

Fission fragments observables measured at the LOHENGRIN spectrometer

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Abstract. Nuclear fission yields are key data for reactor studies, such as spent fuel inventory or decay heat, and for understanding fission process. Despite a significant effort allocated to measure fission yields during the last decades, the recent evaluated libraries still need improvements in particular in the reduction of the uncertainties. Moreover, some discrepancies between these libraries must be explained.

Additional measurements provide complementary information and estimations of experimental correlations, and new kinds of measurements enable to test the models used during the nuclear data evaluation process. A common effort by the CEA, the LPSC and the ILL aims at tackling these issues by providing precise measurements of isotopic and isobaric fission yields with the related variance-covariance matrices. Additionally, the experimental program involves a large range of observables requested by the evaluations, such as kinetic energy dependency of isotopic yields and odd-even effect in order to test the sharing of total excitation energy and the spin generation mechanism. Another example is the complete range of isotopic distribution per mass that allows the determination of the charge polarization, which has to be consistent for complementary masses (pre-neutron emission). For instance, this information is the key observable for the evaluation of isotopic yields. Finally, ionic charge distributions are indirect measurements of nanosecond isomeric ratios as a probe of the nuclear de-excitation path in the (E^*, J, π) representation.

Measurements for thermal neutron induced fission of ^{241}Pu have been carried out at the ILL in Grenoble, using the LOHENGRIN mass spectrometer. Methods, results and comparison to models calculations will be presented corresponding to a status on fission fragments observables reachable with this facility.

1 Introduction

An accurate knowledge of fission data in the actinide region is important for studies of innovative nuclear reactor concepts. Fission yield measurements supply experimental data to put constraints on fission models and improve their predictive power. In the framework of nuclear data evaluation, these models are indeed necessary to increase the consistency and the precision of the libraries. Despite a real effort on fission yields measurements, current evaluated data still need some improvements on different aspects, such as for instance the uncertainties reduction and the estimation of covariance matrices. A special focus on the heavy and symmetry mass regions is important, since it is where the discrepancies between models (or evaluations) and the few experimental data are mainly observed.

A collaboration between the CEA, the LPSC and the Institute Laue Langevin (ILL) is involved in an experimental program using thermal neutrons of the ILL and the LOHENGRIN spectrometer to study the fission process. We developed different methodologies to obtain absolute isobaric and isotopic yields with the estimation of the covariance matrices associated to the measurements. Besides,

the measurement of different observables combined with a comparison with simulation codes such as FIFRELIN [1] enable to get insight data to better understand the fission process. Isomeric ratios can give an indirect information on the fragments spin distribution, and their kinetic energy dependency enlightens on the validity of the models in use. An exhaustive set of isotopic yields per mass enables the charge polarisation estimation, which has to be consistent between complementary masses. A final example of such indirect data measured by our collaboration is the extraction of nanosecond isomeric ratios determined from the ionic charge distributions per isotope.

2 The LOHENGRIN spectrometer

The LOHENGRIN mass spectrometer [2] is a nuclear physics instrument from the ILL research reactor facility which allows to study fragment distributions from thermal neutron induced fission with a very high mass resolution ($\Delta A/A \approx 1/400$). A fissile actinide target is placed close to the reactor core, in a thermal neutron flux reaching 5×10^{14} neutron.cm⁻².s⁻¹.

Fission fragments emerge from the target with an ionic charge distributed around an average ionic charge state of

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about 20 to 23. Those fragments that are emitted along the beam tube axis undergo a horizontal deflection in a magnetic field, directly followed by a vertical deflection in an electric field. These combined fields separate ions according to their A/q and E_k/q ratios, with A , q and E_k the mass, ionic charge state and kinetic energy of the ions respectively. These ratios can be achieved with different triplets (A, E_k, q) leading to a possible degeneracy.

At the spectrometer exit, different detection systems can be installed, such as a dual anode Frisch grid ionisation chamber for mass yield measurements, or two clovers of four high purity Germanium crystals that are used with an additional magnet whose aim is to focus the ion beam. A schematic view of the spectrometer is shown in Fig. 1.

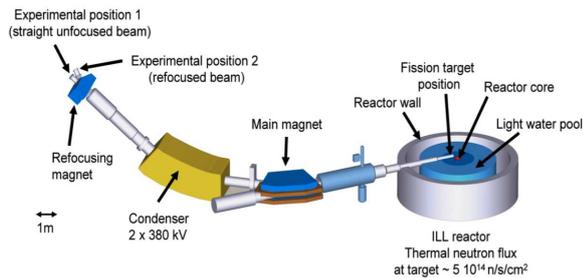


Figure 1. Schematic overview of the LOHENGRIN spectrometer at ILL.

3 Mass and isotopic yields measurements of the $^{241}\text{Pu}(n_{th},f)$ reaction

The first goal of our collaboration measurement campaign consists in the precise measurements of mass, isotopic and isomeric yields, with a control of the systematic effects and the determination of the covariance matrices associated to the analyses. For these observables, their dependency with fission fragment kinetic energy increases significantly the retrieved information on the fission process.

3.1 Mass yields

Isobaric yields are obtained from experimental position 1 (see Fig. 1) after an integration over the kinetic energy and the ionic charge distributions of the count rates measured with the ionisation chamber. A new measurement method and consequent analysis path have been developed and are detailed in Ref. [3–6]. Among the special features of this method are the self-normalisation of our data and the calculation of the experimental covariance matrices. Provided that all the heavy mass rates are measured, it is possible to self-normalise the data by defining to 100% the sum of the whole heavy peak yields. As a consequence, these new measurements are independent from another experiment or assessment and may be compared directly with the existing data and evaluations.

The results for $^{241}\text{Pu}(n_{th},f)$ are shown in Fig. 2, where they are compared to the JEFF-3.3 [7] and ENDF/B-VII.1 [8] libraries. The whole heavy peak and an important part

of the light one were measured. Our results are slightly higher than the libraries for the light mass region, and a structure around mass 140 is observed in the heavy region. Our experimental uncertainties are around 5% on average and below the ones indicated in the two libraries.

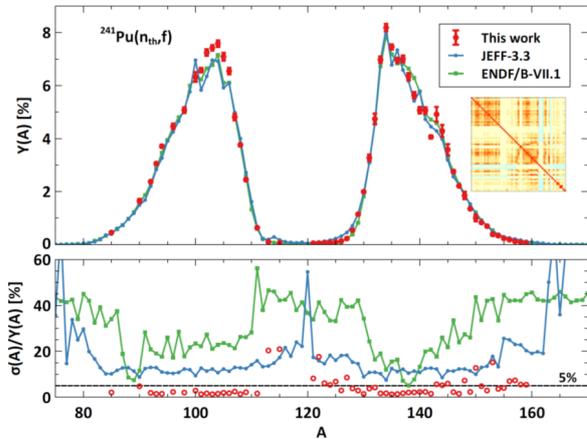


Figure 2. Mass yields for the $^{241}\text{Pu}(n_{th},f)$ reaction and their relative uncertainties.

3.2 Isotopic yields

Isotopic and long-lived isomeric (few μs to ms) fission yields are measured by gamma spectrometry. Experimental position 2 (see Fig. 1) is now used. The ion beam is deposited on a moving tape inside a vacuum chamber and a cumulative measurement with duration of about 30 min per point over the ionic charge distribution is achieved. The tape moves at the end of the measurement to clean the environment and start a new measurement. After corrections of the Bateman equations and the estimation of the contribution of the isotopes from the tape only, we obtain the isotopic distributions per mass. As for the mass yields, a particular effort is made on the determination of the systematic uncertainties and the covariance associated to the measurement process.

Fig. 3 shows a scheme summarizing the measurement procedure and the isotopic yields obtained for the chains 137 and 139 are shown in Fig. 4, along with the covariance matrices. It is important to note that the uncertainties are dominated by the nuclear structure data. Thus current yields measurements can be improved by increasing the nuclear structure knowledge. Fig. 5 helps to understand the construction of the experimental covariance matrix at the main steps of the analysis as illustration of the uncertainty propagation effects.

4 Indirect data measurements

Besides isobaric and isotopic yields, other fission observables are achievable with the LOHENGRIN spectrometer and give important complementary information for the study of the fission process. The kinetic energy dependency of the isotopic and isomeric yields has been already

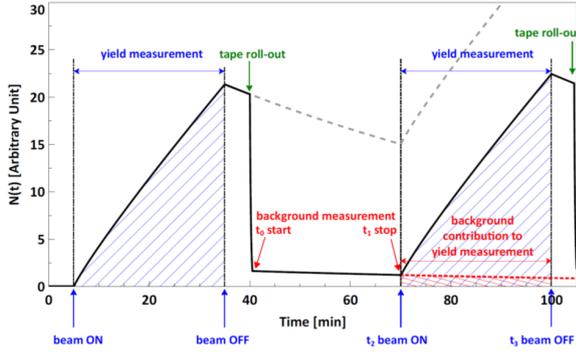


Figure 3. Evolution of a typical isotopic yields measurement procedure.

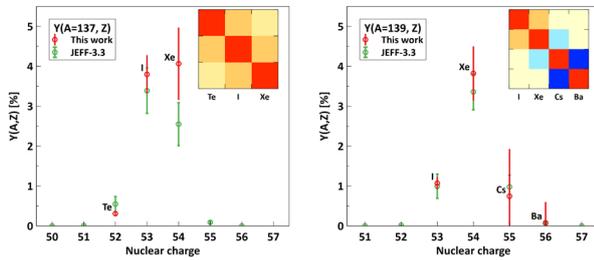


Figure 4. Isotopic yields of the $^{241}\text{Pu}(n_{th}, f)$ reaction for masses 137 and 139, along with the experimental covariance matrices

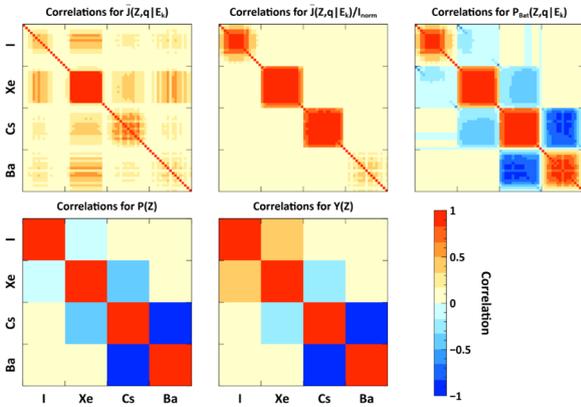


Figure 5. The correlation matrix for mass 139 at different steps of the analysis. 1st step: average on the different gamma rays. 2nd step: Division by I^{norm} . 3rd step: Independent production rate calculation. 4th step: sum over the ionic charges. 5th step: absolute yields after self normalisation.

discussed in ref. [9, 10]. This paper will focus on the description of the charge polarisation and the estimation of nanosecond isomeric ratios.

4.1 Nuclear charge polarisation

The charge polarisation can be extracted from the combination of the isotopic and isobaric yields measurements. It is defined as the difference between the measured mean

nuclear charge and the fragment nuclear charge in the Un-changed Charge Density (UCD) hypothesis.

Fig. 6 shows the measured charge polarisation for the concerned masses in the heavy peak region, compared with the JEFF-3.1.1 library and previous experimental data from Schillebeeckx et al. [11]. We observe a good agreement for the mass 130 and around mass 140, but a strong structure appears for the masses 132, 136 and 138. Complementary measurements on the neighbouring masses are planned by the collaboration to better understand this phenomenon.

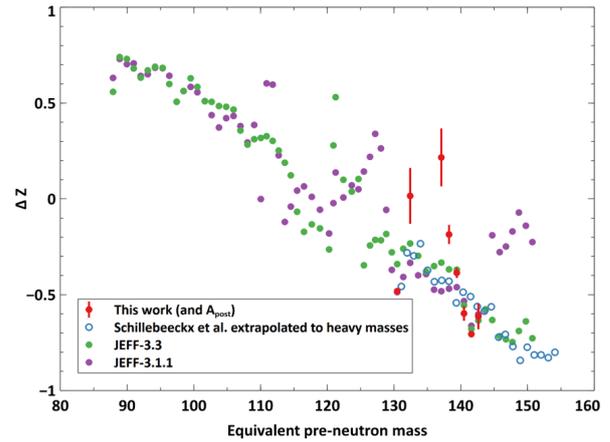


Figure 6. Charge polarisation measured in the heavy peak region as a function of the pre-neutron mass, compared with the JEFF-3.1.1 library and previous experimental data from Schillebeeckx et al. [11].

4.2 Nanosecond converted isomeric ratios

Indirect measurements of nanosecond IR's can be determined from the ionic charge distributions [12, 13]. The method consists in the deconvolution of the ionic charge distribution per isotope obtained by gamma spectrometry after correction from Bateman equations. The converted isomeric ratio (CIR) is defined as the converted isomer population over the total ionic population (converted and unconverted).

$$CIR = \frac{N(A, Z, m \rightarrow e^-)}{N(A, Z, GS) + N(A, Z, m \rightarrow \gamma) + N(A, Z, m \rightarrow e^-)} \quad (1)$$

According to the statistical models from H. Betz [14], we assume that the ionic charge distribution associated to the unconverted population follows a Gaussian distribution due to the charge equilibrium in the cover of the target (a Nickel foil in this work). A deviation from this Gaussian distribution indicates a charge modification due to the conversion from ps and ns isomers to groundstate (see Fig. 7). In most of the cases the deconvolution is achieved using two Gaussian distributions and a Monte Carlo simulation to deduce the CIR in order to consider the covariance terms between the Gaussian integrals. CIR measurements for $^{241}\text{Pu}(n_{th}, f)$ and $^{233}\text{U}(n_{th}, f)$ are shown in Fig. 8 and compared to FIFRELIN calculations for $^{241}\text{Pu}(n_{th}, f)$. We note

