

# Observation of the $\gamma\gamma \rightarrow \tau^+\tau^-$ production in ultraperipheral PbPb collisions with the CMS experiment

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**Abstract.** We report the observation of the photon-induced  $\gamma\gamma \rightarrow \tau^+\tau^-$  production based on a data sample of  $404 \mu\text{b}^{-1}$  collected by the CMS experiment at a per nucleon center-of-mass energy of 5.02 TeV. The cross section is measured in a fiducial phase space region via the decay process involving one muon and three charged hadrons, and is found to be  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-) = 4.8 \pm 0.6$  (stat)  $\pm 0.5$  (syst)  $\mu\text{b}$ , in agreement with leading-order quantum electrodynamics predictions. The  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-)$  measurement is used to determine the anomalous magnetic moment of the  $\tau$  lepton  $a_\tau$ , which is currently poorly constrained.

## 1 Introduction

Ultraperipheral collisions (UPC) of heavy ions provide an extremely clean environment to probe the  $\tau$  lepton anomalous magnetic moment  $a_\tau$  at the LHC, hence offering ample room for searching for physics beyond the standard model. Here, we present [1] the observation of the photoproduction of pairs of  $\tau$  leptons in lead-lead (PbPb) collisions  $\text{PbPb}(\gamma\gamma) \rightarrow \text{Pb}^{(*)} + \text{Pb}^{(*)} \tau^+\tau^-$  (hereafter referred to as  $\gamma\gamma \rightarrow \tau^+\tau^-$ ) by analyzing a data sample collected by the CMS experiment [2] at a per nucleon center-of-mass energy of  $\sqrt{s_{\text{NN}}} = 5.02$  TeV in 2015, and corresponding to an integrated luminosity of  $404 \mu\text{b}^{-1}$ . The  $\tau$  leptons are reconstructed in a decay process involving in the final state one muon and three particles (“prongs”) assumed to be charged  $\pi$  mesons, and the cross section  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-)$  is measured in a fiducial phase space region defined by the muon and pion transverse momenta ( $p_{\text{T}}$ ) and their pseudorapidities ( $\eta$ ). The ATLAS Collaboration has also reported a measurement of  $\gamma\gamma \rightarrow \tau^+\tau^-$  using a larger PbPb data sample with an integrated luminosity of  $1.44 \text{ nb}^{-1}$  [3].

## 2 Event selection, signal and background estimation

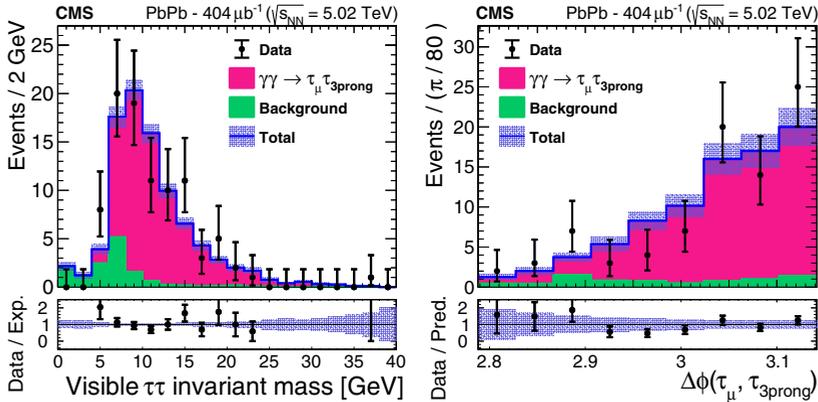
The exchange of photons in UPC does not primarily disassociate the lead ions but scatters them at small angles close to the beam direction, such that the final state is uniquely characterized by low track multiplicity. We thus filter events from UPC by requiring in real time (“online”) a single muon with no  $p_{\text{T}}$  threshold requirement and minimum event activity above the noise threshold in the forward hadron (HF) calorimeter ( $3.0 < |\eta| < 5.2$ ). A further event selection is applied offline for the “ $\tau_\mu$ ” and “ $\tau_{3\text{prong}}$ ” candidates, which at the same time defines the  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-)$  fiducial phase space region as summarized in Table 1.

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A dedicated  $\gamma\gamma \rightarrow \tau^+\tau^-$  [4] Monte Carlo (MC) sample is generated with MADGRAPH5\_AMC@NLO (v2.6.5) [5], decayed and hadronized with PYTHIA8 (v2.1.2) [6], and simulated by GEANT4 [7] to study the detector effects. The background contamination is estimated in a data-driven way using uncorrelated phase space regions defined based on the track multiplicity and the HF activity. In all cases, good agreement is observed between the measured distributions and the sum of the signal simulation and background estimation, e.g., as shown in Fig. 1 (left) for the invariant mass of the  $\tau_{3\text{prong}}$  candidate.

**Table 1.** Summary of the kinematic event selection, which also defines the fiducial phase space region for the  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-)$  measurement [1]. Additional quality criteria are applied for the  $\tau_\mu$  candidate and all selected tracks.

|                                  |   |
|----------------------------------|---|
| $\tau_\mu$ selection             | $p_T > 3.5 \text{ GeV}$ for $ \eta  < 1.2$<br>$p_T > 2.5 \text{ GeV}$ for $1.2 <  \eta  < 2.4$                              |
| Track selection                  | $p_T > 0.5 \text{ GeV}$ for the leading- $p_T$<br>$p_T > 0.3 \text{ GeV}$ for the (sub-)subleading- $p_T$<br>$ \eta  < 2.5$ |
| $\tau_{3\text{prong}}$ selection | $p_T^{\text{vis}} > 2 \text{ GeV}$ and $m_\tau^{\text{vis}} < 1.5 \text{ GeV}$  |



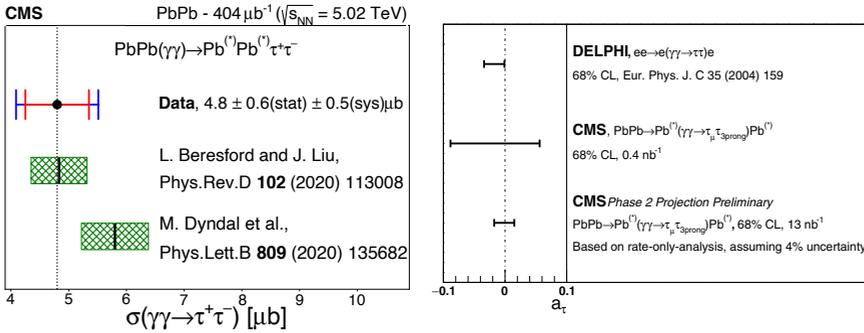
**Figure 1.** Left: Invariant mass of the three pions forming the  $\tau_{3\text{prong}}$  candidate [1]. Right: Difference in azimuthal opening angle between the  $\tau_\mu$  and  $\tau_{3\text{prong}}$  candidates [1]. The lower panels show the ratios of data to the signal-plus-background prefit expectation (left) or postfit prediction (right).

### 3 Results

A binned maximum likelihood method is used for the signal extraction. The difference in opening azimuthal angle between the  $\tau_\mu$  and  $\tau_{3\text{prong}}$  leptons,  $\Delta\phi(\tau_\mu, \tau_{3\text{prong}})$ , is used as the final discriminant. The sum in quadrature of systematic uncertainties, which may affect the normalization and the shape of the  $\Delta\phi(\tau_\mu, \tau_{3\text{prong}})$  distribution, is found to be 9.7%, with the statistical component contributing with an almost equal amount. To compare the compatibility of the data with the background-only and signal plus background hypotheses, where

the signal is allowed to be scaled by some factor  $r$ , we construct a test statistic based on the profile likelihood ratio. The best fit value of the signal strength is  $r = 0.99^{+0.16}_{-0.14}$ , with the number of signal events of  $N_{\text{sig}} = 77 \pm 12$ . The post-fit  $\Delta\phi(\tau_\mu, \tau_{3\text{prong}})$  distribution is shown in Fig. 1 (right).

The cross section is found to be  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-) = 4.8 \pm 0.6$  (stat)  $\pm 0.5$  (syst)  $\mu\text{b}$ , and it is measured in the fiducial phase space region, following the kinematic requirements listed in Table 1. The result, summarized in Fig. 2 (left), is compared to leading-order quantum electrodynamics (QED) predictions [4, 8]. Further assuming the correction factor of Ref. [4] to extrapolate the fiducial cross section measurement to the full phase space region, we then use the dependence of the total  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-)$  as a function of  $a_\tau$  to extract the model-dependent value of  $a_\tau = 0.001^{+0.055}_{-0.089}$ . The comparison of the world best  $a_\tau$  constraints [9] to the ones extracted from this analysis and its expected performance [10] in the High-Luminosity LHC program is given in Fig. 2 (right).



**Figure 2.** Left: The cross section,  $\sigma(\gamma\gamma \rightarrow \tau^+\tau^-)$ , measured in a fiducial phase space region at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV [1], and compared to theoretical predictions at leading-order accuracy in QED [4, 8]. Right: Comparison between the world best  $a_\tau$  constraints [9] and the ones extracted from this analysis [1] and its performance [10] using the integrated PbPb luminosity expected from the High-Luminosity LHC program.

## 4 Summary

In summary, an observation of the  $\tau$  lepton pair production in ultraperipheral nucleus-nucleus collisions is reported. Events with one muon and three particles identified as charged pions in the final state are reconstructed from a lead-lead data sample collected by the CMS experiment at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV in 2015, and corresponding to an integrated luminosity of  $404 \mu\text{b}^{-1}$ . This measurement introduces a novel experimental strategy using heavy ion data recorded by the LHC in order to measure the  $\tau$  lepton magnetic moment  $a_\tau$ . This approach can be used to extract constraints on  $a_\tau$  that surpass the current best from previous lepton-lepton colliders.

## Acknowledgements

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## References

- [1] CMS Collaboration, “Observation of  $\tau$  lepton pair production in ultraperipheral lead-lead collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV”, (2022). [arXiv:2206.05192](#). Accepted by *Phys. Rev. Lett.*
- [2] CMS Collaboration, “The CMS Experiment at the CERN LHC”, *JINST* **3** (2008) S08004, .
- [3] ATLAS Collaboration, “Observation of the  $\gamma\gamma \rightarrow \tau\tau$  process in PbPb collisions and constraints on the  $\tau$  lepton anomalous magnetic moment with the ATLAS detector”, (2022). [arXiv:2204.13478](#). Accepted by *Phys. Rev. Lett.*
- [4] L. Beresford and J. Liu, “New physics and  $\tau g - 2$  using LHC heavy ion collisions”, *Phys. Rev. D* **102** (2020) 113008, , [arXiv:1908.05180](#). [Erratum: [10.1103/PhysRevD.106.039902](#)].
- [5] J. Alwall et al., “The automated computation of tree-level and next-to-leading order differential cross sections, and their matching to parton shower simulations”, *JHEP* **07** (2014) 079, , [arXiv:1405.0301](#).
- [6] T. Sjöstrand et al., “An introduction to PYTHIA 8.2”, *Comput. Phys. Commun.* **191** (2015) 159, , [arXiv:1410.3012](#).
- [7] GEANT4 Collaboration, “GEANT4—a simulation toolkit”, *Nucl. Instrum. Meth. A* **506** (2003) 250, .
- [8] M. Dyndal, M. Klusek-Gawenda, M. Schott, and A. Szczurek, “Anomalous electromagnetic moments of  $\tau$  lepton in  $\gamma\gamma \rightarrow \tau^+\tau^-$  reaction in PbPb collisions at the LHC”, *Phys. Lett. B* **809** (2020) 135682, , [arXiv:2002.05503](#).
- [9] DELPHI Collaboration, “Study of  $\tau$  pair production in photon-photon collisions at LEP and limits on the anomalous electromagnetic moments of the  $\tau$  lepton”, *Eur. Phys. J. C* **35** (2004) 159, , [arXiv:hep-ex/0406010](#).
- [10] CMS Collaboration, “Snowmass White Paper Contribution: Physics with the Phase-2 ATLAS and CMS Detectors”, CMS Physics Analysis Summary CMS-PAS-FTR-22-001, 2022.