

Terahertz generation from single and multiple filaments in air

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Terahertz (THz) photonics is entering an era of strong electromagnetic fields in the THz frequency range [1]. Optical breakdown plasma takes a specific place among THz sources. With the appearance of compact commercial terawatt laser systems and the actualization of remote sensing tasks, two-color optical breakdown in air has turned into a powerful mechanism of broadband THz generation at the observation point [2]. There are several problems associated with THz remote sensing applications. First, the directional diagram of emitted THz radiation usually has a conical structure [3]. Nevertheless, as the filament length increases, the divergence angle decreases [4]. Second, the limitation of air-breakdown THz applications to remote-sensing tasks is due to low signal-to-noise ratio in the amplitude analysis of THz radiation reflected from the object under study. Therefore, one can use polarization analysis of useful signals whose effective application requires thorough control of the state of polarization of THz radiation. Third, further optimization of the THz yield from two-color filaments is required to approach the signal level obtained from optical rectification in crystals [1]. Control over the relative temporal (and spatial) properties of the optical pump fields, e.g. via the pulse chirp, offers an important means to control and optimize the optical-to-THz conversion process. With increasing pulse peak power and transition to multiple filamentation more overall THz energy is generated however the proper regularization of the high-peak-power beam is needed to provide for the maximum THz yield in the narrow angle.

In this paper we present a series of studies for the control and optimization of THz directionality, polarization and energy yield.

We observe linear-to-elliptical transformation of THz radiation through independent control of initially linear polarization directions of ω and 2ω beams. Strongly elliptical THz radiation appears abruptly at the threshold angle of $\sim 85^\circ$ between the ω and 2ω light fields. The physical reason for THz beam ellipticity was identified by us as the nonlinear response of the birefringent medium induced by an 800 nm pump. We found THz ellipticity and energy yield using a vectorial nonparaxial propagation equation (Fig. 1).

We have demonstrated and justified the mechanism which dictates the optimum chirp to maximize the THz yield from a two-color filament. We have experimentally measured the THz energy yield using 800-nm pulses in a single filament regime, being chirped from 35 fs up to 250 fs both positively and

negatively. We found that the positively chirped pulse proves to be a more advantageous choice since the ω - 2ω pulse overlap is larger by ~ 15 fs as compared with the transform-limited pulse (Fig. 2).

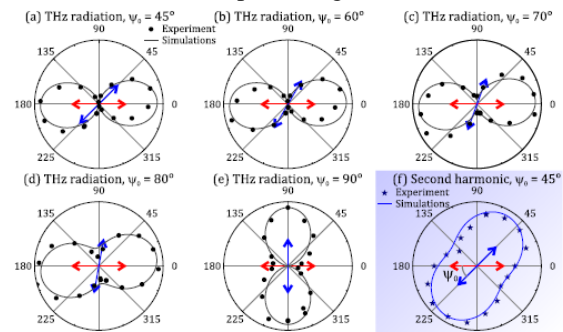


Fig. 1 (a)–(e) THz energy after transmission through THz analyzer, in dependence on its rotation angle, for different angles ψ_0 , for different angles ψ_0 between polarization directions of the 800 nm (ω) and 400 nm (2ω) beams. (f) Second-harmonic energy after its transmission through polarization analyzer. Red (blue) arrow indicates the incident polarization of the ω (2ω) beam.

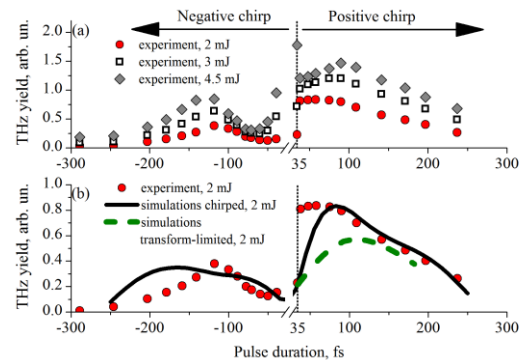


Fig. 2 (a) The experimentally measured dependence of the THz energy on the pump (800 nm) chirped pulse duration. The pump energy is fixed for each chirped pulse duration variation in the range from -300 fs (FWHM, negatively chirped) to $+250$ fs (FWHM, positively chirped). (b) The experimental vs. simulated dependence of the THz energy on the 800-nm 2-mJ chirped pulse duration. Green dashed curve shows simulated THz yield for the transform-limited pulse.

We have shown that a $2D$ array of multiple plasma filaments in atmospheric density gases leads to the efficient increase of terahertz radiation emitted in the forward direction. As the number of filaments in an $N \times N$ array is augmented, the THz energy increase per unit polar angle in the central cone exceeds the standard gain factor N^2 [4,5]. The two-dimensional multiple filament array can be produced from an actu-

al non-ideal elliptical beam using a phase plate to introduce a phase barrier and block the interaction between the neighboring filaments. Such regularized structure of plasma channels is used to control the directional diagram of THz radiation from a sub-terawatt-peak-power 800 nm + 400 nm femtosecond laser pulse.

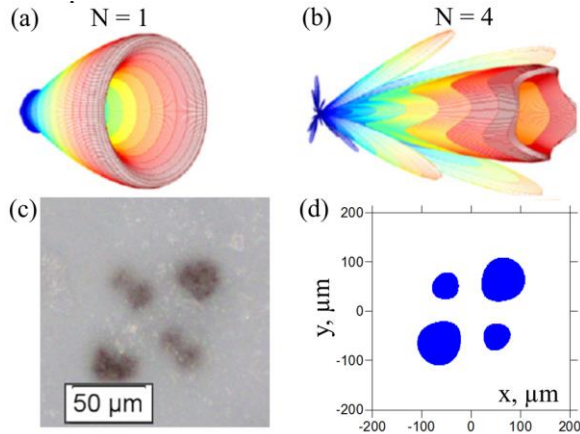


Fig. 3. (a, b) Full 3D angular THz intensity distribution from an array of N multiple filaments with the length $L = 1$ cm. (c) Single-laser-shot burnt photo paper images after the geometrical focus of 1.1-mJ 35-fs 800 nm + 400 nm elliptical beam transported through a phase mask in air. The reproduction of the plasma channels in 3D+time carrier resolved simulations is given in (d).

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