

# Approaching maximal precision of Hong-Ou-Mandel interferometry with non-perfect visibility

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**Abstract.** This work explores precision limits in two-photon Hong-Ou-Mandel interferometry under non-perfect visibility. A theoretical model is developed and experimentally validated using different quantum states. A remarkable ratio of 0.97 between the experimental precision and the quantum limit is observed, establishing a new benchmark in the field.

## 1 Introduction

In quantum mechanics, the precision achieved in parameter estimation using a quantum state as a probe is determined by the measurement strategy employed. The quantum precision limit, a fundamental boundary, is defined by the intrinsic characteristics of the state and its dynamics. Theoretical results have revealed that in interference measurements with two possible outcomes, like the Hong-Ou-Mandel interference, this limit can be reached under ideal conditions of perfect visibility and zero losses [1]. However, in practice, this cannot be achieved, so precision never reaches the quantum limit. But how do experimental setups approach precision limits under realistic circumstances? In this work, we provide a general model for precision limits in two-photon Hong-Ou-Mandel interferometry for non-perfect visibility and validate it experimentally using different quantum states. We reveal the existence of an optimal scenario, where we observe a remarkable ratio of 0.97 between the experimental precision and the quantum limit, establishing a new benchmark in the field [2].

## 2 Main Text

### 2.1 Theoretical model

In a typical metrological protocol, a probe evolves dynamically depending on a parameter  $\theta$  to be estimated, and its measurement outcome is used to estimate  $\theta$ . Precision is bounded by the Cramer-Rao bound  $\delta\theta \geq 1/\sqrt{NF(\theta)}$ , where  $F(\theta)$  is the Fisher information (FI) and  $N$  is the number of measurements. When using quantum resources, like individual photons in a Hong-Ou-Mandel (HOM) experiment, one can define the quantum Fisher information (QFI). In this case, precision is limited by the quantum Cramer Rao (QCR)  $\delta\theta \geq 1/\sqrt{V\mathcal{F}}$ , where  $\mathcal{F}$  is

the QFI. The QCR bound represents the quantum precision limit. Finding an experimental measurement strategy where  $F(\theta) = \mathcal{F}$  is not straightforward.

In the HOM experiment, the QCR bound can be reached under perfect visibility. Attempts to reach this bound have led to remarkable precision in time delay  $\tau$  estimations [3, 4]. The probe used is a biphoton state  $|\psi\rangle = \iint d\omega_1 d\omega_2 f(\omega_1, \omega_2) a_1^\dagger(\omega_1) a_2^\dagger(\omega_2) |0\rangle$  where  $f(\omega_1, \omega_2)$  is the joint spectral amplitude. We model the dependency of the coincidence probability  $P_c(\tau)$  with the HOM visibility as:

$$P_c(\tau) = \frac{1}{2} - \frac{V}{2} W(0, \tau) \quad (1)$$

where  $W(0, \tau)$  is the Wigner function associated with the function  $f_-(\omega_-)$ , defined by :  $f(\omega_1, \omega_2) = f_-(\omega_-) f_+(\omega_+)$ , with  $\omega_\pm = \omega_1 \pm \omega_2$ . The FI is then given by:

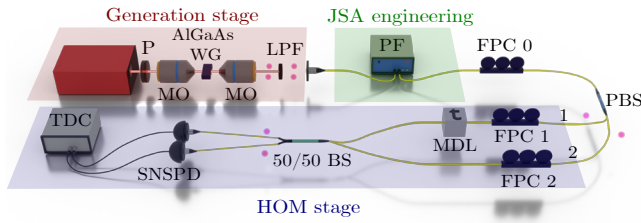
$$F(V, \tau) = V^2 \frac{(W'(0, \tau))^2}{1 - V^2 W^2(0, \tau)} \quad (2)$$

In this work, we are interested in understanding how  $\tilde{F}_V = \max_\tau F(V, \tau)$  approaches  $\mathcal{F}$  as a function of the visibility  $V$  for different quantum states. Eq.2 allows us to identify a relation with the effective phase space occupation of the state's Wigner function: indeed, the scaling of  $\tilde{F}_V$  with  $V$  is connected to how far the quantum state is from saturating the time-frequency Heisenberg uncertainty principle. For this reason, while Gaussian states minimize the FI's scaling with visibility, states with large effective areas in time frequency phase space, like Schrödinger cat (SC)-like states, maximize it, making them remarkably robust in HOM-based quantum metrology.

### 2.2 Experimental demonstration

We validate our model by implementing a HOM interferometer using two-photon states described by various

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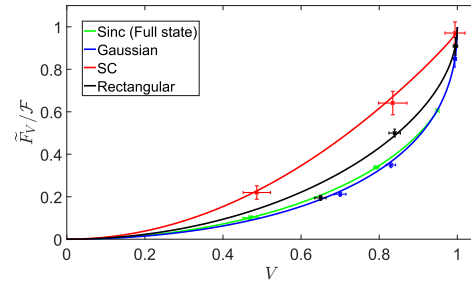


**Figure 1.** Experimental setup for investigating the metrological performance of the Hong-Ou-Mandel (HOM) experiment.

functions  $f_-(\omega_-)$ , demonstrating different scaling behaviors with visibility  $V$  (Fig. 1). An AlGaAs Bragg reflection waveguide [5] is used to generate polarization-entangled photon pairs via type II Spontaneous Parametric Down-Conversion (SPDC) at telecom wavelengths. The joint spectral amplitude (JSA) is engineered by postmanipulation through a programmable filter. Our study includes three different filter shapes: a rectangular filter, a Gaussian filter, and a combination of two rectangular filters corresponding to energy-matched channels, allowing the engineering of a Schrodinger Cat (SC)-like state. Photon pairs, leaving the filter, are separated by a polarizing beam splitter (PBS) and enter the HOM interferometer. Control over polarization distinguishability and temporal delay is done with fibered polarization controllers (FPC) and a motorized optical delay line (MDL), respectively. Photons are recombined at a 50/50 beam splitter (BS) and detected by superconducting nanowires single photon detectors (SNSPD), with temporal correlations analyzed by a time-to-digital converter (TDC).

The experimental results (points) are reported in (Fig2), together with the theoretical prediction (lines). Remarkably high visibilities exceeding 99% are achieved with all the engineered states. The SC-like state exhibits the most favorable scaling behavior, while the Gaussian state dis-

plays a less optimal one. For a visibility value of about 99.4%, the ratio drops from 0.97 for the SC-like state to 0.85 for the Gaussian state, demonstrating the effectiveness of our model across different state configurations.



**Figure 2.** Scaling of the ratio  $\tilde{F}_V/\mathcal{F}$  for different biphoton states with respect to the HOM interferogram visibility  $V$ .

## References

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