

# Top-quark pair production in Pb+Pb collisions in the ATLAS experiment

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**Abstract.** In relativistic heavy-ion collisions, top quarks are expected to be attractive candidates for probing the quark-gluon plasma as well as to bring unique information about the time evolution of strongly interacting matter. We report the first study of top-quark pair production in lead–lead collisions at the centre-of-mass energy of 5.02 TeV with the ATLAS experiment at the LHC. The dataset was recorded in 2015 and 2018, amounting to an integrated luminosity of  $1.9 \text{ nb}^{-1}$ . Top-quark pair production cross section is studied in the  $e\mu$  channel. The result is compared to theory predictions based on different nuclear PDF sets.

## 1 Introduction

The top quark is the only elementary particle, along with the Higgs boson, that has yet to be directly observed in lead–lead (Pb+Pb) collisions with a significance exceeding five standard deviations. As the heaviest particle carrying colour charge, the top quark is expected to provide novel experimental insights into the quark-gluon plasma (QGP) [1]. In particular, hadronically decaying  $W$  bosons from top-quark decays could serve as probes of the time evolution of the QGP [2]. An observation of  $t\bar{t}$  production in Pb+Pb collisions represents a critical initial milestone toward enabling more detailed investigations of the QGP.

The full Pb+Pb dataset recorded with the ATLAS detector [3] during the Run 2 campaign allows for the first observation of top-quark pair production with a significance of five standard deviations [4]. The data were obtained in 2015 and 2018 at a nucleon–nucleon centre-of-mass energy of  $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$ , corresponding to the total integrated luminosity of  $0.5 \text{ nb}^{-1}$  and  $1.4 \text{ nb}^{-1}$ , respectively. The data were collected under low-pileup conditions with an average number of interactions per bunch crossing of  $\langle\mu\rangle = 1.76 \cdot 10^{-3}$  and  $\langle\mu\rangle = 2.58 \cdot 10^{-3}$  in 2015 and 2018, respectively.

## 2 Event selection

Top-quark pairs candidates are reconstructed using events with exactly one electron and one muon. Single-electron and single-muon triggers are employed, with minimum  $p_{\text{T}}$  requirements of 15 and 8 GeV, respectively [5]. Exactly one primary vertex is required, reconstructed

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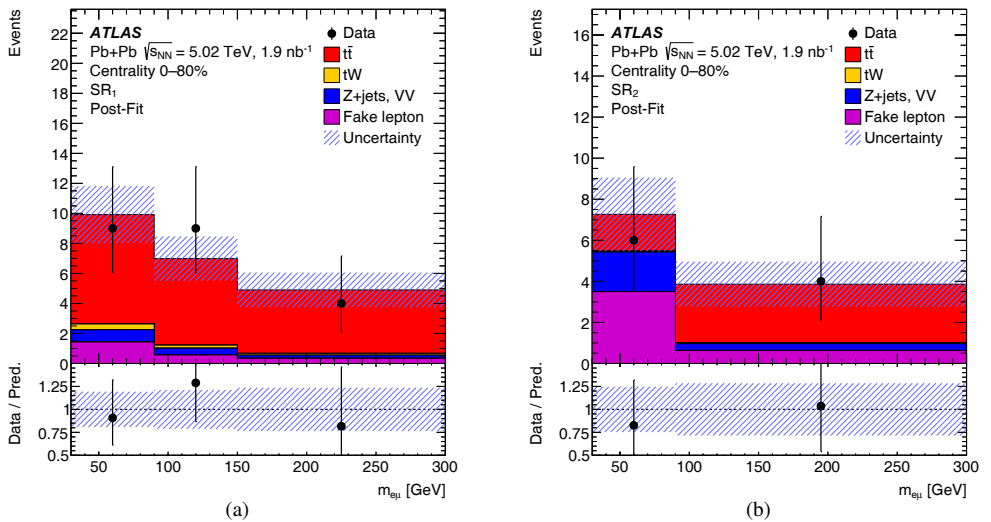
from at least two good-quality charged-particle tracks with  $p_T > 0.5$  GeV. Electron (muon) candidates must have  $p_T > 18$  (15) GeV and  $|\eta| < 2.47$  (2.5), and satisfy ‘Loose’ identification and isolation criteria [6, 7]. Jets are reconstructed using the anti- $k_r$  algorithm [8] with a radius parameter  $R = 0.4$ , and are required to have  $p_T > 35$  GeV and  $|\eta| < 2.5$ . The background energy from the underlying event is subtracted on an event-by-event basis. A fake-lepton background contribution is evaluated from data using data-driven techniques.

Events in the electron–muon ( $e\mu$ ) channel of the  $t\bar{t}$  decay are studied in the measurement. Only events in the 0–80% centrality interval are used, to avoid contributions from photon-induced processes. Final states containing exactly one oppositely charged  $e\mu$  pair are selected. The invariant mass of the  $e\mu$  pair,  $m_{e\mu}$ , has to be above 30 GeV, to minimise the fake-lepton background contribution. At least two jets are required in a given event. Two signal regions (SR<sub>1</sub> and SR<sub>2</sub>) are constructed, where events with dilepton transverse momentum  $p_T^{e\mu} > 40$  GeV form SR<sub>1</sub>, and the remaining events constitute SR<sub>2</sub>.

### 3 Analysis strategy

The signal strength,  $\mu_{t\bar{t}}$ , is defined as the ratio of the observed  $t\bar{t}$  cross section to the Standard Model expectation with no nuclear parton distribution function (nPDF) effects involved. A profile-likelihood method [9] is used to extract  $\mu_{t\bar{t}}$  using the  $e\mu$  channel of  $t\bar{t}$  decays. The value of  $\mu_{t\bar{t}}$  is obtained by the simultaneous fit to the  $m_{e\mu}$  distributions in the two signal regions. Figure 1 shows a comparison between the  $m_{e\mu}$  distributions in data and a sum of predictions in SR<sub>1</sub> and SR<sub>2</sub> after the fit.

Systematic uncertainties affecting the measurement arise from electron, muon, and jet reconstruction, fake-lepton background estimation, the signal and background modelling, and the integrated luminosity. The main sources of systematic uncertainty originate from the modelling of the  $t\bar{t}$  process and jet performance. The total relative systematic uncertainty in the measurement amounts to 18%.



**Figure 1.** Post-fit distributions of the the  $m_{e\mu}$  variable in (a) SR<sub>1</sub> and (d) SR<sub>2</sub>, with total uncertainties represented by the hatched area. The bottom panels show ratios of data and a sum of predictions [4].

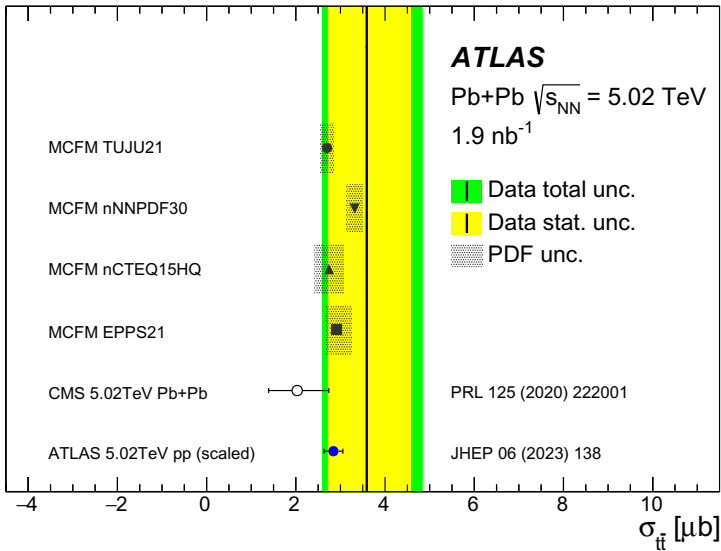
## 4 Results

The measured  $\mu_{t\bar{t}}$  value is converted into the inclusive  $t\bar{t}$  production cross section in Pb+Pb collisions in the 0-100% centrality interval. The measured  $t\bar{t}$  cross section in Pb+Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV is

$$\sigma_{t\bar{t}}^{\text{Pb+Pb}} = 3.6^{+1.0}_{-0.9} \text{ (stat.) }^{+0.8}_{-0.5} \text{ (syst.) } \mu\text{b} = 3.6^{+1.2}_{-1.0} \text{ (tot.) } \mu\text{b}. \quad (1)$$

The total relative uncertainty amounts to 31% and is dominated by the statistical component of 26%. The observed significance of the measurement is 5.0 standard deviations, leading to the first observation of the  $t\bar{t}$  process in Pb+Pb collisions.

Figure 2 shows a comparison of the obtained  $t\bar{t}$  cross section with other measurements and theoretical models. Good agreement is found with predictions based on four nPDF sets: TUJU21 [10], nNNPDF3.0 [11], nCTEQ15HQ [12], and EPPS21 [13]. The result is consistent within the total uncertainties with the  $t\bar{t}$  cross section reported by CMS [14]. The measurement is also in agreement with the  $t\bar{t}$  cross section in  $pp$  collisions at  $\sqrt{s} = 5.02$  TeV by ATLAS [15], scaled to the Pb+Pb system by the lead mass number squared,  $A_{\text{Pb}}^2$  ( $A_{\text{Pb}} = 208$ ).



**Figure 2.** Comparison of the  $t\bar{t}$  cross section with the evidence reported by CMS [14] and the scaled measurement in  $pp$  collisions by ATLAS [15], along with theoretical predictions based on four nPDF sets [4].

## 5 Conclusions

The  $t\bar{t}$  process is observed in the  $e\mu$  channel in Pb+Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV with the significance of 5.0 standard deviations for the first time. The inclusive  $t\bar{t}$  cross section is measured with the total relative uncertainty of 31%, resulting in the most precise measurement of the  $t\bar{t}$  process in Pb+Pb collisions to date.

This result paves a new way for further studies of QGP properties at the LHC. The first observation of the  $t\bar{t}$  process establishes analysis methods for top-quark production in heavy-ion collisions and opens a possibility of studying the QGP time evolution via hadronically decaying  $W$  bosons from top-quark decays.

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