

# Multistability and high reflectance of a mono-layer of three-level quantum emitters with a doublet in the excited state

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**Abstract.** We study theoretically the nonlinear optical response of a mono-layer of three-level quantum emitters with a doublet in the excited state. It is shown that the layer's response exhibits multistability. In a certain frequency range, the monolayer operates as a perfect bistable mirror.

We conduct a theoretical study of the steady-state optical response of a monolayer of regularly spaced three-level quantum emitters (QEs) with a doublet in the excited state. The total (retarded) dipole-dipole interaction of QEs is taken into account. This interaction provides a positive feedback. The interplay of the latter and the immanent nonlinearity of QE's gives rise to a multistability of the monolayer optical response. In a certain frequency range, the system operates as a nanometric bistable mirror.

It is assumed that the monolayer undergoes an action of a CW external field of a Rabi amplitude  $\Omega_0$  and frequency  $\omega_0$ , which is quasi-resonant with the QE's allowed transitions. A constituent QE is modelled by a three-level V-type quantum system with the ground state  $|1\rangle$ , and a doublet  $|2\rangle$  and  $|3\rangle$  in the excited states. The allowed optical transitions are  $|1\rangle \leftrightarrow |2\rangle$  and  $|1\rangle \leftrightarrow |3\rangle$ . They are characterized by the transition dipole moments  $d_{21}$  and  $d_{31}$ , transition frequencies  $\omega_{21}$  and  $\omega_{31}$ , and spontaneous decay constants  $\gamma_{21}$  and  $\gamma_{31}$ . The doublet is described by the splitting  $\Delta_{32}$  and the relaxation constant  $\gamma_{32}$ .

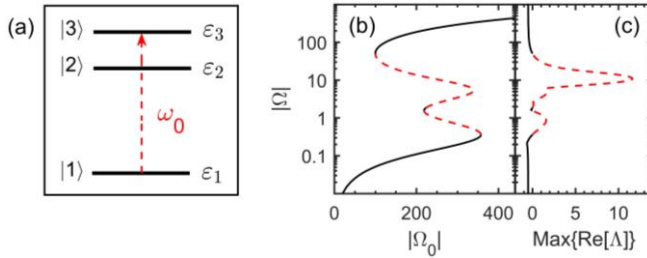
The optical dynamics of a constituent QE is governed by the 3x3 density matrix  $\rho_{\alpha\beta}$  ( $\alpha, \beta = 1, 2, 3$ ). The total field  $\Omega$  acting on a given QE in the monolayer represents a sum of the external field  $\Omega_0$  and the field produced by all others QEs in place of the given one. In this way, the total (retarded) QE-QE dipole-dipole interaction is taken into account. The near-zone (far-zone) part of the QE-QE interaction gives rise to a dynamic renormalization of the transition frequencies  $\omega_{21}$  and  $\omega_{32}$  (relaxation constants  $\gamma_{21}$  and  $\gamma_{31}$ ), depending on the population difference of corresponding transitions [1,2]. The effects are described by the constants  $\Delta_L$  (shift) and  $\gamma_R$  (relaxation). These parameters govern a positive feedback which is responsible for a sophisticated nonlinear optical properties of the monolayer.

In Fig. 1, we present the results of the steady-state calculations performed for the case when the external field is tuned into the resonance with the transition  $|1\rangle \leftrightarrow |3\rangle$  ( $\Delta_{31} = \omega_{31} -$

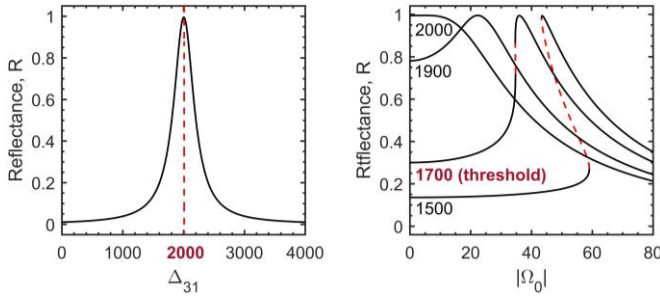
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$\omega_0 = 0$ ). Panel (b) shows the acting field magnitude  $|\Omega|$  as a function of the external field magnitude  $|\Omega_0|$  for the set of parameters typical for two-dimensional supercrystals built up of semiconductor quantum dots (SQD) [3]. As is seen from the plot,  $|\Omega|$  appears to be a multi-valued function of  $|\Omega_0|$ , signaling multistability. The stability of different parts of the steady-state solution has been checked by analyzing the spectrum of Lyapunov exponents  $\Lambda_k$  ( $k = 1, 2, \dots, 8$ ). The maximal real part of  $\{\Lambda_k\}$ ,  $\text{Max}\{\text{Re}[\Lambda]\}$ , is plotted in panel (c).



**Fig. 1.** (a) – Excitation scheme of a QE ( $\Delta_{31} = 0$ ). (b) – Steady-state solution for  $|\Omega|$  as function of  $|\Omega_0|$ . Solid (dashed) parts of the curves indicate stable (unstable) regions of  $|\Omega|$ . (c) – Real part of the major Lyapunov exponent  $\text{Max}\{\text{Re}[\Lambda]\}$  as a function of  $|\Omega|$ . Parameters of calculations are:  $\Delta_{32} = 10$ ,  $\gamma_R = 100$ ,  $\Delta_L = 1000$ ,  $\gamma_{32} = 0.01$ . All quantities is given in units of  $\gamma_{31}$ .



**Fig. 2.** Left panel – the linear reflection coefficient  $R$  as a function of the detuning  $\Delta_{31}$ . Right panel – the field dependence of  $R$  computed for a set of detunings  $\Delta_{31}$  shown in the plot. The value  $\Delta_{31} = 1700$  is the threshold for bistability to occur. The rest of parameters are the same as in Fig. 1.

Fig. 2 shows the detuning and field dependence of the reflectance  $R = |\Omega_{\text{refl}}/\Omega_0|^2$  (left and right panels, respectively),  $\Omega_{\text{refl}} = \gamma_R(\rho_{31} + \rho_{21})$  is the Rabi amplitude of the reflected field. As follows from the left plot, the linear reflectance (for a weak  $|\Omega_0|$ , left panel) has a maximum at  $\Delta_{31} = 2000\gamma_{31}$ . Moreover, at this point  $R$  is approaches unity, i.e. the monolayer almost totally reflects the input field. The right panel in Fig. 2 demonstrates arising three solutions for  $R$  at a given  $|\Omega_0|$ , which means bistability of the reflectance.

Summarizing, we believe that a monolayer comprising V-type QEs may serve as a nanometric bistable mirror. These features might be of interest for nanophotonics. Supercrystals built up of SQDs with the degenerate valence band, e.g. CdSe, placed in magnetic field [4], can be considered as candidates for realization of such systems.

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## References

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