from boundaries limiting the usability of the ideal gas model for steam water), but the accuracy of the calculation of the dry air properties by the ideal gas model will be reduced; in reversible adiabatic compression it is possible to cross the boundary where the ideal gas model is still usable for calculating the dry air properties.

3) In reversible adiabatic expansion, the accuracy of the ideal gas model for the dry air contained in moist air increases; the accuracy of the ideal gas model for the water vapor contained in moist air reduces, but not too significantly - the curve limiting the pressure of moist air for using the ideal gas model for the water vapor has only a slightly greater slope than the isentropic curves; The main "danger" in reversible adiabatic expansion in moist air is the possibility of reaching the condensation boundary.

**Conclusion**

In the equilibrium $p - T$ diagram, the boundary of range of usability of the ideal gas model for the calculation of properties of components of moist air were constructed for different relative accuracy values (1%, 2%, 5%). Once the boundary is reached, a more accurate model is required to use (for calculation of properties of the components of moist air) for preserve the accuracy of calculation.

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