INTRODUCTION OF OSCAR-4 AT THE HIGH FLUX REACTOR (PETTEN)

B. Erasmus1,*, J.A. Hendriks1, A. Hogenbirk1, S.C. van der Marck1, and N.L. Asquith1

1 Nuclear Research and consultancy Group
P.O. Box 25, 1755 ZG, Petten, The Netherlands
erasmus@nrg.eu*, hendriks@nrg.eu, hogenbirk@nrg.eu, vandermarck@nrg.eu, asquith@nrg.eu

ABSTRACT

Since 2005 the nodal diffusion based code system, OSCAR-3, was used for reactor support calculations of operational cycles of the High Flux Reactor in Petten, The Netherlands. OSCAR uses a two-step deterministic calculation, in which homogenized cross sections are generated in lattice environments using neutron transport simulations, and then passed to a nodal diffusion core simulator to model the full reactor. Limitations in OSCAR-3 led to the need for improved modelling capabilities and better physics models for components present in the reactor core. OSCAR-4 offers improvements over OSCAR-3 in its approach to homogenization, and the new version of the diffusion core simulator allows for better modelling of movable components such as control rods. Fuel inventories calculated using OSCAR-4 can also easily be exported to MCNP, which allows the calculation of individual plate powers and local reaction rates amongst others. For these reasons OSCAR-4 is currently being introduced as a core support tool at the High Flux Reactor. In this work the steps that were followed to validate the reactor models are presented, and include results of validation calculations from both OSCAR-4 and MCNP6 over multiple reactor cycles. In addition differences in cross section library evaluations and their impact on the results are presented for the MCNP model.

KEYWORDS: validation, nodal diffusion, OSCAR-4, MCNP

1. INTRODUCTION

The High Flux Reactor (HFR) is extensively used for the production of a variety of isotopes, fuel testing and material irradiation. This means that the reactor core loading not only varies from cycle to cycle, but also changes during the cycle as isotope targets are loaded and unloaded during operation of the reactor. Contributing to the challenge of modelling the reactor accurately are the following characteristics: the presence of localized burnable absorbers (in the form of cadmium wires in the fuel assemblies); the fuel-follower design of the control rods, which sees parts of the fuel extending below the reactor; and increased power peaking as a result of the 17 in-core irradiation positions.

The nodal diffusion based code system, OSCAR-3 [1], was used for core follow and core design calculations for each operational cycle of the HFR since 2005. OSCAR uses a two-step deterministic calculation, by first generating homogenized cross sections in lattice environments with neutron transport simulations, and then using these cross sections in a nodal diffusion core simulator. Although the introduction of OSCAR-3 was a major improvement at the time, the increased use of the reactor, and
limitations in version three of the OSCAR software, led to the need for improved modelling capabilities and better physics models. The fourth version of this software, OSCAR-4 [2, 3], offers improvements over OSCAR-3 in its approach to homogenized cross sections, and improved physics models in the nodal diffusion core simulator MGRAC.

MGRAC uses the multi-group analytic nodal method and generalized equivalence theory, to solve the transverse integrated 3D diffusion equation. The equivalence parameters needed by MGRAC, such as homogenized cross sections, discontinuity factors and flux form factors, collectively referred to as homogenized cross sections here, are generated in the homogenization step in the lattice calculations. Furthermore, this new version of MGRAC allows for the modelling of the fuel follower control elements to their full extent, by capturing the power and flux generated by the fuel followers below the active fuel section of the core. In addition, fuel inventories calculated using OSCAR-4 can easily be exported to create an associated MCNP model, when detailed calculations such as fuel plate power distributions and localized reaction rates are required for copper wires activation experiments.

In the following sections the steps that were followed to validate the OSCAR-4 and associated MCNP6 [4] models are described, and results of validation calculations from both models are presented for multiple reactor cycles.

2. MODELLING APPROACH

For the purpose of generating homogenized cross sections, Serpent-2.1.23 [5] and MCNP4C3 [6, 7] were used in conjunction with HEADE, the lattice code that forms part of the OSCAR-4 code system. Serpent and MCNP were used to generate cross sections for ex-core and in-core non-fuel components, respectively, while HEADE was used to generate cross sections for fueled components. Nuclear data libraries used with each of the transport codes during cross section homogenization are listed in Table I. The collection of libraries used with MCNP4C3 together make up the HFR standard library, currently used for routine reactor support calculations.

At the time these models were developed, the functionality to use microscopic burnup dependent cross sections from Serpent for fueled assemblies was not yet in place. This restriction can be lifted in future versions of the model, where Serpent can be used to generate homogenized cross sections for both fueled and non-fuel components. All homogenized cross sections were generated in a 7 group energy structure.

<table>
<thead>
<tr>
<th>Transport code</th>
<th>Nuclear data library</th>
</tr>
</thead>
<tbody>
<tr>
<td>HEADE</td>
<td>172-group WIMS-E library, based mainly on JEFF-2.2 [8]</td>
</tr>
<tr>
<td>Serpent-2.1.23</td>
<td>ENDF/B-VII.1 [9]</td>
</tr>
<tr>
<td></td>
<td>ENDF/B-VI.5 [10] for most isotopes; JEFF-2.2 used for Al-27, U-235 and H bound in H$_2$O; JENDL-3.3 [11] for Cadmium isotopes. (This combination of libraries is designated as the HFR standard library)</td>
</tr>
<tr>
<td>MCNP4C3</td>
<td></td>
</tr>
</tbody>
</table>

Fuel assemblies were homogenized with HEADE, using two different lattice environments: the first containing a fresh fuel element with two partially spent fuel elements on either side (see Figure 1(a)), and
a second lattice environment in which a representation of an in-core experiment was modelled on either side of the fuel assembly (see Figure 1(b)). In both lattice environments isotropic reflection was used as (white) boundary conditions. These two environments aim to approximate the core environment.

**Figure 1. Extended lattice environments used to generate homogenized cross sections for the fuel assemblies.**

Homogenized cross sections for the fuel follower part of the control rods were also generated with HEADE, while the cross sections for the cadmium absorber part were generated with Serpent. In both cases, homogenized cross sections were generated in a 3x3 lattice environment shown in Figure 2(a) for the fuel follower section and Figure 2(b) for the cadmium absorber section.

**Figure 2. Extended 3x3 lattice environment used to generate homogenized cross sections for the fuel follower (a) and cadmium absorber section (b) of the control rods.**

In-core non-fuel component cross sections were generated using MCNP4C3, in 3x3 lattice environments including fuel driver zones, depending on their location in the core. These models were also used
historically with the OSCAR-3 model of the HFR, and do not include discontinuity factors. In contrast to this, in-core beryllium reflector cross sections and ex-core structure and pool cross sections were generated from a full core Serpent model as depicted in Figure 3. In the figure, the nodal diffusion mesh in the $x$ and $y$ directions used in MGRAC, is superimposed over the geometry in red. Homogenized cross sections were generated for all components in various layers, accounting for material changes in the $z$ direction.

![Figure 3. Full core Serpent model with nodal overlay mesh, used to generate homogenized cross sections for the non-fuel components.](image)

After the homogenization step, assembly models were constructed by stacking layers of homogenized cross sections on top of each other, which correspond to the different axial zones in the assemblies. Finally, these assembly models were combined into a single MGRAC model, describing the entire reactor.

MCNP6 model, full detail, OSCAR

### 3. SIMULATION RESULTS

#### 3.1. Core Reactivity

In order to evaluate the calculation of core reactivity, two time points were identified in 17 different operational cycles, each corresponding to a critical configuration. For each of the identified time points, a MGRAC and MCNP model was built representing the critical configuration, taking into account the fuel depletion, control rod settings, in-core and ex-core experiments and isotope production loadings. Figure 4 shows the calculated core $k_{eff}$ values for the chosen critical configurations. For the OSCAR-4 model, the average $k_{eff}$ for the 34 critical states is $0.99974 \pm 0.00177$, while the average $k_{eff}$ value for the MCNP model with the HFR standard library is $0.99546 \pm 0.00110$. This difference between OSCAR-4 and
MCNP models decreases by 150 pcm, when an explicit list of nuclides is used to represent all fission products instead of a B-10 equivalent number density for a lumped fission product, when exporting the fuel inventory from OSCAR to MCNP.

![Figure 4](https://doi.org/10.1051/epjconf/202124710029)

**Figure 4.** Calculated core reactivity comparison between OSCAR-4, MCNP with the HFR nuclear data library, and MCNP with the JEFF-3.3 nuclear data library for critical states in cycles 2017-01 to 2018-08.

In order to investigate the effect of nuclear data libraries on the criticality results, the HFR standard library was replaced with the JEFF-3.3 library [12] in the MCNP model. Changing libraries resulted in a lower average $k_{\text{eff}}$ of 0.98989 ± 0.00092 over the 34 critical states. This difference was further investigated and it was found that U-235 and Al-27 were the main isotopes contributing to this difference. For a selected critical case, replacing Al-27 and U-235 in the HFR standard library with those from JEFF-3.3, resulted in a 391 pcm and 232 pcm decrease in $k_{\text{eff}}$ respectively. Other isotopes also contribute to the difference but their impact is much less significant. The difference in the average $k_{\text{eff}}$ of the MCNP model, when switching from the HFR library to the JEFF-3.3 library is 565 pcm.

### 3.2. Control Rod Calibrations

Prior to reactor cycle 2018-08, control rod calibration experiments were performed. In these experiments the differential and integral worth of the control rods were measured over a parts of their travel distance. For control rod S1, nearly the entire travel distance was covered, while for control rods S2 to S6 measurements were done for the top 15 cm of travel distance. These experiments were then simulated with the OSCAR-4 and MCNP models. Figure 5 shows a comparison of the measured differential and integral worth curves of control rod S1, with those calculated with the OSCAR-4 and MCNP models. Values for the measured and calculated integral worth of all of the control rods are given in Table II. The values presented in brackets for MCNP are the statistical uncertainties. Measurement uncertainty data was only provided for the measurement of control rod S1, and is also reported in brackets in Table II.

Both the OSCAR-4 and MCNP models show good agreement with the measured worth of control rod S1, with the MCNP model slightly overestimating the total worth. The OSCAR-4 model overestimates the partial worth of control rods S3 and S4, while underestimating the partial worth of control rod S5.
Because only partial worth measurements were available for rods S2 to S6, no conclusions can be drawn on how well the models predict the worth of these rods over the entire travel distance.

![Graph](image)

**Figure 5. Comparison of the OSCAR-4, MCNP and measured differential and integral worth of control rod S1.** The differential worth is represented by scatter plots, while the integral worth is represented by solid line plots.

**Table II. Integral control rod worth:** Comparison of OSCAR-4, MCNP and measured integral control rod worth for all six control rods.

<table>
<thead>
<tr>
<th>Control rod label</th>
<th>OSCAR-4 (cents)</th>
<th>MCNP (cents)</th>
<th>Measured (cents)</th>
</tr>
</thead>
<tbody>
<tr>
<td>S1</td>
<td>493.14</td>
<td>530.25 (3.25)</td>
<td>491.95 (3.58)</td>
</tr>
<tr>
<td>S2</td>
<td>24.66</td>
<td>26.75 (3.61)</td>
<td>23.04</td>
</tr>
<tr>
<td>S3</td>
<td>15.75</td>
<td>10.71 (3.42)</td>
<td>13.45</td>
</tr>
<tr>
<td>S4</td>
<td>13.42</td>
<td>9.31 (3.71)</td>
<td>10.95</td>
</tr>
<tr>
<td>S5</td>
<td>24.01</td>
<td>30.73 (3.71)</td>
<td>30.10</td>
</tr>
<tr>
<td>S6</td>
<td>20.38</td>
<td>21.14 (3.42)</td>
<td>22.61</td>
</tr>
</tbody>
</table>

3.3. **Copper Wire Activation**

Prior to chosen operational cycles, flux mapping measurements were performed by activation of copper wires. For these experiments, copper wires were placed at two or three different locations in fifteen fuel assemblies and centrally in four of the control rods. These experiments measured the relative Cu-63 (n, γ) reaction rates, which were then normalized to the to the MCNP results.

In the MCNP model, the Cu-63 (n, γ) reaction rates were tallied at the position of the copper wires, and compared with the measured reaction rates. For the purpose of comparison, the measured and calculated reaction rates were integrated over ±28 cm with respect to center line core, from which C/E values were calculated. Results for two flux mapping experiments are given in Figure 6 and Figure 7, which were done prior to operational cycles 2017-07 and 2018-08, respectively.
Figure 6. Relative differences in the calculated and experimental Cu-63 reaction rates for cycle 2017-07 (C/E-1, in %). Results are integrated between ±28 cm with respect to core center line.

Figure 7. Relative differences in the calculated and experimental Cu-63 reaction rates for cycle 2018-08 (C/E-1, in %). Results are integrated between ±28 cm with respect to core center line.
The calculated reaction rates are typically within ±5% of the measured values, with notable outliers for the control rods in positions B4 and B6, and the wires positioned in the southern part of assemblies H3 and D7. In both the results from cycle 2017-07 and 2018-08, an east-west tilt in the flux is visible, with reaction rates in elements towards the east generally being under predicted, while the reaction rates in elements to the west generally being over predicted by the MCNP model. The cause of this tilt over the core is currently unknown and requires further investigation.

4. CONCLUSIONS

The results presented for critical states over multiple operational cycles of the HFR, show that the OSCAR-4 and MCNP models produce reasonably stable $k_{\text{eff}}$ values. It was also shown that changing from the HFR standard nuclear data library to JEFF-3.3 libraries leads to a decrease in the average $k_{\text{eff}}$ values of 565 pcm. The main isotopes contributing to this difference were found to be Al-27 and U-235.

Simulation results from both OSCAR and MCNP show acceptable comparison with measured values for various experiments, making the models suitably predictive for reactor support. However, in MCNP model there is a minor east-west tilt in the flux across the core, as seen from the copper wire activation experiments. The cause of the tilt in the flux is possibly due to depletion of the cadmium sections of the control rods being neglected. Causes of the bias present in the models, when using a modern nuclear data library, require further investigation.

REFERENCES