

# Scaling properties of background- and chiral-magnetically-driven charge separation: implications for the chiral magnetic effect in heavy ion collisions

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**Abstract.** The scaling properties of the  $R_{\Psi_2}(\Delta S)$  correlator and the  $\Delta\gamma$  correlator are used to investigate a possible chiral-magnetically-driven (CME) charge separation in  $p$ +Au,  $d$ +Au, Ru+Ru, Zr+Zr, and Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV, and in  $p$ +Pb ( $\sqrt{s_{NN}} = 5.02$  TeV) and Pb+Pb collisions at  $\sqrt{s_{NN}} = 5.02$  and 2.76 TeV. The results for  $p$ +Au,  $d$ +Au,  $p$ +Pb, and Pb+Pb collisions, show the  $1/N_{ch}$  scaling for background-driven charge separation. However, the results for Au+Au, Ru+Ru, and Zr+Zr collisions show scaling violations which indicate a CME contribution in the presence of a large background. In mid-central collisions, the CME accounts for approximately 27% of the signal + background in Au+Au and roughly a factor of two smaller for Ru+Ru and Zr+Zr, which show similar magnitudes.

Metastable domains of gluon fields with non-trivial topological configurations can form in the magnetized chiral relativistic quark-gluon plasma (QGP) [1] produced in collisions at RHIC and the LHC. The colliding ions generate the magnetic field ( $\vec{B}$ ) at early times [2]. The interaction of chiral quarks with the gluon fields can drive a chiral imbalance resulting in an electric current  $\vec{J}_V = \frac{N_c e \vec{B}}{2\pi^2} \mu_A$ , along the  $\vec{B}$ -field, i.e., perpendicular to the reaction plane;  $N_c$  is the color factor, and  $\mu_A$  is the axial chemical potential that quantifies the imbalance between right- and left-handed quarks. The resulting final-state charge separation, termed the chiral magnetic effect (CME) [1], is of great experimental and theoretical interest. However, its experimental characterization has been hampered by significant background.

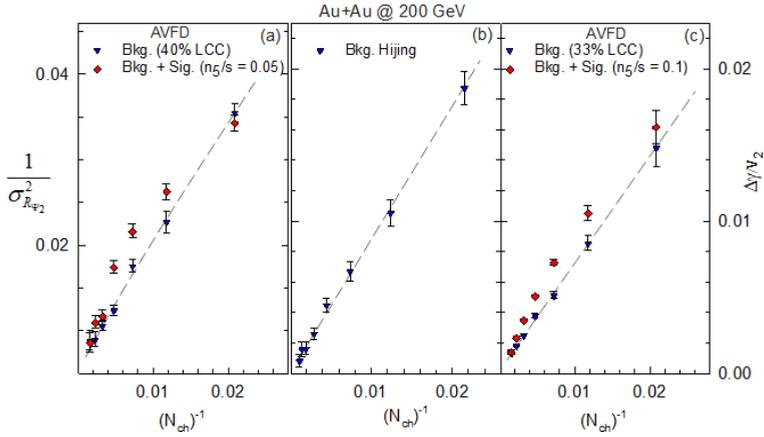
The charge separation can be quantified via the  $P$ -odd sine term  $a_1$ , in the Fourier decomposition of the charged-particle azimuthal distribution [3]:

$$\frac{dN_{ch}}{d\phi} \propto 1 + 2 \sum_n (v_n \cos(n\Delta\phi) + a_n \sin(n\Delta\phi) + \dots) \quad (1)$$

where  $\Delta\phi = \phi - \Psi_{RP}$  gives the particle azimuthal angle with respect to the reaction plane (RP) angle, and  $v_n$  and  $a_n$  denote the coefficients of the  $P$ -even and  $P$ -odd Fourier terms, respectively. A direct measurement of  $a_1$ , is not possible due to the strict global  $\mathcal{P}$  and  $\mathcal{CP}$  symmetry of QCD. However, their fluctuation and/or variance  $\tilde{a}_1 = \langle a_1^2 \rangle^{1/2}$  can be measured with charge-sensitive correlators such as the  $\gamma$ -correlator [3] and the  $R_{\Psi_2}(\Delta S)$  correlator [4].

The  $\gamma$ -correlator measures charge separation as:  $\gamma_{\alpha\beta} = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_2) \rangle$ ,  $\Delta\gamma = \gamma_{OS} - \gamma_{SS}$ , where  $\Psi_2$  is the azimuthal angle of the 2<sup>nd</sup>-order event plane which fluctuates

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**Figure 1.**  $\sigma_{R_{\Psi_2}}^{-2}$  vs.  $1/N_{\text{ch}}$  (a) and  $\Delta\gamma/v_2$  vs.  $1/N_{\text{ch}}$  (b) and (c) for simulated Au+Au collisions at  $\sqrt{s_{\text{NN}}} = 200$  GeV. The results for  $\sigma_{R_{\Psi_2}}^{-2}$  and  $\Delta\gamma/v_2$  from the AVFD model are shown for background and for signal + background as indicated. The Hijing model results are only shown for the background. The dashed lines represent linear fits to the background values.

about the RP,  $\phi$  denote the particle azimuthal emission angles,  $\alpha, \beta$  denote the electric charge (+) or (-) and SS and OS represent same-sign (++, --) and opposite-sign (+-) charges. Measurements of the quotient  $\Delta\gamma/v_2$  with the 2<sup>nd</sup>-order anisotropy coefficient  $v_2$ , are usually employed to aid quantification of the background-driven charge separation.

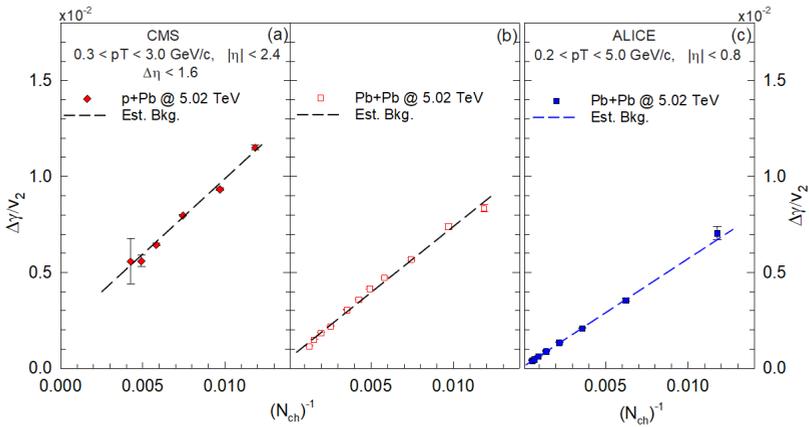
The  $R_{\Psi_2}(\Delta S)$  correlator measures charge separation relative to  $\Psi_2$  via the ratio:  $R_{\Psi_2}(\Delta S) = C_{\Psi_2}(\Delta S)/C_{\Psi_2}^{\perp}(\Delta S)$ , where  $C_{\Psi_2}(\Delta S)$  and  $C_{\Psi_2}^{\perp}(\Delta S)$  are correlation functions that quantify charge separation  $\Delta S$ , approximately parallel and perpendicular (respectively) to the  $\vec{B}$ -field. The charge-shuffling procedure used to construct the correlation functions ensures identical properties for their numerator and denominator, except for the charge-dependent correlations, which are of interest [4].  $C_{\Psi_2}(\Delta S)$  measures both CME- and background-driven charge separation while  $C_{\Psi_2}^{\perp}(\Delta S)$  measures only the background. After correcting the  $R_{\Psi_2}(\Delta S)$  distributions for the effects of particle-number fluctuations and the event-plane resolution, their inverse variance  $\sigma_{R_{\Psi_2}}^{-2}$  are used to quantify the charge separation [4].

In this work, we use model simulations to chart the scaling properties of  $\sigma_{R_{\Psi_2}}^{-2}$  and  $\Delta\gamma/v_2$  for the background and signal + background, respectively, in A+A collisions. We then leverage these scaling properties to identify and characterize a possible CME-driven charge separation using previously published data for  $p$ +Au,  $d$ +Au, Ru+Ru, Zr+Zr and Au+Au collisions at RHIC [5–10], and  $p$ +Pb and Pb+Pb collisions at the LHC [11–14].

Figure 1 shows the results for  $\sigma_{R_{\Psi_2}}^{-2}$  and  $\Delta\gamma/v_2$  obtained with the AVFD and Hijing models for Au+Au collisions. Note that these models emphasize different sources for the charge-dependent non-flow background; the initial axial charge density  $n_5/s$  and the degree of local charge conservation (LCC) regulate the magnitude of the CME- and background-driven charge separation in the AVFD model. The solid triangles in Fig. 1 show that the background scales as  $1/N_{\text{ch}}$  – the expected trend for the charge-dependent non-flow correlations. By contrast, the signal (Sig.) + background values (solid diamonds) indicate positive deviations from the background scaling [16, 17]. This dependence can be represented as;  $\Delta\gamma/v_2 = a + b/(N_{\text{ch}})^{1-c}$ , and  $\sigma_{R_{\Psi_2}}^{-2} = a' + b'/(N_{\text{ch}})^{1-c'}$ , for the small values of  $n_5/s$  indicated in Fig. 1. Here,  $a, b$  and  $c$  are parameters;  $c$  characterizes the degree of the scaling violation.

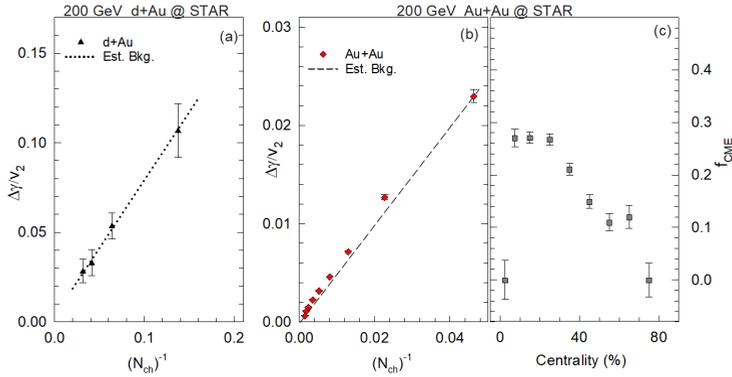
Note that for  $c \sim 0$  the  $1/N_{\text{ch}}$  scaling for the background is retrieved, as demonstrated with the AVFD model in Fig. 1.

The scaling violation gives a direct signature of the CME-driven contributions to the charge separation (Figs. 1 (a) and (c)). It can be quantified via the fraction:  $f_{\text{CME}}^{\Delta\gamma} = [\Delta\gamma/v_2(\text{Sig.} + \text{Bkg.}) - \Delta\gamma/v_2(\text{Bkg.})]/[\Delta\gamma/v_2(\text{Sig.} + \text{Bkg.})]$ ,  $f_{\text{CME}}^R = [\sigma_{R\psi_2}^{-2}(\text{Sig.} + \text{Bkg.}) - \sigma_{R\psi_2}^{-2}(\text{Bkg.})]/[\sigma_{R\psi_2}^{-2}(\text{Sig.} + \text{Bkg.})]$ . The scaling patterns in Fig. 1 suggest that the observation of  $1/N_{\text{ch}}$  scaling for the experimental  $\sigma_{R\psi_2}^{-2}$  and  $\Delta\gamma/v_2$  measurements would strongly indicate background-driven charge separation with little room for a CME contribution. However, observing a violation of this  $1/N_{\text{ch}}$  scaling would indicate the CME-driven contribution. Figs. 1 (a) and (c) also indicate comparable background and signal + background  $\sigma_{R\psi_2}^{-2}$  and  $\Delta\gamma/v_2$  values in central and peripheral collisions, suggesting that the background dominates over that of the CME-driven contributions in these collisions. Note the reduction of  $\vec{B}$  in central collisions and the enhanced de-correlation between the event plane and the  $\vec{B}$ -field in peripheral collisions. Since the background dominates in central and peripheral collisions, the  $\sigma_{R\psi_2}^{-2}$  and  $\Delta\gamma/v_2$  measurements for these collisions can be leveraged with  $1/N_{\text{ch}}$  scaling to obtain a quantitative estimate of the background over the entire centrality span (cf. Fig. 1). Here, an important proviso is to experimentally establish that the background in  $p(d)+A$  and  $A+A$  collisions scale over the full centrality span.



**Figure 2.**  $\Delta\gamma/v_2$  vs.  $1/N_{\text{ch}}$  for  $p+Pb$  (a) and  $Pb+Pb$  [(c) and (d)] collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. The dashed lines indicate an estimate of the background. The data are taken from Refs. [12, 13, 15].

The  $v_2$  and  $\Delta\gamma$  values reported for  $p+Au$ ,  $d+Au$ ,  $Ru+Ru$ ,  $Zr+Zr$  and  $Au+Au$  collisions at RHIC [5–10], and  $p+Pb$  and  $Pb+Pb$  collisions at the LHC [12–15] were used to investigate the scaling properties of  $\Delta\gamma/v_2$ . Fig. 2 shows the results for  $p+Pb$  and  $Pb+Pb$  collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. They indicate that  $\Delta\gamma/v_2$  essentially scales as  $1/N_{\text{ch}}$  ( $c \approx 0$ ), suggesting negligible CME contributions in these collisions. They also confirm that the combined sources of background (LCC, resonances, back-to-back jets, ...) which should be substantial, especially for  $p+Pb$ , scale as  $1/N_{\text{ch}}$ . Note as well that the CME contribution is negligible in  $p(d)+A$  collisions because of significant reductions in  $\vec{B}$ , and the sizable de-correlation between the event plane and the  $\vec{B}$ -field [12]. Thus, the scaling patterns of  $\Delta\gamma/v_2$  for these systems' sizable backgrounds give a direct experimental constraint on the validity of  $1/N_{\text{ch}}$  scaling of the background.



**Figure 3.**  $\Delta\gamma/v_2$  vs.  $1/N_{ch}$  [(a) and (b)] and  $f_{CME}$  vs. centrality (c) for  $d+Au$  and  $Au+Au$  collisions at  $\sqrt{s_{NN}} = 200$  GeV. The dotted and dashed lines indicate an estimate of the background contributions. The data are taken from Refs. [7, 9? ].

The scaling results for  $Au+Au$  collisions at  $\sqrt{s_{NN}} = 200$  GeV are shown in Fig. 3. The  $1/N_{ch}$  scaling apparent for  $d+Au$  collisions (Fig. 3 (a)) confirms the expectation that the CME is negligible in these collisions. It also confirms that the combined sources of background (LCC, resonances, back-to-back jets, ...) which could be substantial in  $d+Au$  collisions, show  $1/N_{ch}$  scaling. In contrast to  $d+Au$ , the results for  $Au+Au$  (Fig. 3(b)) show visible indications of a violation ( $c > 0$ ) to the  $1/N_{ch}$  scaling observed for background-driven charge separation in  $p(d)+A$  collisions. Similar violations were observed for  $Ru+Ru$  and  $Zr+Zr$  [17]. The scaling violation is similar to that observed for signal + background in Figs. 1 (a) and (c), suggesting an unambiguous non-negligible CME contribution to the measured  $\Delta\gamma/v_2$  in  $Au+Au$ ,  $Ru+Ru$ , and  $Zr+Zr$  collisions. The estimates of the background for all three systems are obtained by leveraging the  $\Delta\gamma/v_2$  measurements for peripheral and central collisions with  $1/N_{ch}$  scaling [17]. Here, it is noteworthy that the simulated results from the AVFD and HIJING models, as well as the measurements presented in Figs. 2 and 3(a), provide strong constraints that the combined sources of background, scale as  $1/N_{ch}$  over the full centrality span. The background estimates were used to extract  $f_{CME}$  values for  $Au+Au$ , (Fig. 3 (c))  $Ru+Ru$  and  $Zr+Zr$  collisions respectively. They indicate non-negligible  $f_{CME}$  values that vary with centrality. In mid-central collisions,  $f_{CME} \sim 27\%$  for  $Au+Au$  collisions, which is roughly a factor of two larger than the values for  $Ru+Ru$  and  $Zr+Zr$ . Within the uncertainties, no significant difference between the values for  $Ru+Ru$  and  $Zr+Zr$  was observed, suggesting that  $\Delta\gamma/v_2$  is sensitive to CME-driven charge separation in  $Ru+Ru$  and  $Zr+Zr$  collisions but may be insensitive to the signal difference between them [17].

In summary, the scaling properties of the  $R_{\Psi_2}(\Delta S)$  and the  $\Delta\gamma$  correlators have been used to characterize the CME in several colliding systems at RHIC and the LHC. The results for  $p+Au$  and  $d+Au$  collisions at  $\sqrt{s_{NN}} = 200$  GeV and  $p+Pb$  ( $\sqrt{s_{NN}} = 5.02$  TeV) and  $Pb+Pb$  collisions at  $\sqrt{s_{NN}} = 5.02$  and 2.76 TeV, scales as  $1/N_{ch}$  consistent with background-driven charge separation. However, the results for  $Au+Au$ ,  $Ru+Ru$  and  $Zr+Zr$  collisions ( $\sqrt{s_{NN}} = 200$  GeV) show scaling violations which indicate a CME-driven contribution in the presence of significant background. In mid-central collisions,  $f_{CME} \sim 27\%$  for  $Au+Au$  collisions and approximately a factor of two smaller in  $Ru+Ru$  and  $Zr+Zr$  collisions but with similar magnitudes for the two isobars.

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