

ANALYSIS OF THE OECD/NEA SFR BENCHMARK WITH ANTS REDUCED-ORDER NODAL DIFFUSION SOLVER AND THE SERPENT MONTE CARLO CODE

Marton Szogradi

VTT Technical Research Centre of Finland Ltd
P.O. Box 1000, FI-02044 VTT, Finland

marton.szogradi@vtt.fi

ABSTRACT

In order to meet modern industrial and scientific demands the Kraken multi-physics platform's development was recently launched at VTT Technical Research Centre of Finland. The neutronic solver of the framework consists of two calculation chains, providing full core solutions by the Serpent high fidelity code (1) and the AFEN/FENM-based reduced-order diffusion solver called Ants (2) capable of handling square and hexagonal geometries in steady-state. Present work introduces the simulation of a large 3600 MWth Sodium-cooled Fast Reactor (SFR) described within the activities of the Working Party on Scientific Issues of Reactor Systems (WPRS) of OECD. Full-core 3D results were obtained by Serpent for carbide- and oxide-fuel cores, moreover group constants were generated for Ants utilizing 2D super-cell and single assembly infinite lattice models of Serpent. The continuous-energy Monte Carlo method provided the reference results for the verification of the reduced-order method. Implementing the spatially homogenized properties, 3D solutions were obtained by Ants as well for both core configurations. Comparison was made between the various core designs and codes based on reactivity feedbacks (Doppler constant, sodium voiding, control rod worth) considering power distributions. Regarding reactivity sensitivity on geometry, axial fuel- and radial core expansion coefficients were evaluated as well.

KEYWORDS: SFR, nodal diffusion solver, Ants, Serpent

1. INTRODUCTION

Benchmark data from report [1] was used to construct 3D full-core SFR models in Serpent with respect to given boundary conditions. Two configurations, one of carbide (CA) and one of oxide (OX) fuel was defined with Beginning of Cycle (BoC) isotopic composition. In Section 2 the core design is discussed and the performance of the models under various perturbations e.g. Doppler effect, sodium voiding and diagrid expansion. In Section 3 few-group constants' generation is described with 2D spatial homogenization highlighting the sensitivity of various condensation methods. In Section 4 the diffusion solver is introduced briefly regarding theoretical background and capabilities. Finally the assessment of Ants will be carried out with respect to 3D Serpent results.

2. SERPENT FULL-CORE 3D MODELS

Both studied cores have medium power density, structural materials are identical (e.g. Oxide Strengthened Steel - ODS cladding) and helium bonded fuel pellets. Each design bears benefits and disadvantages therefore considering both configurations gives a wider perspective regarding code versatility and reliability. The CA core's main motivation was to achieve a low linear fuel temperature gradient providing an increased safety margin for pellet-centerline temperature. The OX configuration was conceived with the "fat pin with small wire" approach i.e. the core could reach self-breeding without fertile blanket. The models represent BoC conditions thus the helium bonded fuel pellets are not in contact with the cladding.

2.1. Core Configurations

The CA core contains 487 fuel assemblies (206 inner / 201 outer), 270 radial reflector elements and 27 control assemblies whereas 18 and 9 assemblies represent the primary and secondary control system, respectively. The (U,Pu)C pellets are whole, the fuel cladding is made of ODS alloy while the shroud tube and the axial reflector elements consist of EM10 SS (Euralloy). The OX core has 453 fuel assemblies (225 inner / 228 outer) and the same radial reflector shield and control systems as the CA. The (U,Pu)O₂ pellets are hollow with ODS cladding and similar EM10 structures as the CA fuel. Major geometry and material properties of the cores are tabulated in Table 1. The core layouts are shown in Fig. 1. Regarding the control rods, both system contains B₄C pellets as absorber material, whereas the primary rods have natural isotopic composition while the secondary assemblies are loaded with enriched (90 % ¹⁰B) boron carbide.

Table 1: Comparison of Core Design.

Property	CA	OX
Overall height (cm)	301.70	311.16
Active core height (cm)	100.56	100.56
Subassembly pitch (cm)	20.9889	21.2205
Number of fuel pins (-)	469	271
Pellet radius (mm)	3.319	4.742
Inner/Outer core Pu content (wt%)	13.76/17.60	15.81/17.60

2.2. Present Models

Full-core 3D model inputs were built for Serpent 2.1.31 [2] based on Ref. [1] whereas He pressure (atomic density) was chosen to be 10 bar according to MOX fuel recommendations reported by the Halden Reactor Project [3]. The fuel pins of the inner and outer core were divided into 5 axial nodes representing different Pu enrichments. The control assemblies were fully withdrawn from the core in every test case except for the control rod worth (CRW) simulations. A total of 1700 cycles (1500 active / 200 inactive cycles) were calculated with 600 000 neutron histories/cycle

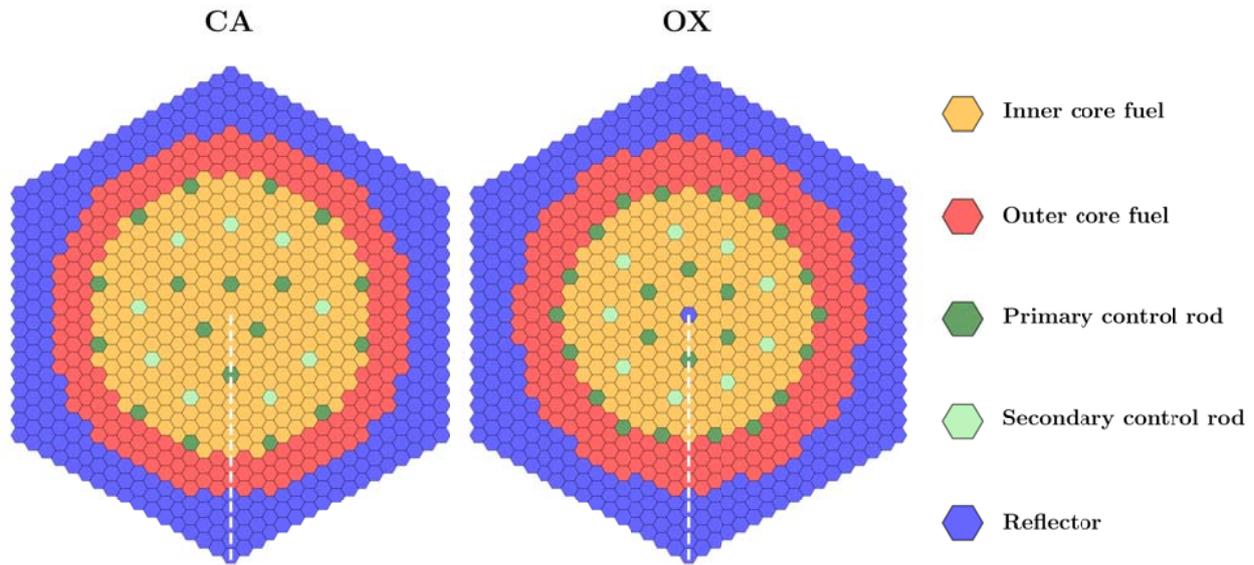


Figure 1: SFR Core Layouts.

utilizing the JEFF-3.1.2 library. Besides the benchmark criticality case Doppler constants (K_D) were calculated with 1000 K fuel temperature raise * (ΔT_f) for each configuration:

$$K_D = T_f \frac{d\rho}{dT} \approx \frac{\rho_1 - \rho_0}{\ln(T_{f,1}/T_{f,0})} \text{ (pcm)}, \quad (1)$$

where index 0 refers to the benchmark case, index 1 denotes the perturbed case and ρ is reactivity. The sodium void reactivity (SVR) and CRW values were obtained by simply taking the difference between the nominal and the perturbed cases as $SVR = \rho_0 - \rho_{SV,100}$ and $CRW = \rho_0 - \rho_{CRs,in}$, respectively. The axial fuel and radial core expansion reactivity feedback was calculated with the same logic as SVR and CRW comparing the two cases whereas the L_f length of the fuel slug had to be modified in the Serpent model by:

$$L_f(T_1) = L_f(T_0) [1 + \alpha_{CA/OX}(T_1 - T_0)] \text{ (cm)}, \quad (2)$$

yielding a $\Delta L_f/L_f$ of + 1.29/1.77 % for CA/OX with corresponding fuel density reduction, where $\alpha_{CA/OX}$ is the linear expansion coefficient of the fuel [4,5]. Radial expansion of core diagrid was derived from the pitch (p) perturbation:

$$p_1(T_1) = p_0(T_0) [1 + \alpha_{rad,g}(T_1 - T_0)] \text{ (m)}, \quad (3)$$

giving a $\Delta p/p$ of + 0.99 % for both cases based on the same 316L SS structural material, where $\alpha_{rad,g}$ was the global linear expansion coefficient of the diagrid.

*The benchmark suggested doubling the fuel temperature for both CA/OX cores however admittedly such temperatures were not physical. In order to approximate these temperatures as close as possible and keep temperatures below fuel melting point the 1000 K increase was applied.

2.3. Serpent Results

Results were compared with benchmark data summarized in [6] and studies conducted by Nikitin et al. [7,8] with Serpent. The benchmark averages are distorted by codes with homogeneous cell description while [7,8] represent a heterogeneous approach. One has to account for inherent differences of code versions and boundary conditions e.g. cross section libraries, Table 2 summarizes the results. In general the discrepancies are acceptable considering the ± 500 pcm uncertainty ranges in [6] however the SVR values highlight larger differences. Axial fuel and radial diagrid expansion feedback indicated that the cores are more sensitive to pitch perturbation than thermal elongation of the fuel slug. Regarding statistics, present Serpent calculations showed an average deviation of ± 1.4 pcm and Ref. [7,8] incorporated ± 2.0 pcm uncertainty. Available standard deviation values ($\pm \sigma$) of the benchmark cases are included in Table 2 below.

Table 2: 3D Full-Core Results for CA/OX Cores.

Reference	k_{eff} (-)	K_D (pcm)	SVR (pcm)	CRW (pcm)	$\Delta\rho_{\text{ax}}$ (pcm)	$\Delta\rho_{\text{rad}}$ (pcm)
[6]	1.0136/1.0096	-1002/-895	2048/1932	-4326/-6092	-	-
$\pm \sigma$ [6]	0.0090/0.0048	167/89	398/171	1124/995	-	-
Present model	1.0067/1.0120	-945/-936	1552/1336	-4217/-6127	-179/-103	-421/-405
[7,8] (OX)	1.0107	-852	1864	-6046	-120	-429

3. FEW-GROUP CONSTANTS' GENERATION FOR ANTS

3.1. 2D Models and Boundary Conditions

Supercell and colorset lattices were composed in order to obtain 2D homogenized few-group cross sections of non-fissile and fissile cells. In cases of interfaces between fissile and non-fissile cells colorset models were written for subchannels e.g. outer core FA with one, two or three adjacent radial reflector cells. Fuel lattices had periodic boundary conditions while other models e.g. control rod models were treated with reflective boundary conditions.

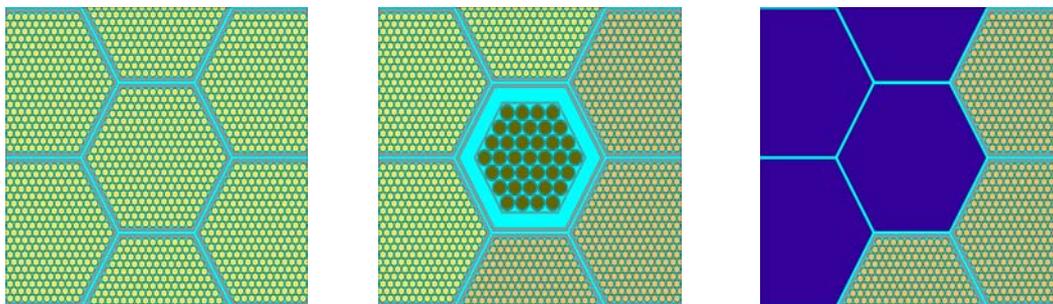


Figure 2: Serpent 2D Inputs (left - inner core infinite lattice, middle - primary CR, right - radial reflector)

Regarding the few-group energy condensation the ECCO-33 structure was considered first, afterwards a simplified 24-group equal-lethargy bin structure was derived lumping the thermal region

groups into one group (24th) hence reducing statistical errors of the lowest energy groups [9]. The SCALE238 structure was used for micro-group condensation.

A sensitivity study of homogenization methods was carried out as well, considering B_1 and P_1 leakage corrections and FM leakage correction with out-scattering diffusion coefficients. The study utilized a CA inner fuel 2D infinite lattice model in order to demonstrate the bias of the methods (see Fig. 2 left). Each run had 2 500 cycles (2 000 active / 500 passive) with a population of 60 000 neutrons. Calculations were performed with varying neutron populations between 2 000 – 100 000 whereas it was concluded that statistical errors did not show noticeable variance with populations over 60 000 particles, later on the same boundary conditions were applied in XS generation for Ants.

3.2. Few-Group Cross Section Generation Sensitivity Study

Group constants e.g. Σ_a , $\nu\Sigma_a$, $\kappa\Sigma_f$ and D describe materials in Ants, Fig. 3 represents the trends of the mentioned terms on the 10^{-11} –20 MeV spectrum. It can be seen that for most cases the various methods showed little to negligible discrepancies except for the 24th energy-group where the B_1 leakage correction had the lowest XSs systematically. Regarding Σ_a , $\nu\Sigma_f$, $\kappa\Sigma_f$ the average and maximum difference was $-0.24 / 0.64$ % between no-leakage and corrected values, respectively. As for the diffusion coefficient the maximum average difference was 1.52 % given by the P_1 method followed by B_1 (0.33 %) and out-scattering (0.07 %) that were noticeably lower than the P_1 error.

These tendencies confirm an earlier conclusion of Faure et al. [10] whereas leakage-corrected XS generation methods were studied in in fast systems. It was reported that the leakage models had a low impact on cross section generation. Such insensitivity can be explained by the fact that the introduction of a leakage term does not change noticeably the structure of the local multi-group flux within a macro-group. Considering scattering cross sections, in general for the scattering matrix the maximum error between leakage-corrected and no-leakage methods was -0.23 % (P_1).

4. ANTS REDUCED-ORDER NODAL DIFFUSION SOLVER

4.1. Code and Input Structure

The development of the Ants reduced-order nodal diffusion solver has started at VTT in 2018 [11,12]. The in-house built code is based on the analytic function expansion nodal (AFEN) - and function expansion nodal methods (FENM) [13,14]. The algorithm tackles the large number of nodal variables without transverse integration which is beneficial considering hexagonal geometry but has a negative impact on solver speed. The expansion of the intranodal flux is based on analytic basis functions which' solutions are obtained from the steady-state multi-group diffusion equation. Regarding the model structure the user has to provide basic geometry data e.g. core lattice pitch, number and length of axial nodes with respect to benchmark data [1]. The material of each node is described by a set of cross sections in a 24-group condensed format, extracted from the 2D supercell results as mentioned above. Finally convergence limits can be defined for k_{eff} and fission source for instance, calculations presented in Section 4.2 had a $\varepsilon_{k_{\text{eff}}} = 10^{-8}$ and $\varepsilon_S = 10^{-6}$ limits.

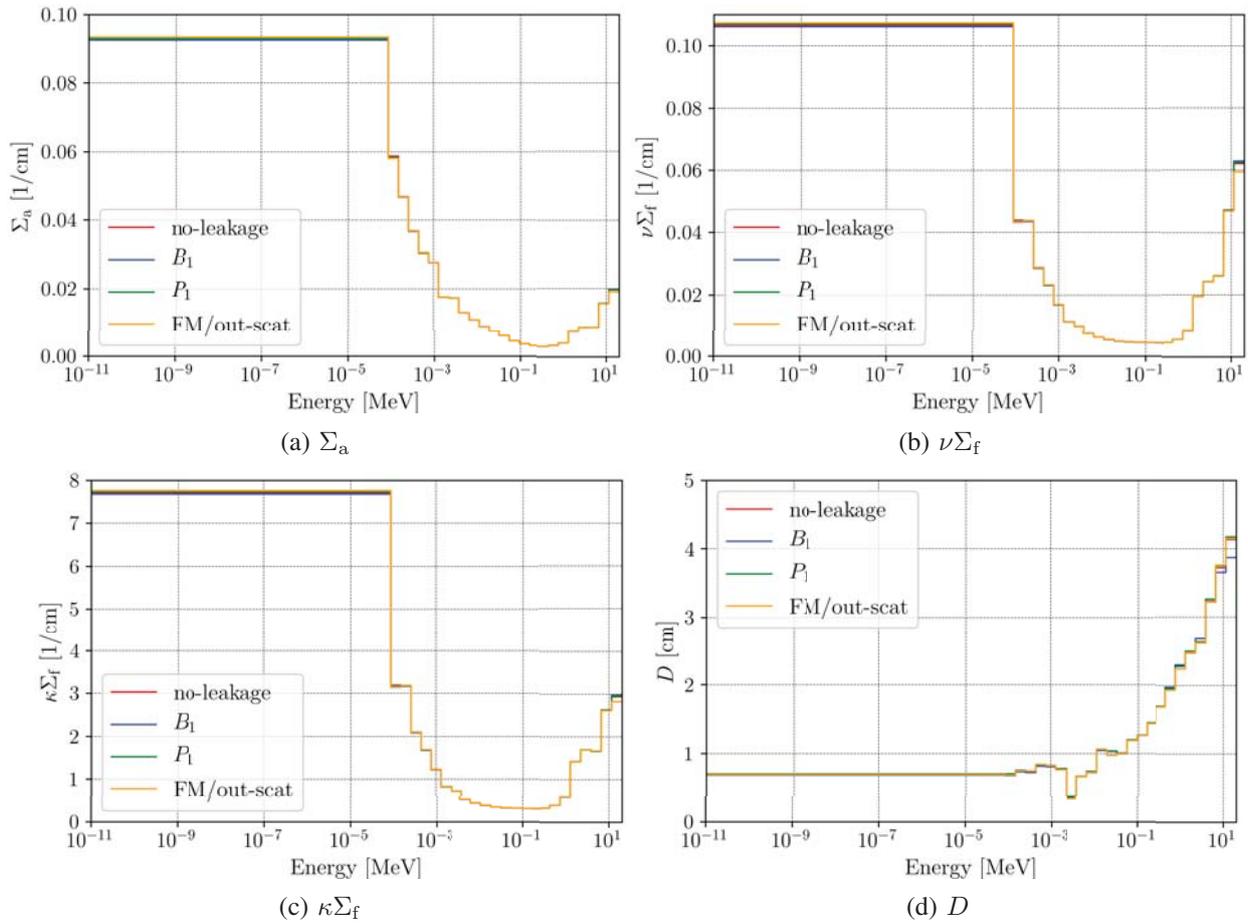


Figure 3: Group Constants in CA Inner Fuel 2D Supercell.

4.2. Comparison of Ants and Serpent

In order to compare power distribution of Ants and Serpent cores the assembly-wise relative powers were normalized to the maximum value in each solution, the differences between these normalized distributions are depicted in Fig. 4. Comparing Serpent to Ants the average and absolute maximum error was 0.00 / 1.52 % in case of CA and 0.19 / 2.66 % for OX, respectively.

The axial power distribution comparison is shown in Fig. 5 highlighting that Ants provided a more symmetric distribution for CA than Serpent meaning that the CRs had smaller worth in the reduced-order model as discrepancies started to increase in the upper section of the core. The OX core followed the axial profile of Serpent better in the upper part, noticeable +0.39 % and +1.85 % for CA and +0.52 % and +2.23 % for OX, respectively.

The Ants radial power profiles (see Fig. 6) show good agreement with the MC results with largest absolute error of 1.26 % for CA and 1.86 % for OX, comparing Ants to MC. The label "Ants-simp." denotes a simplified case where the local geometry was neglected on the radial reflector/outer core boundary i.e. outer core FA 2D infinite lattice cross sections were used for fuel assemblies with

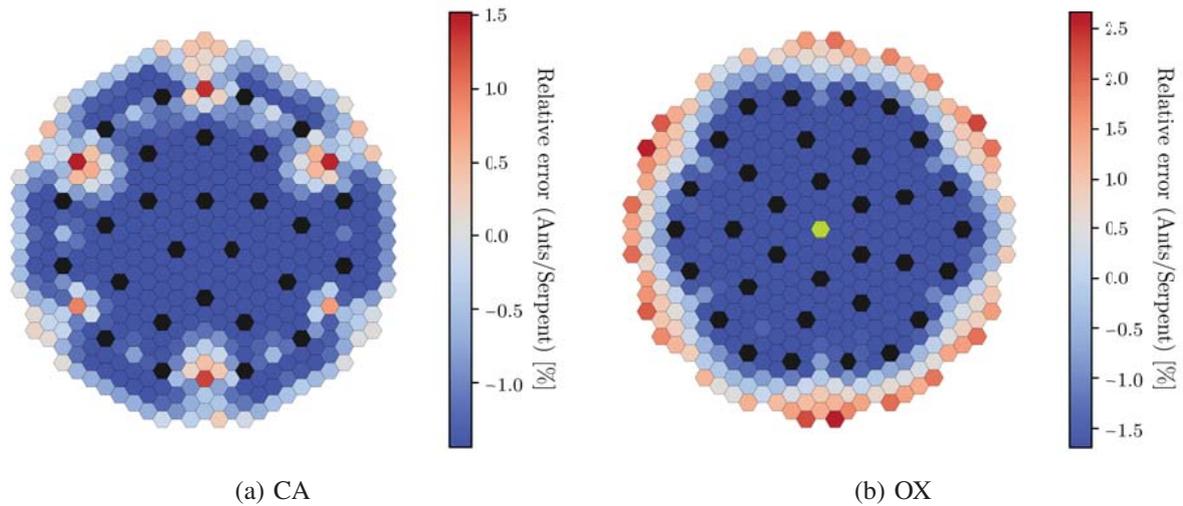


Figure 4: Comparison of Power Distributions in Criticality Case.

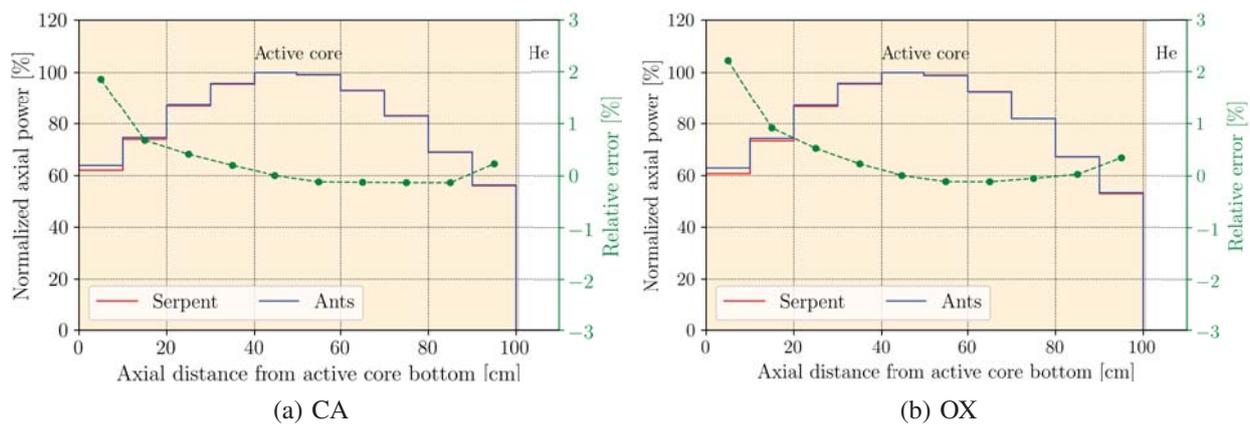


Figure 5: Axial Power Profiles in Criticality Case.

reflector neighbours. The deviation is clear between the two CA Ants results indicating that local heterogeneities affecting flux did alter the global results, moreover, the maximum absolute error increased to 2.17 %. The OX results are similar comparing maximum absolute errors where the finer mesh and the simplified case gave 1.87 % and 3.03 %, respectively. Hence the OX system with larger flux gradient on the outer core edge was more sensitive to perturbation from mesh resolution.

Table 3 summarizes the Ants-Serpent comparisons considering several test cases described in Section 2.3. Regarding the criticality case the discrepancies between k_{eff} results were 7 pcm (CA) and 136 pcm (OX). Even though the OX error is significantly larger than its CA counterpart, earlier studies on the OX core showed similar discrepancy ranges with Serpent-PARCS (84 pcm) and Serpent-DYN3D (128 pcm) [7]. Has to be mentioned that using superhomogenization (SPH) the

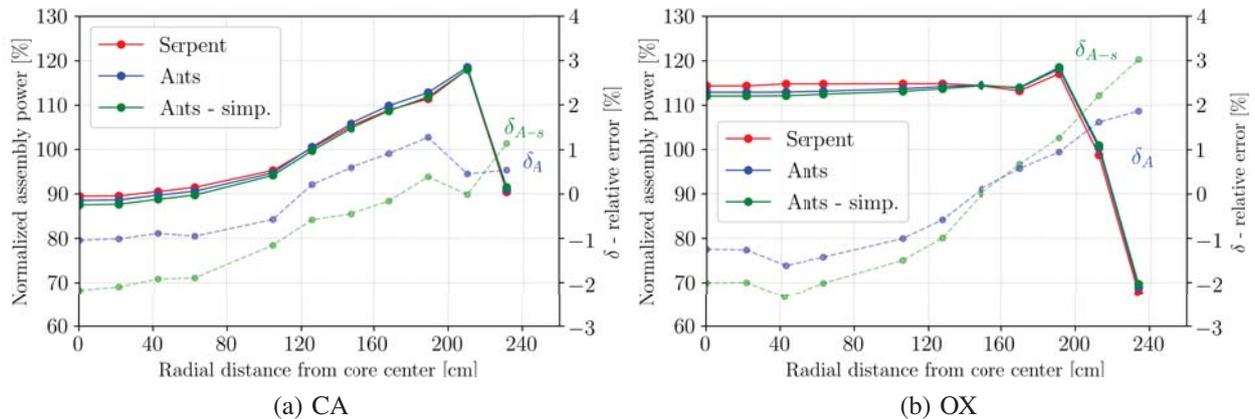


Figure 6: Radial Power Profiles (in Main Diagonal).

nodal solution can be significantly improved as reported in [15]. SVR coefficients show larger departure comparing Ants to Serpent, further analysis is required to study the effect of self-shielding in heterogeneous and homogeneous codes. Looking at CRW, $\Delta\rho_{ax}$ and $\Delta\rho_{rad}$ the trends are acceptable in general although $\Delta\rho_{ax,OX}$ did not show satisfying similarity between the high fidelity and the reduced-order values. This underlines that the depiction (by 2D homogenization) of a composition-perturbation in fuel material can translate into varying discrepancies comparing stochastic and deterministic codes.

Table 3: 3D Full-Core Results for CA/OX with High-Fidelity and Reduced-Order Solvers.

Code	k_{eff} (-)	K_D (pcm)	SVR (pcm)	CRW (pcm)	$\Delta\rho_{ax}$ (pcm)	$\Delta\rho_{rad}$ (pcm)
Serpent	1.00670/1.01198	-945/-936	1552/1336	4217/6127	-179/-106	-421/-405
Ants	1.00806/1.01241	-947/-957	1895/1290	4212/6180	-182/-159	-414/-404

5. CONCLUSIONS

The Serpent full-core 3D simulation of two SFR cores was presented and compared to benchmark data [6] where Serpent showed satisfying agreement with previous works considering heterogeneous geometry description. The 2D supercell models highlighted the insensitivity of few-group constants to leakage-correction methods in a fast system. The methodology and performance of the Ants reduced-order diffusion solver was demonstrated as a potential surrogate of Serpent. Reactivity feedbacks showed good agreement with benchmark data and MC results however larger differences were observed in cases of SVRs and axial fuel expansion of the OX core indicating that the applied homogenization had larger impact on composition-related perturbations compared

to changes in geometry. Further studies are needed to improve the efficiency of group constant generation with special care to fissile and non-fissile cell interfaces.

References

- [1] D. Blanchet, L. Buiron, N. Stauff, T. K. Kim, T. Taiwo, *AEN - WPRS Sodium Fast Reactor Core Definitions (v1.2)*, OECD/NEA, 19th September 2011.
- [2] J. Leppänen, M. Pusa, T. Viitanen, V. Valtavirta, T. Kaltiaisenaho, “The Serpent Monte Carlo code: Status, development and applications in 2013” *Annals of Nuclear Energy*, **Volume 82**, pp. 142-150 (2015).
- [3] Terje Tverberg, *Mixed-oxide (MOX) fuel performance benchmark, Summary of the Results for the Halden Reactor Project MOX Rods, NEA/NSC/DOC(2007)6*, OECD HRP, 2007.
- [4] П. Л. Кириллов, Теплофизические свойства материалов ядерной техники - учебное справочное пособие 2-е изд., Обнинский Институт Атомной Энергетики, Обнинск, Russian Federation (2007).
- [5] Ю. Г. Годин, А. В. Тенишев, Карбидное ядерное топливо, Московский Инженерно-Физический Институт, Москва, Russian Federation (2007).
- [6] L. Buiron et al., “Evaluation of Large 3600 MWth Sodium-Cooled Fast Reactor Neutronic OECD benchmarks” *Proceedings of PHYSOR 2014*, Kyoto, Japan, September 28 - October 3, (2014).
- [7] E. Nikitin, E. Fridman, K. Mikityuk, “Solution of the OECD/NEA neutronic SFR benchmark with Serpent-DYN3D and Serpent-Parcs code systems” *Annals of Nuclear Energy*, **Volume 75**, pp. 492-497 (2014).
- [8] E. Nikitin, E. Fridman, “Extension of the reactor dynamics code DYN3D to SFR applications - Part I: Thermal expansion models” *Annals of Nuclear Energy*, **Volume 119**, pp. 382-389 (2018).
- [9] E. Fridman, “Generation of few-group constants with Serpent: Application examples” *Proceedings of PHYSOR 2014*, Kyoto, Japan, September 28 - October 3, (2014).
- [10] B. Faure, G. Marleau, “Simulation of a sodium fast core: Effect of B_1 leakage models on group constant generation” *Annals of Nuclear Energy*, **Volume 99**, pp. 484-494 (2017).
- [11] V. Sahlberg, A. Rintala, “Development and first results of a new rectangular nodal diffusion solver of Ants” *Proceedings of PHYSOR 2018*, Cancún, Mexico, 22-26 April, (2018).
- [12] A. Rintala, V. Sahlberg, “Extension of nodal diffusion solver of Ants to hexagonal geometry” *Kerntechnik*, **Volume 84**, pp. 252-261 (2019).
- [13] N. Z. Cho, J. M. Noh, “Analytic function expansion nodal method for hexagonal geometry” *Nuclear Science and Engineering*, **Volume 121**, pp. 245-253 (2014).
- [14] B. Xia, Z. Xie, “Flux expansion nodal method for solving multigroup neutron diffusion equations in hexagonal-z geometry” *Annals of Nuclear Energy*, **Volume 33**, pp. 370-376 (2006).
- [15] E. Nikitin, E. Fridman, K. Mikityuk, “On the use of the SPH method in nodal diffusion analyses of SFR cores” *Annals of Nuclear Energy*, **Volume 85**, pp. 544-551 (2015).